quatrième série - tome 47 fascicule 3 mai-juin 2014

> ANNALES
> SCIENTIFIQUES de $L^{\prime}$ ÉCOLE NORMALE SUPÉRIEURE

Steve HOFMANN \& José María MARTELL
Uniform rectifiability and barmonic measure I: Uniform rectifiability implies Poisson kernels in $L^{p}$

# UNIFORM RECTIFIABILITY AND HARMONIC MEASURE I: UNIFORM RECTIFIABILITY IMPLIES POISSON KERNELS IN $L^{p}$ 

by Steve HOFMANN and José María MARTELL


#### Abstract

We present a higher dimensional, scale-invariant version of a classical theorem of F. and M. Riesz [37]. More precisely, we establish scale invariant absolute continuity of harmonic measure with respect to surface measure, along with higher integrability of the Poisson kernel, for a domain $\Omega \subset \mathbb{R}^{n+1}$, $n \geq 2$, with a uniformly rectifiable boundary, which satisfies the Harnack chain condition plus an interior (but not exterior) Corkscrew condition. In a companion paper to this one [28], we also establish a converse, in which we deduce uniform rectifiability of the boundary, assuming scale invariant $L^{q}$ bounds, with $q>1$, on the Poisson kernel.


Résumé. - On présente une version invariante par échelles et en dimension supérieure à 3 , d'un théorème classique de F. et M. Riesz [37]. Plus précisément, on établit l'absolue continuité de la mesure harmonique par rapport à la mesure de surface, ainsi qu'un gain d'intégrabilité pour le noyau de Poisson, pour un domaine $\Omega \subset \mathbb{R}^{n+1}, n \geq 2$, à bord uniformément rectifiable, vérifiant une condition de chaîne de Harnack et une condition de type «points d'ancrage» ou «Corkscrew» intérieure (mais pas extérieure). L'article associé [28] établit une réciproque, c'est-à-dire l'uniforme rectifiabilité du bord en supposant des estimées invariantes par échelle $L^{q}$ pour $q>1$ sur le noyau de Poisson.

## 1. Introduction

In [37], F. and M. Riesz showed that for a simply connected domain in the complex plane with a rectifiable boundary, harmonic measure is absolutely continuous with respect to arc length measure. A quantitative version of this theorem was obtained by Lavrentiev [35]. More generally, if only a portion of the boundary is rectifiable, Bishop and Jones [8] have shown that harmonic measure is absolutely continuous with respect to arclength on that portion.

[^0]They also present a counter-example to show that the result of [37] may fail in the absence of some topological hypothesis (e.g., simple connectedness).

In this paper we extend the results of [37] and [35] to higher dimensions, without imposing extra assumptions on either the exterior domain or the boundary, as has been done previously. Our extension (Theorem 1.26 below) is "scale-invariant", i.e., assuming scaleinvariant analogues of the hypotheses of [37], we show that harmonic measure satisfies a scale-invariant version of absolute continuity, namely the weak- $A_{\infty}$ condition (cf. Definition 1.19 below). More precisely, let $\Omega \subset \mathbb{R}^{n+1}, n \geq 2$, be a connected, open set. We establish the weak $-A_{\infty}$ property of harmonic measure, assuming that $\partial \Omega$ is uniformly rectifiable (cf. (1.13) below), and that $\Omega$ satisfies interior (but not necessarily exterior) Corkscrew and Harnack chain conditions (cf. Definitions 1.4 and 1.6 below). Uniform rectifiability is the scale-invariant version of rectifiability, while the Corkscrew and Harnack chain conditions are scale invariant analogues of the topological properties of openness and path connectedness, respectively. We emphasize that in contrast to previous work in this area in dimensions $n+1 \geq 3$, we impose no restriction on the geometry of the exterior domain $\Omega_{\text {ext }}:=\mathbb{R}^{n+1} \backslash \bar{\Omega}$, nor any extra condition on the geometry of the boundary, beyond uniform rectifiability. In particular, we do not require that any component of $\Omega_{\text {ext }}$ satisfy a Corkscrew condition (as in [29], [39, 6]) or even an $n$-disk condition as in [17]; nor do we assume that $\partial \Omega$ contains "Big Pieces" of the boundaries of Lipschitz sub-domains of $\Omega$, as in [7]. The absence of such assumptions is the main advance in the present paper.

In addition, in a companion paper to this one [28], written jointly with I. Uriarte-Tuero, we establish a converse, Theorem 1.28 , in which we deduce uniform rectifiability of the boundary, given a certain scale invariant local $L^{q}$ estimate, with $q>1$, for the Poisson kernel (cf. (1.24)). The method of proof in [28] may be of independent interest, as it entails a novel use of " $T b$ " theory to obtain a free boundary result.

Taken together, the main results of the present paper and of [28], namely Theorems 1.26 and 1.28 below, may be summarized as follows (the terminology and notation used in the statement will be clarified or cross-referenced immediately afterwards):

THEOREM 1.1. - Let $\Omega \subset \mathbb{R}^{n+1}$, $n \geq 2$, be a connected open set which satisfies interior Corkscrew and Harnack chain conditions, and whose boundary $\partial \Omega$ is n-dimensional AhlforsDavid regular. Then the following are equivalent:

1. $\partial \Omega$ is uniformly rectifiable.
2. For every surface ball $\Delta=\Delta(x, r) \subset \partial \Omega$, with radius $r \lesssim \operatorname{diam} \partial \Omega$, the harmonic measure $\omega^{X_{\Delta}} \in$ weak $-A_{\infty}(\Delta)$.
3. $\omega \ll \sigma$, and there is $a q>1$ such that the Poisson kernel $k^{X_{\Delta}}$ satisfies the scale invariant $L^{q}$ bound (1.24), for every $\Delta=\Delta(x, r) \subset \partial \Omega$, with radius $r \lesssim \operatorname{diam} \partial \Omega$.

Remark 1.2. - By the counter-example of [8], one would not expect to obtain the implication $(1) \Longrightarrow(2)$, without some sort of connectivity assumption; for us, the interior Harnack chain condition plays this role.

Given a domain $\Omega \subset \mathbb{R}^{n+1}$, a "surface ball" is a set $\Delta=\Delta(x, r):=B(x, r) \cap \partial \Omega$, where $x \in \partial \Omega$, and $B(x, r)$ denotes the standard $(n+1)$-dimensional Euclidean ball of radius $r$ centered at $x$. For such a surface ball $\Delta$, we let $\omega^{X_{\Delta}}$ denote harmonic measure for $\Omega$, with pole at the
$4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - No 3
"Corkscrew point" $X_{\Delta}$ (see Definition 1.4). The Corkscrew and Harnack chain conditions, as well as the notions of Ahlfors-David regularity (ADR), uniform rectifiability (UR) and weak $-A_{\infty}$, are described in Definitions 1.4, 1.6, 1.7, 1.9, and 1.19 below.

The present paper treats the direction (1) implies (2). That (2) implies (3) is well known (see the discussion following Definition 1.19). The main result in [28] is that (3) implies (1). We mention also that we obtain in the present paper an extension of (1) implies (2), in which our hypotheses are assumed to hold only in an "interior big pieces" sense (cf. Definition 1.14 and Theorem 1.27 below).

To place Theorem 1.1 in context, we review previous related work in dimension $n+1 \geq 3$. We recall that in [29], the authors introduce the notion of a "non-tangentially accessible" (NTA) domain: $\Omega$ is said to be NTA if it satisfies the Corkscrew and Harnack chain conditions ("interior Corkscrew and Harnack chain conditions"), and also if the exterior domain, $\Omega_{\text {ext }}:=\mathbb{R}^{n+1} \backslash \bar{\Omega}$ (which need not be connected), satisfies the Corkscrew condition ("exterior Corkscrew condition"). The latter was relaxed in [17] to allow a sort of "weak exterior Corkscrew" condition in which the analogue of the exterior Corkscrew point is the center merely of an $n$-dimensional disk in $\Omega_{\text {ext }}$, rather than of a full Euclidean ball. A key observation made in [17] was that the weak exterior Corkscrew condition is still enough to obtain local Hölder continuity at the boundary of harmonic functions which vanish on a surface ball. In [17], the authors prove that, in the presence of Ahlfors-David regularity of the boundary, the NTA condition of [29] or even its relaxed version with "weak exterior Corkscrews", implies that $\Omega$ satisfies an "interior big pieces" of Lipschitz sub-domains condition (cf. Definition 1.14 below). By a simple maximum principle argument (plus the deep result of [15]), one then almost immediately obtains a certain lower bound for harmonic measure, to wit that there are constants $\eta \in(0,1)$ and $c_{0}>0$ such that for each surface ball $\Delta \subset \partial \Omega$, and any Borel subset $A \subset \Delta$, we have

$$
\begin{equation*}
\omega^{X_{\Delta}}(A) \geq c_{0}, \quad \text { whenever } \sigma(A) \geq \eta \sigma(\Delta) \tag{1.3}
\end{equation*}
$$

In turn, still given NTA, or at least the relaxed version of [17], the latter bound self-improves to an $A_{\infty}$ estimate for harmonic measure, via the comparison principle. The same $A_{\infty}$ conclusion was also obtained by a different argument in [39], under the full NTA condition of [29]. In [7], the authors impose an interior Corkscrew condition, but in lieu of the Harnack chain and exterior (or weak exterior) Corkscrew conditions, the authors assume instead the consequence of these conditions deduced in [17], namely, that $\Omega$ satisfies the aforementioned condition concerning "interior big pieces" of Lipschitz sub-domains. The bound (1.3) (suitably interpreted) then holds almost immediately (again by the maximum principle), but the self-improvement argument, in the absence of the Harnack chain and exterior (or weak exterior) Corkscrew conditions, is now more problematic (indeed, the usual proofs of the comparison principle rely on Harnack's inequality and local Hölder continuity at the boundary), and the authors conclude in [7] only that $\omega$ is weak- $A_{\infty}$. On the other hand, they give an example to show that this conclusion is best possible (that is, they construct a domain which satisfies the "interior big pieces" condition, but whose harmonic measure fails to be doubling). We mention also in this context the recent paper [6], in which the geometric conclusion of [17], namely the existence of "interior big pieces" of Lipschitz sub-domains, is shown to hold assuming the full NTA condition (with two-sided Corkscrews), but in
which only the lower (but not the upper) bound is required in the Ahlfors-David condition (cf. (1.8)).

In the present paper, we improve the results of [7] and of [17] by removing the "big pieces of Lipschitz sub-domains" hypothesis, as well as all assumptions regarding the exterior domain. That is, in Theorem 1.26, we assume only that $\Omega$ satisfies interior Corkscrew and Harnack chain conditions, and that its boundary is uniformly rectifiable. More generally, in Theorem 1.27, we suppose only that these hypotheses hold in an appropriate "interior big pieces" sense (in particular, our results include those of [7] as a special case, since their Lipschitz sub-domains clearly satisfy our hypotheses). The difficulty now, and the heart of the proof, is to establish (1.3); with the latter in hand, the self-improvement to weak $A_{\infty}$ proceeds as in [7]. We mention that by an unpublished example of Hrycak, UR does not, in general, imply big pieces of Lipschitz graphs ${ }^{(1)}$ (that the opposite implication does hold for ADR sets is easy, and well known). Moreover, in [28] we obtain a converse which shows that the UR hypothesis is optimal. In this connection, we mention also the following observation, which was brought to our attention by M. Badger and T. Toro. Let $F \subset \mathbb{R}^{2}$ denote the " 4 corners Cantor set" of J. Garnett (see, e.g., [19, p. 4]), and let $F^{*}:=F \times \mathbb{R} \subset \mathbb{R}^{3}$ be the "cylinder" above $F$. Then $\Omega:=\mathbb{R}^{3} \backslash F^{*}$ satisfies the (interior) Corkscrew and Harnack chain conditions, and has a 2-dimensional ADR boundary, but the boundary is not UR, and therefore its harmonic measure is not weak- $A_{\infty}$.

We conclude this historical survey by providing some additional context for our work here and in [28], namely, that our results may be viewed as a "large constant" analogue of the work of Kenig and Toro $[32,33,34]$. The latter, taken collectively, say that in the presence of a Reifenberg flatness condition and Ahlfors-David regularity, one has that $\log k \in \mathrm{VMO}$ if and only if $\nu \in \mathrm{VMO}$, where $k$ is the Poisson kernel with pole at some fixed point, and $\nu$ is the unit normal to the boundary. Moreover, given the same background hypotheses, the condition that $\nu \in \mathrm{VMO}$ is equivalent to a uniform rectifiability (UR) condition with vanishing trace, thus $\log k \in \mathrm{VMO} \Longleftrightarrow$ vanishing $U R$. On the other hand, our large constant version "almost" says " $\log k \in \mathrm{BMO} \Longleftrightarrow U R$ ", given interior Corkscrews and Harnack chains. Indeed, it is well known that the $A_{\infty}$ condition (i.e., weak- $A_{\infty}$ plus the doubling property) implies that $\log k \in \mathrm{BMO}$, while if $\log k \in \mathrm{BMO}$ with small norm, then $k \in A_{\infty}$.

In order to state our results precisely, we shall first need to discuss some preliminary matters.

### 1.1. Notation and definitions

- We use the letters $c, C$ to denote harmless positive constants, not necessarily the same at each occurrence, which depend only on dimension and the constants appearing in the hypotheses of the theorems (which we refer to as the "allowable parameters"). We shall also

[^1]$4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - $\mathrm{N}^{\mathrm{o}} 3$
sometimes write $a \lesssim b$ and $a \approx b$ to mean, respectively, that $a \leq C b$ and $0<c \leq a / b \leq C$, where the constants $c$ and $C$ are as above, unless explicitly noted to the contrary. At times, we shall designate by $M$ a particular constant whose value will remain unchanged throughout the proof of a given lemma or proposition, but which may have a different value during the proof of a different lemma or proposition.

- Given a domain $\Omega \subset \mathbb{R}^{n+1}$, we shall use lower case letters $x, y, z$, etc., to denote points on $\partial \Omega$, and capital letters $X, Y, Z$, etc., to denote generic points in $\mathbb{R}^{n+1}$ (especially those in $\mathbb{R}^{n+1} \backslash \partial \Omega$ ).
- The open $(n+1)$-dimensional Euclidean ball of radius $r$ will be denoted $B(x, r)$ when the center $x$ lies on $\partial \Omega$, or $B(X, r)$ when the center $X \in \mathbb{R}^{n+1} \backslash \partial \Omega$. A "surface ball" is denoted $\Delta(x, r):=B(x, r) \cap \partial \Omega$.
- Given a Euclidean ball $B$ or surface ball $\Delta$, its radius will be denoted $r_{B}$ or $r_{\Delta}$, respectively.
- Given a Euclidean or surface ball $B=B(X, r)$ or $\Delta=\Delta(x, r)$, its concentric dilate by a factor of $\kappa>0$ will be denoted by $\kappa B:=B(X, \kappa r)$ or $\kappa \Delta:=\Delta(x, \kappa r)$.
- For $X \in \mathbb{R}^{n+1}$, we set $\delta(X):=\operatorname{dist}(X, \partial \Omega)$.
- We let $H^{n}$ denote $n$-dimensional Hausdorff measure, and let $\sigma:=\left.H^{n}\right|_{\partial \Omega}$ denote the "surface measure" on $\partial \Omega$.
- For a Borel set $A \subset \mathbb{R}^{n+1}$, we let $1_{A}$ denote the usual indicator function of $A$, i.e., $1_{A}(x)=1$ if $x \in A$, and $1_{A}(x)=0$ if $x \notin A$.
- For a Borel set $A \subset \mathbb{R}^{n+1}$, we let $\operatorname{int}(A)$ denote the interior of $A$. If $A \subset \partial \Omega$, then $\operatorname{int}(A)$ will denote the relative interior, i.e., the largest relatively open set in $\partial \Omega$ contained in $A$. Thus, for $A \subset \partial \Omega$, the boundary is then well defined by $\partial A:=\bar{A} \backslash \operatorname{int}(A)$.
- For a Borel set $A$, we denote by $\mathscr{C}(A)$ the space of continuous functions on $A$, by $\mathscr{C}_{c}(A)$ the subspace of $\mathscr{C}(A)$ with compact support in $A$, and by $\mathscr{C}_{b}(A)$ the space of bounded continuous functions on $A$. If $A$ is unbounded, we denote by $\mathscr{C}_{0}(A)$ the space of continuous functions on $A$ converging to 0 at infinity.
- For a Borel subset $A \subset \partial \Omega$, we set $f_{A} f d \sigma:=\sigma(A)^{-1} \int_{A} f d \sigma$.
- We shall use the letter $I$ (and sometimes $J$ ) to denote a closed $(n+1)$-dimensional Euclidean cube with sides parallel to the co-ordinate axes, and we let $\ell(I)$ denote the side length of $I$. We use $Q$ to denote a dyadic "cube" on $\partial \Omega$. The latter exists, given that $\partial \Omega$ is ADR (cf. $[18,13]$ ), and enjoy certain properties which we enumerate in Lemma ?? below.

Definition 1.4 (Corkscrew condition). - Following [29], we say that a domain $\Omega \subset \mathbb{R}^{n+1}$ satisfies the "Corkscrew condition" if for some uniform constant $c>0$ and for every surface ball $\Delta:=\Delta(x, r)$, with $x \in \partial \Omega$ and $0<r<\operatorname{diam}(\partial \Omega)$, there is a ball $B\left(X_{\Delta}, c r\right) \subset B(x, r) \cap \Omega$. The point $X_{\Delta} \subset \Omega$ is called a "Corkscrew point" relative to $\Delta$. We note that we may allow $r<C \operatorname{diam}(\partial \Omega)$ for any fixed $C$, simply by adjusting the constant $c$.

Remark 1.5. - We note that, on the other hand, every $X \in \Omega$, with $\delta(X)<\operatorname{diam}(\partial \Omega)$, may be viewed as a Corkscrew point, relative to some surface ball $\Delta \subset \partial \Omega$. Indeed, set $r=K \delta(X)$, with $K>1$, fix $x \in \partial \Omega$ such that $|X-x|=\delta(X)$, and let $\Delta:=\Delta(x, r)$.

Definition 1.6 (Harnack chain condition). - Again following [29], we say that $\Omega$ satisfies the Harnack chain condition if there is a uniform constant $C$ such that for every $\rho>0$, $\Lambda \geq 1$, and every pair of points $X, X^{\prime} \in \Omega$ with $\delta(X), \delta\left(X^{\prime}\right) \geq \rho$ and $\left|X-X^{\prime}\right|<\Lambda \rho$, there is a chain of open balls $B_{1}, \ldots, B_{N} \subset \Omega, N \leq C(\Lambda)$, with $X \in B_{1}, X^{\prime} \in B_{N}, B_{k} \cap B_{k+1} \neq \varnothing$ and $C^{-1} \operatorname{diam}\left(B_{k}\right) \leq \operatorname{dist}\left(B_{k}, \partial \Omega\right) \leq C \operatorname{diam}\left(B_{k}\right)$. The chain of balls is called a "Harnack chain".

We remark that the Corkscrew condition is a quantitative, scale invariant version of the fact that $\Omega$ is open, and the Harnack chain condition is a scale invariant version of path connectedness.

Definition 1.7 (Ahlfors-David regular). - We say that a closed set $E \subset \mathbb{R}^{n+1}$ is $n$-dimensional ADR (or simply ADR) ("Ahlfors-David regular") if there is some uniform constant $C$ such that

$$
\begin{equation*}
\frac{1}{C} r^{n} \leq H^{n}(E \cap B(x, r)) \leq C r^{n}, \forall r \in\left(0, R_{0}\right), x \in E \tag{1.8}
\end{equation*}
$$

where $R_{0}$ is the diameter of $E$ (which may be infinite). When $E=\partial \Omega$, the boundary of a domain $\Omega$, we shall sometimes for convenience simply say that " $\Omega$ has the ADR property" to mean that $\partial \Omega$ is ADR.

Definition 1.9 (Uniform Rectifiability). - Following David and Semmes [18, 19], we say that a closed set $E \subset \mathbb{R}^{n+1}$ is $n$-dimensional UR (or simply UR) ("Uniformly Rectifiable"), if it satisfies the ADR condition (1.8), and if for some uniform constant $C$ and for every Euclidean ball $B:=B\left(x_{0}, r\right), r \leq \operatorname{diam}(E)$, centered at any point $x_{0} \in E$, we have the Carleson measure estimate

$$
\begin{equation*}
\iint_{B}\left|\nabla^{2} \& 1(X)\right|^{2} \operatorname{dist}(X, E) d X \leq C r^{n} \tag{1.10}
\end{equation*}
$$

where $\& f$ is the single layer potential of $f$, i.e.,

$$
\begin{equation*}
\& f(X):=c_{n} \int_{E}|X-y|^{1-n} f(y) d H^{n}(y) \tag{1.11}
\end{equation*}
$$

Here, the normalizing constant $c_{n}$ is chosen so that $\mathscr{E}(X):=c_{n}|X|^{1-n}$ is the usual fundamental solution for the Laplacian in $\mathbb{R}^{n+1}$. When $E=\partial \Omega$, the boundary of a domain $\Omega$, we shall sometimes for convenience simply say that " $\Omega$ has the UR property" to mean that $\partial \Omega$ is UR.

We note that there are numerous characterizations of uniform rectifiability given in [18, 19]; the one stated above will be most useful for our purposes, and appears in [19, Chapter 3, Part III]. We remark that the UR sets are precisely those for which all "sufficiently nice" singular integrals are bounded on $L^{2}$ (see [18]).

We recall that "Uniform Rectifiability" is the scale invariant analogue of rectifiability; in particular, using an idea of P. Jones [30], one may derive, for UR sets, a quantitative version
of the fact that rectifiability may be characterized in terms of existence a.e. of approximate tangent planes. For $x \in E, t>0$, we set

$$
\begin{equation*}
\beta_{2}(x, t) \equiv \inf _{P}\left(\frac{1}{t^{n}} \int_{B(x, t) \cap E}\left(\frac{\operatorname{dist}(y, P)}{t}\right)^{2} d H^{n}(y)\right)^{1 / 2} \tag{1.12}
\end{equation*}
$$

where the infimum runs over all $n$-planes $P$. Then a closed, ADR set $E$ is UR if and only if the following Carleson measure estimate holds on $E \times \mathbb{R}_{+}$:

$$
\begin{equation*}
\sup _{x_{0} \in E, r>0} r^{-n} \int_{0}^{r} \int_{B\left(x_{0}, t\right) \cap E} \beta_{2}(x, t)^{2} d H^{n}(x) \frac{d t}{t}<\infty \tag{1.13}
\end{equation*}
$$

Again see [18] for details.

Definition 1.14 ("Interior Big Pieces"). - Given a domain $\Omega \subset \mathbb{R}^{n+1}$, with ADR boundary, and a collection $\&$ of domains in $\mathbb{R}^{n+1}$, we say that $\Omega$ has "interior big pieces of $\delta$ " (denoted $\Omega \in I B P(\phi)$ ) if there are constants $\alpha>0, K>1$ such that for every $X \in \Omega$, with $\delta(X)<\operatorname{diam}(\partial \Omega)$, there is a point $x \in \partial \Omega$, with $|x-X|=\delta(X)$, and a domain $\Omega^{\prime} \in \varnothing$ for which, with $r:=K \delta(X)$, we have

1. $\Omega^{\prime} \subset \Omega$.
2. $H^{n}\left(\partial \Omega^{\prime} \cap \Delta(x, r)\right) \geq \alpha H^{n}(\Delta(x, r)) \approx \alpha r^{n}$.
3. $X$ is a Corkscrew point for $\Omega^{\prime}$, relative to $\Delta_{\star}(y, 2 r):=B(y, 2 r) \cap \partial \Omega^{\prime}$, for some $y \in \partial \Omega^{\prime} \cap \Delta$ (we note that $X$ is also a Corkscrew point for $\Omega$, relative to $\Delta$, by construction; cf. Remark 1.5).

Lemma 1.15 (Existence and properties of the "dyadic grid" $[18,19,13]$ )
Suppose that $E \subset \mathbb{R}^{n+1}$ satisfies the $A D R$ condition (1.8). Then there exist constants $a_{0}>0$, $\eta>0$ and $C_{1}<\infty$, depending only on dimension and the $A D R$ constants, such that for each $k \in \mathbb{Z}$, there is a collection of Borel sets ("cubes")

$$
\mathbb{D}_{k}:=\left\{Q_{j}^{k} \subset E: j \in \mathfrak{I}_{k}\right\}
$$

where $\mathfrak{I}_{k}$ denotes some (possibly finite) index set depending on $k$, satisfying
(i) $E=\bigcup_{j} Q_{j}^{k}$ for each $k \in \mathbb{Z}$.
(ii) If $m \geq k$ then either $Q_{i}^{m} \subset Q_{j}^{k}$ or $Q_{i}^{m} \cap Q_{j}^{k}=\varnothing$.
(iii) For each $(j, k)$ and each $m<k$, there is a unique $m$ such that $Q_{j}^{k} \subset Q_{i}^{m}$.
(iv) Diameter $\left(Q_{j}^{k}\right) \leq C_{1} 2^{-k}$.
(v) Each $Q_{j}^{k}$ contains some "surface ball" $\Delta\left(x_{j}^{k}, a_{0} 2^{-k}\right):=B\left(x_{j}^{k}, a_{0} 2^{-k}\right) \cap E$.
(vi) $H^{n}\left(\left\{x \in Q_{j}^{k}: \operatorname{dist}\left(x, E \backslash Q_{j}^{k}\right) \leq \tau 2^{-k}\right\}\right) \leq C_{1} \tau^{\eta} H^{n}\left(Q_{j}^{k}\right)$, for all $k, j$ and for all $\tau \in\left(0, a_{0}\right)$.

A few remarks are in order concerning this lemma.

- In the setting of a general space of homogeneous type, this lemma has been proved by Christ [13]. In that setting, the dyadic parameter $1 / 2$ should be replaced by some constant $\delta \in(0,1)$. It is a routine matter to verify that one may take $\delta=1 / 2$ in the presence of the Ahlfors-David property (1.8) (in this more restrictive context, the result already appears in $[18,19])$.
- For our purposes, we may ignore those $k \in \mathbb{Z}$ such that $2^{-k} \gtrsim \operatorname{diam}(E)$, in the case that the latter is finite.
- We shall denote by $\mathbb{D}=\mathbb{D}(E)$ the collection of all relevant $Q_{j}^{k}$, i.e.,

$$
\mathbb{D}:=\bigcup_{k} \mathbb{D}_{k}
$$

where, if $\operatorname{diam}(E)$ is finite, the union runs over those $k$ such that $2^{-k} \lesssim \operatorname{diam}(E)$.

- Properties (iv) and (v) imply that for each cube $Q \in \mathbb{D}_{k}$, there is a point $x_{Q} \in E$, a Euclidean ball $B\left(x_{Q}, r\right)$ and a surface ball $\Delta\left(x_{Q}, r\right):=B\left(x_{Q}, r\right) \cap E$ such that $r \approx 2^{-k} \approx \operatorname{diam}(Q)$ and

$$
\begin{equation*}
\Delta\left(x_{Q}, r\right) \subset Q \subset \Delta\left(x_{Q}, C r\right) \tag{1.16}
\end{equation*}
$$

for some uniform constant $C$. We shall denote this ball and surface ball by

$$
\begin{equation*}
B_{Q}:=B\left(x_{Q}, r\right), \quad \Delta_{Q}:=\Delta\left(x_{Q}, r\right), \tag{1.17}
\end{equation*}
$$

and we shall refer to the point $x_{Q}$ as the "center" of $Q$.

- Let us now specialize to the case that $E=\partial \Omega$, with $\Omega$ satisfying the Corkscrew condition. Given $Q \in \mathbb{D}(\partial \Omega)$, we shall sometimes refer to a "Corkscrew point relative to $Q$ ", which we denote by $X_{Q}$, and which we define to be the Corkscrew point $X_{\Delta}$ relative to the ball $\Delta:=\Delta_{Q}$ (cf. (1.16), (1.17) and Definition 1.4). We note that

$$
\begin{equation*}
\delta\left(X_{Q}\right) \approx \operatorname{dist}\left(X_{Q}, Q\right) \approx \operatorname{diam}(Q) \tag{1.18}
\end{equation*}
$$

- For a dyadic cube $Q \in \mathbb{D}_{k}$, we shall set $\ell(Q)=2^{-k}$, and we shall refer to this quantity as the "length" of $Q$. Evidently, $\ell(Q) \approx \operatorname{diam}(Q)$.
- For a dyadic cube $Q \in \mathbb{D}$, we let $k(Q)$ denote the "dyadic generation" to which $Q$ belongs, i.e., we set $k=k(Q)$ if $Q \in \mathbb{D}_{k}$; thus, $\ell(Q)=2^{-k(Q)}$.

Definition $1.19\left(A_{\infty}, A_{\infty}^{\text {dyadic }}\right.$ and weak- $\left.A_{\infty}\right)$. - Given a surface ball $\Delta=B \cap \partial \Omega$, a Borel measure $\omega$ defined on $\partial \Omega$ is said to belong to the class $A_{\infty}(\Delta)$ if there are positive constants $C$ and $\theta$ such that for every $\Delta^{\prime}=B^{\prime} \cap \partial \Omega$ with $B^{\prime} \subseteq B$, and every Borel set $F \subset \Delta^{\prime}$, we have

$$
\begin{equation*}
\omega(F) \leq C\left(\frac{\sigma(F)}{\sigma\left(\Delta^{\prime}\right)}\right)^{\theta} \omega\left(\Delta^{\prime}\right) \tag{1.20}
\end{equation*}
$$

If we replace the surface balls $\Delta$ and $\Delta^{\prime}$ by a dyadic cube $Q$ and its dyadic subcubes $Q^{\prime}$, with $F \subset Q^{\prime}$, then we say that $\omega \in A_{\infty}^{\text {dyadic }}(Q)$ :

$$
\begin{equation*}
\omega(F) \leq C\left(\frac{\sigma(F)}{\sigma\left(Q^{\prime}\right)}\right)^{\theta} \omega\left(Q^{\prime}\right) \tag{1.21}
\end{equation*}
$$

Similarly, $\omega \in$ weak- $A_{\infty}(\Delta)$, with $\Delta=B \cap \partial \Omega$, if for every $\Delta^{\prime}=B^{\prime} \cap \partial \Omega$ with $2 B^{\prime} \subseteq B$, we have

$$
\begin{equation*}
\omega(F) \leq C\left(\frac{\sigma(F)}{\sigma\left(\Delta^{\prime}\right)}\right)^{\theta} \omega\left(2 \Delta^{\prime}\right) \tag{1.22}
\end{equation*}
$$

As is well known [14, 21], [38], the $A_{\infty}$ (resp. weak- $A_{\infty}$ ) condition is equivalent to the property that the measure $\omega$ is absolutely continuous with respect to $\sigma$, and that its density satisfies a reverse Hölder (resp. weak reverse Hölder) condition. In this paper, we are interested in the case that $\omega=\omega^{X}$, the harmonic measure with pole at $X$. In that setting, we let $k^{X}:=d \omega^{X} / d \sigma$ denote the Poisson kernel, so that (1.20) is equivalent to the reverse Hölder estimate

$$
\begin{equation*}
\left(f_{\Delta^{\prime}}\left(k^{X}\right)^{q} d \sigma\right)^{1 / q} \leq C f_{\Delta^{\prime}} k^{X} d \sigma \tag{1.23}
\end{equation*}
$$

for some $q>1$ and for some uniform constant $C$. In particular, when $\Delta^{\prime}=\Delta$, and $X=X_{\Delta}$, a Corkscrew point relative to $\Delta$, the latter estimate reduces to

$$
\begin{equation*}
\int_{\Delta}\left(k^{X_{\Delta}}\right)^{q} d \sigma \leq C \sigma(\Delta)^{1-q} \tag{1.24}
\end{equation*}
$$

Similarly, (1.22) is equivalent to

$$
\begin{equation*}
\left(f_{\Delta^{\prime}}\left(k^{X}\right)^{q} d \sigma\right)^{1 / q} \leq C f_{2 \Delta^{\prime}} k^{X} d \sigma \tag{1.25}
\end{equation*}
$$

Assuming that the latter bound holds with $\Delta^{\prime}=\Delta$, and with $X=X_{\Delta}$, then one again obtains (1.24).

### 1.2. Statement of the main results

Our main results are as follows. We shall use the terminology that a connected open set $\Omega \subset \mathbb{R}^{n+1}$ is a $l$-sided NTA domain if it satisfies interior (but not necessarily exterior) Corkscrew and Harnack chain conditions ${ }^{(2)}$.

THEOREM 1.26. - Let $\Omega \subset \mathbb{R}^{n+1}$, $n \geq 2$, be a 1 -sided NTA domain whose boundary $\partial \Omega$ is n-dimensional $U R$. Then for each surface ball $\Delta$, the harmonic measure $\omega^{X_{\Delta}}$ belongs to weak $-A_{\infty}(\Delta)$, with uniform weak $-A_{\infty}$ constants depending only on dimension and on the constants in the ADR, UR, Corkscrew and Harnack chain conditions.

We emphasize again that we impose no hypothesis (as in [29, 39], [17]) on the geometry of the exterior domain, nor do we assume as in [7] that the boundary has "Big Pieces" of boundaries of Lipschitz subdomains of $\Omega$.

We shall also obtain a certain "self-improvement" of Theorem 1.26 , in which the hypotheses are assumed to hold only in an appropriate "big pieces" sense.

[^2]Theorem 1.27. - Let $\Omega \subset \mathbb{R}^{n+1}, n \geq 2$, be a connected open set whose boundary $\partial \Omega$ is $n$-dimensional ADR. Suppose further that $\Omega \in \operatorname{IBP}(\$)$ (cf. Definition 1.14), where $\&$ is a collection of 1 -sided NTA domains with UR boundaries, with uniform control of all of the relevant Corkscrew, Harnack chain, $A D R$ and UR constants. Then for each surface ball $\Delta=\Delta(x, r)$, and for every $X \in \Omega \backslash B(x, r)$, the harmonic measure $\omega^{X}$ belongs to weak $-A_{\infty}(\Delta)$, with uniform weak- $A_{\infty}$ constants that depend only on dimension, on the constants in the ADR and interior big pieces conditions, and on the relevant constants for the subdomains.

Remark. - We note that in Theorem 1.27, we have obtained that $\omega^{X}$ belongs to weak- $A_{\infty}(\Delta(x, r))$, for all $X \in \Omega \backslash B(x, r)$. In the presence of the Harnack chain Condition, as in Theorem 1.26, one may obtain the same conclusion for $X=X_{\Delta}$, the Corkscrew point relative to $\Delta$. On the other hand, in Theorem 1.26, we of course also obtain that $\omega^{X}$ belongs to weak- $A_{\infty}(\Delta(x, r))$, for all $X \in \Omega \backslash B(x, r)$.

In a companion paper to this one [28], we shall establish the converse to Theorem 1.26:
Theorem 1.28. - Let $\Omega \subset \mathbb{R}^{n+1}, n \geq 2$, be a 1 -sided NTA domain, whose boundary is $n$-dimensional ADR. Suppose also that harmonic measure $\omega$ is absolutely continuous with respect to surface measure and that there is a $q>1$ such that for every surface ball $\Delta=\Delta(x, r)$ with radius $r \lesssim \operatorname{diam} \partial \Omega$, the Poisson kernel satisfies the scale invariant estimate (1.24). Then $\partial \Omega$ is $U R$.

We also mention that in [28] we obtain a "big pieces" version of the previous result in the following sense. Let $E \subset \mathbb{R}^{n+1}$ be a closed set and assume that $E$ is $n$-dimensional ADR. Assume that there exists $q>1$ such that $E$ has "big pieces of boundaries of $\phi$ " (i.e., for every surface ball $B(x, r) \cap E$ there is $\Omega^{\prime} \in \&$ whose boundary has an "ample" contact with $E \cap B(x, r)$ ), where $\&$ is a collection of domains $\Omega^{\prime}$ each of them satisfying the hypotheses of Theorem 1.28 (with $q$ fixed) and with uniform control on the relevant constants. Then $E$ is UR. See [28] for the precise statement.

## Acknowledgements

The first named author wishes to thank John Lewis for helpful comments concerning the paper [7]. He also thanks Misha Safonov and Tatiana Toro for bringing to our attention the work of Aikawa [1, 2].

## 2. Outline of the strategy of the proof

Let us sketch the strategy of the proofs of Theorems 1.26 and 1.27. We shall do most of our analysis in certain approximating domains which enjoy additional qualitative properties. Given these qualitative properties, we shall prove some a priori estimates for the Green function $G$ and for harmonic measure $\omega$, beginning with Lemma 3.30 in Section 3, whose proofs rely on being able to "hide" certain small quantities, which must therefore be known in advance to be finite. An interesting feature of these a priori estimates is that they permit us to deduce the doubling property for $\omega$, as well as a comparison principle for $G$, in the absence of an exterior disk or Corkscrew condition (the exterior conditions enable one to
$4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - No 3
prove boundary Hölder continuity of solutions vanishing on a surface ball). We obtain these properties for $G$ and $\omega$ without establishing boundary Hölder continuity. We note that, by the work of Aikawa [1, 2], some of the preliminary estimates that we prove in Section 3, in particular, the "Carleson estimate" Lemma 3.37, and the Comparison Principle (aka "Boundary Harnack Principle") Lemma 3.64, are known, but we include our own relatively short proofs here for the sake of self-containment.

We also establish several geometric preliminaries as follows. In Section 4, we use the Harnack chain property to prove a Poincaré inequality (Lemma 4.8), which we use in turn, in Section 5, to obtain a criterion for the existence of exterior Corkscrew points in the complement of certain "sawtooth" regions (Lemma 5.10). This criterion stipulates that the Carleson measure (cf. (1.10))

$$
\begin{equation*}
\left|\nabla^{2} \& 1(X)\right|^{2} \operatorname{dist}(X, \partial \Omega) d X \tag{2.1}
\end{equation*}
$$

be sufficiently small in the relevant sawtooth region. We then present in Section 6 a variant of the "sawtooth lemma" of [16] (Lemma 6.15), which roughly speaking allows for a comparison, in the sense of $A_{\infty}$, between the respective harmonic measures, $\omega$ and $\omega_{\Omega_{\mathcal{F}}}$, for the original domain and for the sawtooth domain (more precisely, our version of the sawtooth lemma allows us to transfer the dyadic $A_{\infty}$ property of $\omega_{\Omega_{\mathcal{g}}}$ to $\mathscr{P}_{g} \omega$, where $\mathscr{P}_{\mathscr{G}}$ is a sort of "conditional expectation" projection operator, with respect to some collection $\mathcal{F}$ of nonoverlapping dyadic cubes from which the sawtooth was constructed). The arguments of Section 6 are an extension, to the present context, of our previous work in the Euclidean setting [27].

With these preliminary matters in hand, we proceed to the heart of our proof, which will exploit the technique of "extrapolation (i.e., bootstrapping) of Carleson measures", as it appears in our previous work [27] (see also [26]), but originating in [12] and [36]. We now describe the application of this technique in our setting. By a Corona type stopping time construction delineated in Section 7, plus an induction scheme (formalized in Lemma 8.5), we reduce matters to verifying that $\mathscr{P}_{\mathcal{G}} \omega$ (that is, the projection of harmonic measure mentioned above) enjoys the dyadic $A_{\infty}$ property, in sawtooth domains $\Omega_{\mathscr{F}}$ in which the Carleson measure (2.1) has sufficiently small Carleson norm. In turn, we establish this property for $\mathscr{P}_{\mathscr{G}} \omega$, by using the preliminary facts noted above: by the smallness of (2.1) in the sawtooth, we deduce that the complement of the sawtooth enjoys an exterior Corkscrew condition. Thus, we may apply the results of [17] to the sawtooth, to obtain that $\omega_{\Omega_{g}}$, the harmonic measure for the sawtooth domain, belongs to $A_{\infty}$ with respect to surface measure on the boundary of the sawtooth. Then, invoking our version of the sawtooth lemma, we find that $\mathscr{P}_{g} \omega$ belongs to dyadic $A_{\infty}$, as desired. The "extrapolation" technology (i.e., Lemma 8.5) now allows us to conclude that $\omega$ belongs to $A_{\infty}$ with respect to surface measure, in a local, but scale invariant way. However, at this point, we have only reached this conclusion in our approximating domains $\Omega_{N}$, albeit with $A_{\infty}$ constants independent of $N$. Here $\left\{\Omega_{N}\right\}$ is a nested increasing sequence of sub-domains of $\Omega$, each of which enjoys the qualitative properties mentioned above, such that $\Omega_{N} \nearrow \Omega$. It is not clear whether the $A_{\infty}$ property of harmonic measure, or even the doubling property, are transmitted in the limit to harmonic measure on $\Omega$. However, a maximum principle argument (in the case of Theorem 1.27, there are two separate maximum principle arguments) allows us to transfer, at least, the property that there are uniform
constants $c_{0}, \eta \in(0,1)$ such that for any Borel subset $A \subset \Delta$,

$$
\begin{equation*}
\sigma(A)>\eta \sigma(\Delta) \quad \Longrightarrow \quad \omega^{X_{\Delta}}(A) \geq c_{0} \tag{*}
\end{equation*}
$$

The fact that $(*)$ holds, in the absence of assumptions on the exterior domain $\Omega_{\text {ext }}$ or on $\partial \Omega$ (beyond UR), is really the main result of this paper. Given $(*)$, we obtain the conclusion of Theorems 1.26 and 1.27 by invoking the arguments of [7].

## 3. Some fundamental estimates

In this section we recall or establish certain fundamental estimates for harmonic measure and the Green function. In the sequel, $\Omega \subset \mathbb{R}^{n+1}, n \geq 2$, will be a connected, open set, $\omega^{X}$ will denote harmonic measure for $\Omega$, with pole at $X$, and $G(X, Y)$ will be the Green function. At least in the case that $\Omega$ is bounded, we may, as usual, define $\omega^{X}$ via the maximum principle and the Riesz representation theorem, after first using the method of Perron (see, e.g., [22, pp. 24-25]) to construct a harmonic function "associated" to arbitrary continuous boundary data. ${ }^{(3)}$ For unbounded $\Omega$, we may still define harmonic measure via a standard approximation scheme as follows. Given $R>0$, set $\Omega_{R}:=\Omega \cap B\left(x_{0}, 2 R\right)$, where $x_{0}$ is a fixed point on $\partial \Omega$. Define a smooth cut-off function $\eta \in \mathscr{C}_{0}^{\infty}([-2,2])$, with $0 \leq \eta \leq 1$, $\eta \equiv 1$ on $[-1,1]$, and $\eta$ monotone decreasing on $(1,2)$ and monotone increasing on $(-2,-1)$. Suppose now that $0 \leq f \in \mathscr{C}_{b}(\partial \Omega)$ and set

$$
\begin{equation*}
f_{R}(x):=f(x) \eta\left(\frac{\left|x-x_{0}\right|}{R}\right) \tag{3.1}
\end{equation*}
$$

Extending $f_{R}$ to be zero outside of its support defines a continuous function on $\partial \Omega_{R}$, so we may construct the corresponding Perron solution $u_{R}$ in $\Omega_{R}$. By the maximum principle,

$$
u_{R} \leq u_{R^{\prime}} \text { in } \Omega_{R}, \text { if } R^{\prime}>R, \quad \text { and } \quad \sup _{\Omega_{R}} u_{R} \leq \sup _{\partial \Omega_{R}} f_{R} \leq \sup _{\partial \Omega} f
$$

Consequently, by Harnack's convergence theorem ([22, p. 22]), there is a harmonic function $u$ in $\Omega$ such that

$$
\begin{equation*}
\lim _{R \rightarrow \infty} u_{R}=u \tag{3.2}
\end{equation*}
$$

with the convergence being uniform on compacta in $\Omega$. Moreover, $u$ satisfies the maximum principle

$$
\sup _{\Omega} u \leq \sup _{\partial \Omega} f
$$

Thus, we may again define harmonic measure $\omega^{X}$ for $X \in \Omega$ via the Riesz representation theorem. We note for future reference that $\omega^{X}$ is a non-negative, finite Borel measure which satisfies the outer regularity property

$$
\begin{equation*}
\omega^{X}(A):=\inf _{A \subseteq O} \omega^{X}(O) \tag{3.3}
\end{equation*}
$$

for every Borel set $A \subset \partial \Omega$, where the infimum runs over all (relatively) open $O \subset \partial \Omega$ containing $A$.

[^3]The Green function may now be constructed by setting

$$
\begin{equation*}
G(X, Y):=\mathcal{E}(X-Y)-\int_{\partial \Omega} \mathcal{E}(X-z) d \omega^{Y}(z) \tag{3.4}
\end{equation*}
$$

where $\mathcal{E}(X):=c_{n}|X|^{1-n}$ is the usual fundamental solution for the Laplacian in $\mathbb{R}^{n+1}$. We choose the normalization that makes $\mathscr{E}$ positive. Given this normalization, we shall also have that $G \geq 0$ (cf. Lemma 3.11 below.)

Before proceeding further, let us note one more fact for future reference. Assuming that $\Omega$ is unbounded, and using the notation above, let $\omega_{R}^{X}$ and $G_{R}(X, Y)$ denote, respectively, harmonic measure and Green's function for the approximating domain $\Omega_{R}$. We then have

$$
\begin{equation*}
\lim _{R \rightarrow \infty} G_{R}(X, Y)=G(X, Y) \tag{3.5}
\end{equation*}
$$

with the convergence being uniform on compacta in $\Omega$, in the $Y$ variable with $X \in \Omega$ fixed. Indeed, fixing $X$, choosing $R$ so large that $R \gg\left|X-x_{0}\right|$, and setting $f:=\mathcal{E}(X-\cdot)$, with $f_{R}$ defined as in (3.1), we have that

$$
\int_{\partial \Omega_{R}} f d \omega_{R}^{Y}=\int_{\partial \Omega_{R}} f_{R} d \omega_{R}^{Y}+O\left(R^{1-n}\right):=u_{R}(Y)+O\left(R^{1-n}\right) .
$$

We then obtain (3.5) immediately from (3.2) and the definition of the Green function (3.4).
Lemma 3.6 (Bourgain [9]). - Suppose that $\partial \Omega$ is $n$-dimensional ADR. Then there are uniform constants $c \in(0,1)$ and $C \in(1, \infty)$, such that for every $x \in \partial \Omega$, and every $r \in(0, \operatorname{diam}(\partial \Omega))$, if $Y \in \Omega \cap B(x, c r)$, then

$$
\begin{equation*}
\omega^{Y}(\Delta(x, r)) \geq 1 / C>0 . \tag{3.7}
\end{equation*}
$$

In particular, if $\Omega$ satisfies the Corkscrew and Harnack chain conditions, then for every surface ball $\Delta$, we have

$$
\begin{equation*}
\omega^{X_{\Delta}}(\Delta) \geq 1 / C>0 \tag{3.8}
\end{equation*}
$$

We refer the reader to [9, Lemma 1] for the proof.
We next introduce some notation. We say that a domain $\Omega$ satisfies the qualitative exterior Corkscrew condition if there exists $N \gg 1$ such that $\Omega$ has exterior Corkscrew points at all scales smaller than $2^{-N}$. That is, there exists a constant $c_{N}$ such that for every surface ball $\Delta=\Delta(x, r)$, with $x \in \partial \Omega$ and $r \leq 2^{-N}$, there is a ball $B\left(X_{\Delta}^{\text {ext }}, c_{N} r\right) \subset B(x, r) \cap \Omega_{\text {ext }}$.

Given a ball $B_{0}$ centered on $\partial \Omega$, and $X \in \Omega \backslash B_{0}$, we also introduce the quantity

$$
\begin{equation*}
\Upsilon_{B_{0}}(X):=\sup _{B: 2 B \subseteq B_{0}} \frac{r_{\Delta}^{1-n} \omega^{X}(\Delta)}{G\left(X_{\Delta}, X\right)}, \tag{3.9}
\end{equation*}
$$

where the sup runs over all the balls $B$ centered at $\partial \Omega$ with $2 B \subseteq B_{0}$ and where as usual $\Delta=B \cap \partial \Omega$. We also set $\left\|\Upsilon_{B_{0}}\right\|=\sup _{X \in \Omega \backslash B_{0}} \Upsilon_{B_{0}}(X)$. The quantity $\Upsilon_{B_{0}}$ will enter in the proof of Lemma 3.30 below.

Remark 3.10. - Let us observe that if $\Omega$ satisfies the qualitative exterior Corkscrew condition, then every point in $\partial \Omega$ is regular in the sense of Wiener. Moreover, for 1-sided NTA domains, the qualitative exterior Corkscrew points allow local Hölder continuity at the boundary (albeit with bounds which may depend badly on $N$ ), so that the program of [29] may be followed to prove that $\Upsilon_{B_{0}}(X)$ is a priori finite (possibly depending on $N, X$
and $B_{0}$ ). Eventually, we shall apply Lemmas 3.11 and 3.30 below (and several related lemmas and corollaries) to certain approximating domains $\Omega_{N}$ which will inherit the stated quantitative hypotheses from the original domain $\Omega$, but which also satisfy the qualitative exterior Corkscrew conditions for scales $\lesssim 2^{-N}$. We emphasize that all of the quantitative bounds that we shall establish will depend only upon dimension and on the parameters in the 1 -sided NTA and UR (including ADR) conditions, and thus these bounds will hold uniformly for the entire family of approximating domains.

Lemma 3.11. - There are positive, finite constants $C$, depending only on dimension, and $c(n, \theta)$, depending on dimension and $\theta \in(0,1)$, such that the Green function satisfies

$$
\begin{gather*}
G(X, Y) \leq C|X-Y|^{1-n}  \tag{3.12}\\
c(n, \theta)|X-Y|^{1-n} \leq G(X, Y), \quad \text { if }|X-Y| \leq \theta \delta(X), \theta \in(0,1) \tag{3.13}
\end{gather*}
$$

Moreover, if every point on $\partial \Omega$ is regular in the sense of Wiener, then

$$
\begin{align*}
& G(X, Y) \geq 0, \quad \forall X, Y \in \Omega, X \neq Y  \tag{3.14}\\
& G(X, Y)=G(Y, X), \quad \forall X, Y \in \Omega, X \neq Y \tag{3.15}
\end{align*}
$$

and

$$
\begin{equation*}
\int_{\partial \Omega} \Phi d \omega^{X}=-\iint_{\Omega} \nabla_{Y} G(Y, X) \cdot \nabla \Phi(Y) d Y \tag{3.16}
\end{equation*}
$$

for every $X \in \Omega$ and $\Phi \in C_{0}^{\infty}\left(\mathbb{R}^{n+1}\right)$ with $\Phi(X)=0$.

Proof. - Some of these facts are standard, but we include the simple proof here. Recall that we have chosen the normalization $\mathcal{E}(X):=c_{n}|X|^{1-n}$ with $c_{n}>0$. Inequality (3.12) is then trivial, by Definition (3.4), since $\int_{\partial \Omega} \mathcal{E}(X-z) d \omega^{Y}(z) \geq 0$. We now consider (3.13). Suppose that $0<\theta<1$, and that $|X-Y| \leq \theta \delta(X)$. Then,

$$
\int_{\partial \Omega}|X-z|^{1-n} d \omega^{Y}(z) \leq \delta(X)^{1-n} \leq \theta^{n-1}|X-Y|^{1-n}
$$

Thus, $G(X, Y) \geq c_{n}\left(1-\theta^{n-1}\right)|X-Y|^{1-n}$, as desired.
We now assume that every boundary point is regular in the sense of Wiener. Let us prove (3.14). Suppose first that $\Omega$ is bounded. Fix $X \in \Omega$, and observe that by (3.13), it is enough to consider the case that $Y \in \Omega^{\prime}:=\Omega \backslash \overline{B(X, \delta(X) / 2)}$. Moreover, by (3.13), we have in particular that $G(X, \cdot)>0$ on $\partial B(X, \delta(X) / 2)$. On the other hand, since every boundary point is regular, we have by Definition (3.4) that $G(X, \cdot) \equiv 0$ on $\partial \Omega$. Applying the maximum principle in $\Omega^{\prime}$, we then obtain (3.14), at least when $\Omega$ is bounded. If $\Omega$ is unbounded, we may invoke (3.5).

Next, we establish the symmetry condition (3.15), again assuming that every boundary point is regular in the sense of Wiener. By (3.5), it is enough to treat the case that $\Omega$ is bounded. Specializing to the case of the Laplacian, the Green function constructed in [23], which we denote temporarily by $\widetilde{G}(X, Y)$, is symmetric (see [23, Theorem 1.3]). Therefore, it is enough to verify that our Green function is the same as the one constructed in [23]. To this
$4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - No 3
end, we first recall that by [23, Theorem 1.1] $\widetilde{G}$ is unique among all those real valued, nonnegative functions defined on $\Omega \times \Omega \backslash\{(X, Y) \in \Omega \times \Omega: X=Y\}$, such that for each $X \in \Omega$ and $r>0$,

$$
\begin{gather*}
\widetilde{G}(X, \cdot) \in W^{1,2}(\Omega \backslash B(X, r)) \cap W_{0}^{1,1}(\Omega)  \tag{3.17}\\
\iint_{\Omega} \nabla_{Y} \widetilde{G}(X, Y) \cdot \nabla \phi(Y) d Y=\phi(X), \quad \forall \phi \in C_{0}^{\infty}(\Omega) . \tag{3.18}
\end{gather*}
$$

It is clear that (3.18) holds for our Green function $G(X, Y)$, by Definition 3.4. Thus, we need only show that $G$ satisfies (3.17). As in [31, p. 5], for $X \in \Omega$ fixed, we may construct $v(X, \cdot)$, the variational solution to the Dirichlet problem with data $\mathcal{E}(X-\cdot)$. In particular, $v(X, \cdot) \in W^{1,2}(\Omega)$. Since $\mathcal{E}(X-\cdot)$ is Lipschitz on $\partial \Omega$, and since every point on $\partial \Omega$ is Wiener regular, it follows as in [31, p. 5] that $v(X, \cdot) \in C(\bar{\Omega})$, and therefore

$$
\begin{equation*}
v(X, Y):=\int_{\partial \Omega} \mathcal{E}(X-z) d \omega^{Y}(z) \tag{3.19}
\end{equation*}
$$

(see, e.g., [22], p. 25). Thus, $G(X, Y)=\mathscr{E}(X-Y)-v(X, Y)$ (cf. (3.4)), and since $v \in W^{1,2}(\Omega)$, we obtain (3.17).

Finally we verify (3.16). We begin by reducing matters to the case that $\Omega$ is bounded. Indeed, for the left hand side of (3.16), we may pass immediately from the bounded to the unbounded case by splitting $\Phi$ into positive and negative parts, and using (3.2). To pass to the limit on the right hand side is more delicate, and we proceed as follows. As above, given an unbounded domain $\Omega$, let $G_{R}$ denote the Green function for the domain $\Omega_{R}:=\Omega \cap B\left(x_{0}, 2 R\right)$, for some fixed $x_{0} \in \partial \Omega$. We claim that

$$
\begin{equation*}
\lim _{R \rightarrow \infty} \iint_{\Omega_{R}} \nabla_{Y} G_{R}(Y, X) \cdot \mathbf{h}(Y) d Y=\iint_{\Omega} \nabla_{Y} G(Y, X) \cdot \mathbf{h}(Y) d Y \tag{3.20}
\end{equation*}
$$

for all Lipschitz vector-valued $\mathbf{h}$ with compact support in $\mathbb{R}^{n+1}$. Given the claim, and assuming that (3.16) holds for bounded $\Omega$, we may then pass to the unbounded case by setting $\mathbf{h}=\nabla \Phi$.

Thus, to reduce the proof of (3.16) to the case that $\Omega$ is bounded, it remains to prove (3.20). To this end, we first recall our previous observation that for bounded domains with Wiener regular boundaries, our Green function is the same as that constructed in [23]. Thus, there is a purely dimensional constant $C_{n}$ such that for every $R<\infty$, and $X \in \Omega_{R}, \nabla G_{R}(\cdot, X)$ enjoys the weak- $L^{(n+1) / n}$ estimate

$$
\left|\left\{Y \in \Omega_{R}:\left|\nabla_{Y} G_{R}(Y, X)\right|>\lambda\right\}\right| \leq C_{n} \lambda^{-(n+1) / n}
$$

Consequently, if $A \subset \Omega_{R}$, we have that

$$
\begin{equation*}
\iint_{A}\left|\nabla_{Y} G_{R}(Y, X)\right|^{p} d Y \leq C(n, p,|A|), \quad \forall p<(n+1) / n \tag{3.21}
\end{equation*}
$$

as may be deduced from the weak-type inequality by arguing as in the proof of Kolmogorov's lemma. We emphasize that the constant in the last inequality depends only upon $n, p$ and $|A|$, but not on $R$. Let us now fix a ball $B_{0}:=B\left(X_{0}, R_{0}\right) \subset \mathbb{R}^{n+1}$, and consider a Lipschitz function $\mathbf{h}$ supported in $B_{0}$. We note that
(3.22)

$$
\iint_{\Omega} \nabla_{Y} G_{R}(Y, X) \cdot \mathbf{h}(Y) d Y=\iint_{\Omega} G_{R}(Y, X) \operatorname{div} \mathbf{h}(Y) d Y \rightarrow \iint_{\Omega} G(Y, X) \operatorname{div} \mathbf{h}(Y) d Y
$$

as $R \rightarrow \infty$, where we have used first that $G_{R} \in W_{0}^{1,1}\left(\Omega_{R}\right)$ (again, because $G_{R}$ coincides with the [23] Green function), and then (3.5) (in $\Omega \cap B_{0} \cap\{\delta(Y)>\varepsilon\}$ ), along with (3.12) (to control small errors in the "border strip" $\Omega \cap B_{0} \cap\{\delta(Y) \leq \varepsilon\}$ ). Here we may suppose that $R_{0} \ll R$ so that $B_{0} \cap \Omega \subset \Omega_{R}$. Let us now extend $G(\cdot, X)$ to be zero in $\mathbb{R}^{n+1} \backslash \bar{\Omega}$, and call this extension $\mathscr{G}$. Then from (3.21) and (3.22) it follows that

$$
\begin{equation*}
\left|\iint_{\mathbb{R}^{n+1}} \mathscr{G}(Y, X) \operatorname{div} \mathbf{h}(Y) d Y\right| \leq C\left(n, p,\left|B_{0}\right|\right)\|h\|_{p^{\prime}}, \quad 1<p<(n+1) / n . \tag{3.23}
\end{equation*}
$$

Taking a supremum over all Lipschitz $\mathbf{h}$ supported in $B_{0}$, with $\|\mathbf{h}\|_{p^{\prime}}=1$, we obtain that for $p \in(1,(n+1) / n)$,

$$
\begin{equation*}
\nabla \mathscr{G}(\cdot, X) \in L^{p}\left(B_{0}\right), \quad \text { with }\|\nabla \mathscr{G}\|_{L^{p}\left(B_{0} \backslash\{X\}\right)} \leq C\left(n, p,\left|B_{0}\right|\right) . \tag{3.24}
\end{equation*}
$$

Now let $\psi \in C^{\infty}(\mathbb{R})$, with $0 \leq \psi \leq 1, \psi(t) \equiv 0$ if $t \leq 1, \psi(t) \equiv 1$ if $t \geq 2$. We fix $\mathbf{h}$ as above, let $\varepsilon>0$, and set $\mathbf{h}_{\varepsilon}(Y):=\mathbf{h}(Y) \psi(\delta(Y) / \varepsilon)$. Then, by (3.5),

$$
\begin{align*}
\iint_{\Omega} \nabla_{Y} G_{R}(Y, X) & \cdot \mathbf{h}_{\varepsilon}(Y) d Y=\iint_{\Omega} G_{R}(Y, X) \operatorname{div} \mathbf{h}_{\varepsilon}(Y) d Y  \tag{3.25}\\
& \rightarrow \iint_{\Omega} G(Y, X) \operatorname{div} \mathbf{h}_{\varepsilon}(Y) d Y=\iint_{\Omega} \nabla_{Y} G(Y, X) \cdot \mathbf{h}_{\varepsilon}(Y) d Y
\end{align*}
$$

as $R \rightarrow \infty$. Also, by (3.21), (3.24) and Hölder's inequality, for $1<p<(n+1) / n$, we have

$$
\begin{align*}
\left|\iint_{\Omega} \nabla_{Y} G_{R}(Y, X) \cdot\left(\mathbf{h}-\mathbf{h}_{\varepsilon}\right)(Y) d Y\right|+\left|\iint_{\Omega} \nabla_{Y} G(Y, X) \cdot\left(\mathbf{h}-\mathbf{h}_{\varepsilon}\right)(Y) d Y\right|  \tag{3.26}\\
\leq C\left(n, p,\left|B_{0}\right|\right)\|\mathbf{h}\|_{\infty}\left|\left\{Y \in \Omega \cap B_{0}: \delta(Y)<2 \varepsilon\right\}\right|^{1 / p^{\prime}} \rightarrow 0
\end{align*}
$$

as $\varepsilon \rightarrow 0$, uniformly in $R$. Then by (3.25)-(3.26), we have that

$$
\text { 27) } \begin{align*}
& \lim _{R \rightarrow \infty} \iint_{\Omega} \nabla_{Y} G_{R}(Y, X) \cdot \mathbf{h}(Y) d Y  \tag{3.27}\\
= & \lim _{R \rightarrow \infty} \iint_{\Omega} \nabla_{Y} G_{R}(Y, X) \cdot \mathbf{h}_{\varepsilon}(Y) d Y+o(1)=\iint_{\Omega} \nabla_{Y} G(Y, X) \cdot \mathbf{h}(Y) d Y+o(1),
\end{align*}
$$

as $\varepsilon \rightarrow 0$. Letting $\varepsilon \rightarrow 0$, we obtain (3.20).
We may now assume that $\Omega$ is bounded, and proceed to prove (3.16) in that case. As above, Wiener regularity then guarantees that a given Perron solution, with Lipschitz data, coincides with the corresponding variational solution with the same data. This is true for the function $v(X, Y)$ defined in (3.19), as well as for

$$
u(X):=\int_{\partial \Omega} \Phi d \omega^{X} .
$$

Thus, in particular, $u-\Phi \in W_{0}^{1,2}(\Omega)$, and we claim that

$$
\begin{equation*}
\iint_{\Omega} \nabla_{Y} G(Y, X) \cdot(\nabla u(Y)-\nabla \Phi(Y)) d Y=u(X)-\Phi(X)=u(X) . \tag{3.28}
\end{equation*}
$$

If $u-\Phi$ were in $C_{0}^{\infty}(\Omega)$, the claim would follow immediately from (3.18). We will pass from $C_{0}^{\infty}(\Omega)$ to $W_{0}^{1,2}(\Omega)$ by a density argument, with a slight complication since the Green function is not in $W^{1,2}$ near the pole. To address this technical issue, we multiply $(u-\Phi)$ by a smooth cut-off function supported in a small neighborhood of the pole $X$. For the part near
the pole we may invoke (3.18): $u$ is harmonic, therefore smooth in $\Omega$ and $u$ times a smooth cut-off is $C_{0}^{\infty}(\Omega)$. For the part away from the pole we use (3.17) and (3.18), plus the routine density argument mentioned above. We leave the details, which are standard, to the reader.

At this point, (3.16) follows immediately from (3.28) and the fact that

$$
\begin{equation*}
\iint_{\Omega} \nabla_{Y} G(Y, X) \nabla u(Y) d Y=0 . \tag{3.29}
\end{equation*}
$$

In turn, we may verify the latter identity as follows. For $0<\varepsilon \ll \delta(X)$, set $\phi_{\varepsilon}(Y):=$ $\phi((X-Y) / \varepsilon)$, where $\phi \in C^{\infty}\left(\mathbb{R}^{n+1}\right), \phi \equiv 1$ in $\mathbb{R}^{n+1} \backslash B(0,2), \phi \equiv 0$ in $B(0,1)$. Then

$$
\begin{array}{r}
\iint_{\Omega} \nabla_{Y} G(Y, X) \cdot \nabla u(Y) d Y=\iint_{\Omega} \nabla_{Y}\left(\mathcal{E}(Y-X) \phi_{\varepsilon}(Y)-v(Y, X)\right) \cdot \nabla u(Y) d Y \\
+\iint_{\Omega} \nabla_{Y}\left(\mathcal{E}(Y-X)\left(1-\phi_{\varepsilon}(Y)\right)\right) \cdot \nabla u(Y) d Y=0+O(\varepsilon)
\end{array}
$$

where in the vanishing term we have used the definition of weak solution, since $\mathcal{E}(\cdot-X) \phi_{\varepsilon}(\cdot)-v(\cdot, X) \in W_{0}^{1,2}(\Omega)$. To obtain the $O(\varepsilon)$ bound, we have used standard estimates for the fundamental solution and its gradient, along with the fact that $\nabla u$ is harmonic and therefore locally bounded in $\Omega$. Finally, we obtain (3.29) by letting $\varepsilon \rightarrow 0^{+}$.

Lemma 3.30. - Let $\Omega$ be a 1-sided NTA domain with n-dimensional ADR boundary, and suppose that every $x \in \partial \Omega$ is regular in the sense of Wiener. Fix $B_{0}:=B\left(x_{0}, r_{0}\right)$ with $x_{0} \in \partial \Omega$, and $\Delta_{0}:=B_{0} \cap \partial \Omega$. Let $B:=B(x, r), x \in \partial \Omega$, and $\Delta:=B \cap \partial \Omega$, and suppose that $2 B \subset B_{0}$. Then for $X \in \Omega \backslash B_{0}$ we have

$$
\begin{equation*}
r^{n-1} G\left(X_{\Delta}, X\right) \leq C \omega^{X}(\Delta) \tag{3.31}
\end{equation*}
$$

If, in addition, $\Omega$ satisfies the qualitative exterior Corkscrew condition, then

$$
\begin{equation*}
\omega^{X}(\Delta) \leq C r^{n-1} G\left(X_{\Delta}, X\right) . \tag{3.32}
\end{equation*}
$$

The constants in (3.31) and (3.32) depend only on dimension and on the constants in the $A D R$ and 1 -sided NTA conditions.

Remark. - Let us emphasize that in several results below we will assume that certain domains satisfy the hypotheses of Lemma 3.30; by this we mean that the domains are 1-sided NTA with $n$-dimensional ADR boundary, which moreover satisfy the qualitative exterior Corkscrew condition (in particular then, every boundary point is regular in the sense of Wiener, cf. Remark 3.10). Notice that in such a case (3.32) holds with $C$ depending only on dimension and on the constants in the $A D R$ and 1 -sided NTA conditions. In particular, $C$ does not depend on the parameter $N$ from the qualitative assumption. This will be crucial when applied to approximating domains.

Proof. - The first estimate (3.31) may be obtained by a well known argument (cf. [10] or [31, Lemma 1.3.3]) using (3.8) plus Harnack's inequality, the upper bound for $G(X, Y)$ in (3.12), and the maximum principle in $\Omega \backslash B\left(X_{\Delta}, \delta\left(X_{\Delta}\right) / 2\right)$ (the use of the maximum principle is justified even in the case that $\Omega$ is unbounded, by virtue of the decay of the Green function at infinity). We omit the details.

The proof of the second estimate (3.32) will require a bit more work. In contrast to the case of previous proofs of this estimate [10, 29], we do not use local Hölder continuity at the boundary for solutions vanishing on a surface ball (since this depends on the parameter $N$ in our qualitative assumption). Instead, we proceed as follows. Fix $B_{0}$ centered on $\partial \Omega$, and $X \in \Omega \backslash B_{0}$, and write $\Upsilon_{B_{0}}=\Upsilon_{B_{0}}(X)$ (cf. (3.9)). As observed in Remark 3.10, $\Upsilon_{B_{0}}$ is a priori finite (possibly depending on $N$ ). Thus, it will suffice to show that $\Upsilon_{B_{0}} \leq C_{\varepsilon}+C \varepsilon \Upsilon_{B_{0}}$, for every small $\varepsilon>0$. Choose now $B=B(x, r)$, with $2 B \subseteq B_{0}$, such that

$$
\frac{1}{2} \Upsilon_{B_{0}} \leq \frac{\omega^{X}(\Delta)}{r^{n-1} G\left(X_{\Delta}, X\right)}
$$

where as usual $\Delta=B \cap \partial \Omega$.
Now set $B:=B(x, r)$ and $\widetilde{B}:=B(x, 5 r / 4)$. Taking $\Phi \in C_{0}^{\infty}(\widetilde{B})$, with $0 \leq \Phi \leq 1$, $\Phi(Y) \equiv 1$ on $B(x, r)$, and $\|\nabla \Phi\|_{\infty} \lesssim 1 / r$, we deduce from (3.16) that

$$
\begin{align*}
\omega^{X}(\Delta) & \lesssim \frac{1}{r} \iint_{\Omega \cap \widetilde{B}}\left|\nabla_{Y} G(Y, X)\right| d Y  \tag{3.33}\\
& =\frac{1}{r} \iint_{\Omega \cap \tilde{B} \cap\{\delta(Y)>\varepsilon r\}}\left|\nabla_{Y} G(Y, X)\right| d Y+\frac{1}{r} \iint_{\Omega \cap \tilde{B} \cap\{\delta(Y) \leq \varepsilon r\}}\left|\nabla_{Y} G(Y, X)\right| d Y \\
& \lesssim \frac{1}{r} \iint_{\Omega \cap \tilde{B} \cap\{\delta(Y)>\varepsilon r\}} \frac{G(Y, X)}{\delta(Y)} d Y+\frac{1}{r} \iint_{\Omega \cap \tilde{B} \cap\{\delta(Y) \leq \varepsilon r\}} \frac{G(Y, X)}{\delta(Y)} d Y \\
& =I+I I,
\end{align*}
$$

where $\varepsilon>0$ is at our disposal, and where in the next to last line we have used standard interior estimates for harmonic functions. By Harnack's inequality and the Harnack chain condition, we have that

$$
I \leq C_{\varepsilon} r^{n-1} G\left(X_{\Delta}, X\right)
$$

as desired. To handle term $I I$, choose $y \in \partial \Omega$ such that $|Y-y|=\delta(Y)$, and set $\Delta(Y):=\Delta(y, \delta(Y))$. Then by (3.31) and the Harnack chain condition we have

$$
\begin{aligned}
I I & \leq \frac{1}{r} \iint_{\Omega \cap \tilde{B} \cap\{\delta(Y) \leq \varepsilon r\}} \frac{\omega^{X}(\Delta(Y))}{(\delta(Y))^{n}} d Y \\
& \lesssim \frac{1}{r} \sum_{k: 2^{-k} \lesssim \varepsilon r} \iint_{\Omega \cap \tilde{B} \cap\left\{2^{-k-1}<\delta(Y) \leq 2^{-k}\right\}} \frac{\omega^{X}(\Delta(Y))}{(\delta(Y))^{n}} d Y \\
& \lesssim \frac{1}{r} \sum_{k: 2^{-k} \lesssim \varepsilon r} \sum_{j} \iint_{B_{j}^{k} \cap\left\{2^{-k-1}<\delta(Y) \leq 2^{-k}\right\}} \frac{\omega^{X}(\Delta(Y))}{(\delta(Y))^{n}} d Y,
\end{aligned}
$$

where in the last step $B_{j}^{k}:=B\left(x_{j}^{k}, 2^{-k+1}\right)$, and for each $k$ in the sum, $\mathscr{B}_{k}:=\left\{B_{j}^{k}\right\}_{j}$ is a collection of balls whose doubles have bounded overlaps, such that $x_{j}^{k} \in \partial \Omega$,

$$
\begin{equation*}
\left(\Omega \cap \widetilde{B} \cap\left\{2^{-k-1}<\delta(Y) \leq 2^{-k}\right\}\right) \subset \bigcup_{j} B_{j}^{k}, \quad \text { and } \quad \bigcup_{j} 2 B_{j}^{k} \subset B(x, 3 r / 2) . \tag{3.34}
\end{equation*}
$$

We leave it to the reader to verify that such a collection exists, by virtue of the ADR property of $\partial \Omega$, for all sufficiently small $\varepsilon$. We then have that

$$
\begin{align*}
I I & \lesssim \frac{1}{r} \sum_{k: 2^{-k} \lesssim \varepsilon r} 2^{-k} \sum_{j} \omega^{X}\left(\Delta\left(x_{j}^{k}, 2^{-k+2}\right)\right) \leq C \varepsilon \omega^{X}(\Delta(x, 3 r / 2))  \tag{3.35}\\
& \leq C \varepsilon \sum_{\Delta^{\prime}} \omega^{X}\left(\Delta^{\prime}\right) \leq C \varepsilon \Upsilon_{B_{0}} r^{n-1} \sum_{\Delta^{\prime}} G\left(X_{\Delta^{\prime}}, X\right) \leq C \varepsilon \Upsilon_{B_{0}} r^{n-1} G\left(X_{\Delta}, X\right),
\end{align*}
$$

where in the second, third, fourth and fifth inequalities we have used, respectively, the bounded overlap property of the balls $2 B_{j}^{k}$; the ADR property to cover $\Delta(x, 3 r / 2)$ by a collection $\left\{\Delta^{\prime}\right\}$ of bounded cardinality, such that $r_{\Delta^{\prime}} \approx r$, and $\Delta^{\prime}=B^{\prime} \cap \partial \Omega$, with $B^{\prime}$ centered on $\partial \Omega$ and $2 B^{\prime} \subset B_{0}$; the definition of $\Upsilon_{B_{0}}$; and the Harnack chain condition. We then obtain (3.32) by choosing $\varepsilon$ sufficiently small.

Remark. - Let us observe that from the previous proof it follows that we are not really using the full strength of the qualitative exterior Corkscrew condition, but only that every boundary point is regular in the sense of Wiener and that $\Upsilon\left(B_{0}\right)$ is a priori finite. Although these relaxed qualitative hypotheses suffice for our purposes, the qualitative exterior Corkscrew condition is cleaner, easier to check in practice and holds for the approximating domains introduced below.

Corollary 3.36. - Suppose that $\Omega$ is a 1 -sided NTA domain with $n$-dimensional ADR boundary and that it also satisfies the qualitative exterior Corkscrew condition. Let $B:=B(x, r), x \in \partial \Omega, \Delta:=B \cap \partial \Omega$ and $X \in \Omega \backslash 4 B$. Then there is a uniform constant $C$ such that

$$
\omega^{X}(2 \Delta) \leq C \omega^{X}(\Delta) .
$$

Proof. - The conclusion of the corollary follows immediately from the combination of (3.31) and (3.32), and Harnack's inequality. We omit the details.

Next, we establish a bound of "Carleson-type" for the Green function. The Carleson estimate is already known for arbitrary non-negative harmonic functions vanishing on a surface ball [1, 2]; however, specializing to the Green function, one may give a fairly simple direct proof, based upon that of the previous lemma.

Lemma 3.37. - Suppose that $\Omega$ is a 1 -sided NTA domain with $n$-dimensional ADR boundary and that it also satisfies the qualitative exterior Corkscrew condition. Then there is a uniform constant $C$ such that for each $B:=B(x, r), x \in \partial \Omega, \Delta=B \cap \partial \Omega$, and $X \in \Omega \backslash 2 B$, we have

$$
\begin{equation*}
\sup _{Y \in B \cap \Omega} G(Y, X) \leq C G\left(X_{\Delta}, X\right) . \tag{3.38}
\end{equation*}
$$

Proof. - Fix $B, \Delta$ and $X$ as in the statement of the lemma, and set $u(Y):=G(Y, X)$. Extending $u$ to be zero in $\Omega_{\text {ext }}$, we obtain from (3.16) that $u$ is subharmonic in $B(x, 3 r / 2)$.

Let $B^{\prime}:=B(x, 5 r / 4)$. By the sub-mean value inequality, we have that

$$
\begin{aligned}
\sup _{Y \in B} u(Y) & \lesssim \frac{1}{\left|B^{\prime}\right|} \iint_{B^{\prime}} u=\frac{1}{\left|B^{\prime}\right|} \iint_{B^{\prime} \cap \Omega} u \\
& =\frac{1}{\left|B^{\prime}\right|} \iint_{B^{\prime} \cap \Omega \cap\{\delta(Y)>\varepsilon r\}} u(Y) d Y+\frac{1}{\left|B^{\prime}\right|} \iint_{B^{\prime} \cap \Omega \cap\{\delta(Y) \leq \varepsilon r\}} u(Y) d Y \\
& =: I^{*}+I I^{*} \leq C_{\varepsilon} u\left(X_{\Delta}\right)+I I^{*}
\end{aligned}
$$

where in the last step we have used the Harnack chain condition to estimate term $I^{*}$, and we have fixed a small $\varepsilon$ as in the proof of Lemma 3.30 so that (3.34) holds. Moreover, by definition of $u$,

$$
I I^{*} \lesssim \frac{1}{\left|B^{\prime}\right|} \varepsilon r \iint_{B^{\prime} \cap \Omega \cap\{\delta(Y) \leq \varepsilon r\}} \frac{G(Y, X)}{\delta(Y)} d Y \approx \varepsilon r^{1-n} I I
$$

where $I I$ is exactly the same term as in (3.33). In turn, by (3.35) and the fact that $\Upsilon_{B_{0}}(X)$ is uniformly bounded (the latter fact is simply a restatement of (3.32)), we find that

$$
I I^{*} \leq C \varepsilon^{2} G\left(X_{\Delta}, X\right)
$$

Under the same hypotheses as in the previous two lemmata, we shall obtain a comparison principle for the Green function, again without the use of Hölder continuity at the boundary. In order to state our comparison principle, we shall need to introduce the notion of a Carleson region. Given a "dyadic cube" $Q \in \mathbb{D}(\partial \Omega)$, the discretized Carleson region $\mathbb{D}_{Q}$ is defined to be

$$
\begin{equation*}
\mathbb{D}_{Q}:=\left\{Q^{\prime} \in \mathbb{D}: Q^{\prime} \subseteq Q\right\} \tag{3.39}
\end{equation*}
$$

For future reference, we also introduce discretized sawtooth regions as follows. Given a family $\mathcal{F}$ of disjoint cubes $\left\{Q_{j}\right\} \subset \mathbb{D}$, we define the global discretized sawtooth relative to $\mathcal{F}$ by

$$
\begin{equation*}
\mathbb{D}_{\mathcal{F}}:=\mathbb{D} \backslash \bigcup_{\mathcal{F}} \mathbb{D}_{Q_{j}} \tag{3.40}
\end{equation*}
$$

i.e., $\mathbb{D}_{\mathcal{G}}$ is the collection of all $Q \in \mathbb{D}$ that are not contained in any $Q_{j} \in \mathcal{F}$. Given some fixed cube $Q$, the local discretized sawtooth relative to $\mathcal{F}$ by

$$
\begin{equation*}
\mathbb{D}_{\mathcal{F}, Q}:=\mathbb{D}_{Q} \backslash \bigcup_{\mathcal{F}} \mathbb{D}_{Q_{j}}=\mathbb{D}_{\mathcal{F}} \cap \mathbb{D}_{Q} \tag{3.41}
\end{equation*}
$$

We shall also require "geometric" Carleson regions and sawtooths. Let us first recall that we write $k=k(Q)$ if $Q \in \mathbb{D}_{k}$ (cf. Lemma ??), and in that case the "length" of $Q$ is denoted by $\ell(Q)=2^{-k(Q)}$. We also recall that there is a Corkscrew point $X_{Q}$, relative to each $Q \in \mathbb{D}$ (in fact, there are many such, but we just pick one). Given such a $Q$, we define an associated "Whitney region" as follows. Let $\mathcal{W}=\mathcal{W}(\Omega)$ denote a collection of (closed) dyadic Whitney cubes of $\Omega$, so that the cubes in $W$ form a pairwise non-overlapping covering of $\Omega$, which satisfy

$$
\begin{equation*}
4 \operatorname{diam}(I) \leq \operatorname{dist}(4 I, \partial \Omega) \leq \operatorname{dist}(I, \partial \Omega) \leq 40 \operatorname{diam}(I), \quad \forall I \in \mathcal{W} \tag{3.42}
\end{equation*}
$$

$4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - No 3
(just dyadically divide the standard Whitney cubes, as constructed in [40, Chapter VI], into cubes with side length $1 / 8$ as large) and also

$$
(1 / 4) \operatorname{diam}\left(I_{1}\right) \leq \operatorname{diam}\left(I_{2}\right) \leq 4 \operatorname{diam}\left(I_{1}\right),
$$

whenever $I_{1}$ and $I_{2}$ touch. Let $\ell(I)$ denote the side length of $I$, and write $k=k_{I}$ if $\ell(I)=2^{-k}$. We set

$$
\begin{equation*}
\mathcal{W}_{Q}:=\left\{I \in \mathcal{W}: k(Q)-m_{0} \leq k_{I} \leq k(Q)+1, \text { and dist }(I, Q) \leq C_{0} 2^{-k(Q)}\right\} \tag{3.43}
\end{equation*}
$$

where we may (and do) choose the constant $C_{0}$ and positive integer $m_{0}$, depending only on the constants in the Corkscrew condition and in the dyadic cube construction (cf. Lemma ??), so that $X_{Q} \in I$ for some $I \in \mathcal{W}_{Q}$, and for each dyadic child $Q^{j}$ of $Q$, the respective Corkscrew points $X_{Q^{j}} \in I^{j}$ for some $I^{j} \in \mathcal{W}_{Q}$. In particular, the collection $\mathcal{W}_{Q}$ is nonempty for every $Q \in \mathbb{D}$. Moreover as long as $C_{0}$ is chosen large enough depending on the constant $c$ in the Corkscrew condition, then by the properties of Whitney cubes, we may always find an $I \in \mathcal{W}_{Q}$ with the slightly more precise property that $k(Q)-1 \leq k_{I} \leq k(Q)$. We may further suppose, by choosing $C_{0}$ large enough, that

$$
\begin{equation*}
\mathcal{W}_{Q_{1}} \cap \mathcal{W}_{Q_{2}} \neq \varnothing, \text { whenever } 1 \leq \frac{\ell\left(Q_{2}\right)}{\ell\left(Q_{1}\right)} \leq 2, \text { and } \operatorname{dist}\left(Q_{1}, Q_{2}\right) \leq 1000 \ell\left(Q_{2}\right) \tag{3.44}
\end{equation*}
$$

We omit the details. In the sequel, we shall assume always that $C_{0}$ has been so chosen, and further that $C_{0} \geq 1000 \sqrt{n}$.

We shall need to augment $\mathcal{W}_{Q}$ in order to exploit the Harnack chain condition. It will be convenient to introduce the following notation: given a subset $A \subset \Omega$, we write

$$
X \rightarrow_{A} Y
$$

if the interior of $A$ contains all the balls in a Harnack chain (in $\Omega$ ), connecting $X$ to $Y$, and if, moreover, for any point $Z$ contained in any ball in the Harnack chain, we have

$$
\begin{equation*}
\operatorname{dist}(Z, \partial \Omega) \approx \operatorname{dist}(Z, \Omega \backslash A) \tag{3.45}
\end{equation*}
$$

with uniform control of the implicit constants. We denote by $X(I)$ the center of a cube $I \in \mathbb{R}^{n+1}$, and we recall that $X_{Q}$ denotes a designated Corkscrew point relative to $Q$, which we may, from this point on, assume without loss of generality to be the center of some Whitney cube $I$ such that $\ell(I) \approx \ell(Q) \approx \operatorname{dist}(I, Q)$. More precisely, we note the following.

Remark 3.46. - Having fixed the collection $W$ (the Whitney cubes of $\Omega$ ), by taking the Corkscrew constant $c$ to be slightly smaller, if necessary, we may assume that the Corkscrew point $X_{Q}$ is the center of some $I \in \mathcal{W}$, with $I \subset B_{Q} \cap \Omega$ and $\ell(I) \approx \ell(Q)$.

We now define the augmented collection $\mathcal{W}_{Q}^{*}$ as follows. For each $I \in \mathcal{W}_{Q}$, we form a Harnack chain, call it $H(I)$, from the center $X(I)$ to the Corkscrew point $X_{Q}$. We now denote by $\mathcal{W}(I)$ the collection of all Whitney cubes which meet at least one ball in the chain $H(I)$, and we set

$$
\mathcal{W}_{Q}^{*}:=\bigcup_{I \in \mathcal{W}_{Q}} W(I)
$$

We also define, for $\lambda \in(0,1)$ to be chosen momentarily,

$$
\begin{equation*}
U_{Q}:=\bigcup_{W_{Q}^{*}}(1+\lambda) I=: \bigcup_{I \in W_{Q}^{*}} I^{*} \tag{3.47}
\end{equation*}
$$

By construction, we then have that

$$
\begin{equation*}
\mathcal{W}_{Q} \subset W_{Q}^{*} \subset \mathcal{W} \quad \text { and } \quad X_{Q} \in U_{Q}, X_{Q^{j}} \in U_{Q} \tag{3.48}
\end{equation*}
$$

for each child $Q^{j}$ of $Q$. It is also clear that there are uniform constants $k^{*}$ and $K_{0}$ such that

$$
\begin{align*}
k(Q)-k^{*} & \leq k_{I} \leq k(Q)+k^{*}, \quad \forall I \in W_{Q}^{*} \\
X(I) & \rightarrow_{U_{Q}} X_{Q}, \quad \forall I \in W_{Q}^{*}  \tag{3.49}\\
\operatorname{dist}(I, Q) & \leq K_{0} 2^{-k(Q)}, \quad \forall I \in W_{Q}^{*},
\end{align*}
$$

where $k^{*}, K_{0}$ and the implicit constants in (3.45) (which pertain to the condition $X(I) \rightarrow U_{Q} X_{Q}$ ), depend only on the "allowable parameters" (since $m_{0}$ and $C_{0}$ also have such dependence) and on $\lambda$. Thus, by the addition of a few nearby Whitney cubes of diameter also comparable to that of $Q$, we can "augment" $W_{Q}$ so that the Harnack chain condition holds in $U_{Q}$.

We fix the parameter $\lambda$ so that for any $I, J \in \mathcal{W}$,

$$
\begin{align*}
\operatorname{dist}\left(I^{*}, J^{*}\right) & \approx \operatorname{dist}(I, J) \\
\operatorname{int}\left(I^{*}\right) \cap \operatorname{int}\left(J^{*}\right) & \neq \varnothing \Longleftrightarrow \partial I \cap \partial J \neq \varnothing \tag{3.50}
\end{align*}
$$

(the fattening thus ensures overlap of $I^{*}$ and $J^{*}$ for any pair $I, J \in \mathcal{W}$ whose boundaries touch, so that the Harnack chain property then holds locally, with constants depending upon $\lambda$, in $\left.I^{*} \cup J^{*}\right)$. By choosing $\lambda$ sufficiently small, we may also suppose that there is a $\tau \in(1 / 2,1)$ such that for distinct $I, J \in \mathcal{W}$,

$$
\begin{equation*}
\tau J \cap I^{*}=\varnothing \tag{3.51}
\end{equation*}
$$

We remark that any sufficiently small choice of $\lambda$ (say $0<\lambda \leq \lambda_{0}$ ) will do for our purposes.
Of course, there may be some flexibility in the choice of additional Whitney cubes which we add to form the augmented collection $W_{Q}^{*}$, but having made such a choice for each $Q \in \mathbb{D}$, we fix it for all time. We may then define the Carleson box associated to $Q$ by

$$
\begin{equation*}
T_{Q}:=\operatorname{int}\left(\bigcup_{Q^{\prime} \in \mathbb{D}_{Q}} U_{Q^{\prime}}\right) \tag{3.52}
\end{equation*}
$$

Similarly, we may define geometric sawtooth regions as follows. As above, give a family $\mathcal{F}$ of disjoint cubes $\left\{Q_{j}\right\} \subset \mathbb{D}$, we define the global sawtooth relative to $\mathscr{G}$ by

$$
\begin{equation*}
\Omega_{\mathscr{F}}:=\operatorname{int}\left(\bigcup_{Q^{\prime} \in \mathbb{D}_{\mathscr{G}}} U_{Q^{\prime}}\right) \tag{3.53}
\end{equation*}
$$

and again given some fixed $Q \in \mathbb{D}$, the local sawtooth relative to $\mathscr{F}$ by

$$
\begin{equation*}
\Omega_{\mathcal{G}, Q}:=\operatorname{int}\left(\bigcup_{Q^{\prime} \in \mathbb{D}_{\mathscr{G}, Q}} U_{Q^{\prime}}\right) \tag{3.54}
\end{equation*}
$$

For future reference, we present the following.
$4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - No 3

Lemma 3.55. - Suppose that $\Omega$ is a 1 -sided NTA domain with an $A D R$ boundary. Given $Q \in \mathbb{D}$, let $B_{Q}:=B\left(x_{Q}, r\right), r \approx \ell(Q)$, and $\Delta_{Q}:=B_{Q} \cap \partial \Omega \subset Q$, be as in (1.16) and (1.17). Then for each $Q$, there is a ball $B_{Q}^{\prime}:=B\left(x_{Q}, s\right) \subset B_{Q}$, with $s \approx \ell(Q) \approx r$, such that

$$
\begin{equation*}
B_{Q}^{\prime} \cap \Omega \subset T_{Q} \tag{3.56}
\end{equation*}
$$

Moreover, for a somewhat smaller choice of s $\approx\left(K_{0}\right)^{-1} \ell(Q)$, we have for every pairwise disjoint family $\mathcal{F} \subset \mathbb{D}$, and for each $Q_{0} \in \mathbb{D}$ containing $Q$, that

$$
\begin{equation*}
B_{Q}^{\prime} \cap \Omega_{\mathscr{G}, Q_{0}}=B_{Q}^{\prime} \cap \Omega_{\mathscr{G}, Q} . \tag{3.57}
\end{equation*}
$$

Proof. - We prove (3.56) first. Let $Y \in \Omega$, with $\left|Y-x_{Q}\right|<c r=: s$, where $c>0$ is to be determined. Then $Y \in I \in \mathcal{W}$, with

$$
\ell(I) \approx \delta(Y) \leq\left|Y-x_{Q}\right|<c r .
$$

Fix $Q_{I} \in \mathbb{D}$ such that $\ell\left(Q_{I}\right)=\ell(I)$, and $\operatorname{dist}\left(Q_{I}, I\right) \approx \ell(I)$. In particular, $I \in \mathcal{W}_{Q_{I}} \subset \mathcal{W}_{Q_{I}}^{*}$, so that $I \subset \operatorname{int}\left(I^{*}\right) \subset \operatorname{int}\left(U_{Q_{I}}\right)$. By the triangle inequality, for all $x \in Q_{I}$, we have

$$
\left|x-x_{Q}\right| \leq|x-Y|+\left|Y-x_{Q}\right| \leq C \ell(I)+c r \leq C c r<r,
$$

if $c$ is chosen small enough. Hence, $Q_{I} \subset \Delta_{Q} \subset Q$, so $Y \in \bigcup_{Q^{\prime} \in \mathbb{D}_{Q}} \operatorname{int}\left(U_{Q^{\prime}}\right) \subset T_{Q}$.
We now turn to the proof of (3.57). Since $Q \subset Q_{0}$, the "right to left" containment is trivial, for any choice of $B_{Q}^{\prime}$. We therefore suppose that $Y \in B_{Q}^{\prime} \cap \Omega_{\mathcal{F}, Q_{0}}$, where again $B_{Q}^{\prime}:=B\left(x_{Q}, s\right)$, and $s$ will be chosen momentarily. It is enough to show that $Y \in \Omega_{\mathcal{G}, Q}$, for some choice of $s \approx\left(K_{0}\right)^{-1} \ell(Q)$. Since $Y \in \Omega_{\mathcal{G}, Q_{0}}$, by definition there is some $Q^{\prime} \in \mathbb{D}_{Q_{0}} \cap \mathbb{D}_{\mathcal{G}}$, for which $Y \in I^{*}=(1+\lambda) I$, with $I \in \mathcal{W}_{Q^{\prime}}^{*}$. Then

$$
\ell\left(Q^{\prime}\right) \approx \ell(I) \approx \delta(Y) \leq\left|Y-x_{Q}\right| \leq s
$$

where in the last step we have used that $Y \in B_{Q}^{\prime}$. Moreover, for every $y^{\prime} \in Q^{\prime}$, we have

$$
\left|y^{\prime}-Y\right| \lesssim \operatorname{dist}\left(Y, Q^{\prime}\right) \lesssim \operatorname{dist}\left(I, Q^{\prime}\right) \lesssim K_{0} \ell\left(Q^{\prime}\right) \lesssim K_{0} s
$$

Thus, by the triangle inequality, for every $y^{\prime} \in Q^{\prime}$, we have

$$
\left|y^{\prime}-x_{Q}\right| \lesssim K_{0} s<r \approx \ell(Q)
$$

by choice of $s=c\left(K_{0}\right)^{-1} \ell(Q)$ with $c$ sufficiently small. Thus, $Q^{\prime} \subset \Delta_{Q} \subset Q$, whence $Y \in \bigcup_{Q^{\prime} \in \mathbb{D}_{g, Q}} U_{Q^{\prime}}$, i.e., we have shown that $B_{Q}^{\prime} \cap \Omega_{\mathcal{G}, Q_{0}} \subset \bigcup_{Q^{\prime} \in \mathbb{D}_{g, Q}} U_{Q^{\prime}}$, and therefore $\operatorname{int}\left(B_{Q}^{\prime}\right) \cap \Omega_{\mathscr{G}, Q_{0}} \subset \Omega_{\mathscr{F}, Q}$, by definition. Choosing $s$ slightly smaller, which amounts to replacing $B_{Q}^{\prime}$ by a slightly smaller ball, we obtain (3.57).

We also define as follows the "Carleson box" $T_{\Delta}$ associated to a surface ball $\Delta:=\Delta\left(x_{\Delta}, r\right)$. Let $k(\Delta)$ denote the unique $k \in \mathbb{Z}$ such that $2^{-k-1}<200 r \leq 2^{-k}$, and set

$$
\begin{equation*}
\mathbb{D}^{\Delta}:=\left\{Q \in \mathbb{D}_{k(\Delta)}: Q \cap 2 \Delta \neq \varnothing\right\} \tag{3.58}
\end{equation*}
$$

We then define

$$
\begin{equation*}
T_{\Delta}:=\operatorname{int}\left(\bigcup_{Q \in \mathbb{D}^{\Delta}} \overline{T_{Q}}\right) . \tag{3.59}
\end{equation*}
$$

For future reference, we record the following analogue of Lemma 3.55. Set $B_{\Delta}:=B\left(x_{\Delta}, r\right)$, so that $\Delta=B_{\Delta} \cap \partial \Omega$. Then

$$
\begin{equation*}
\frac{5}{4} B_{\Delta} \cap \Omega \subset T_{\Delta} \tag{3.60}
\end{equation*}
$$

Indeed, let $X \in \Omega$ with $\left|X-x_{\Delta}\right|<5 r / 4$. Then $X \in I \in \mathcal{W}$ with $\ell(I) \approx \delta(X)<5 r / 4$, so that $\ell(I) \leq \ell(Q)$, for each $Q \in \mathbb{D}^{\Delta}$.

Suppose first that $\delta(X)<3 r / 4$. There is an $x_{1} \in \partial \Omega$ such that $\left|X-x_{1}\right|=\delta(X)$, so that by the triangle inequality, $\left|x_{1}-x_{\Delta}\right|<2 r$. Consequently, there is a $Q \in \mathbb{D}^{\Delta}$ for which $x_{1} \in \bar{Q}$, whence there is a $Q^{\prime} \subset Q$, whose closure contains $x_{1}$, such that $\ell\left(Q^{\prime}\right)=\ell(I)$, and $\operatorname{dist}\left(Q^{\prime}, I\right) \leq \delta(X) \leq 41 \operatorname{diam}(I) \ll C_{0} \ell\left(Q^{\prime}\right)$ (cf (3.42)-(3.43)). Thus, $I \in \mathcal{W}_{Q^{\prime}}$, so $X \in \operatorname{int}\left(I^{*}\right) \subset \operatorname{int}\left(U_{Q^{\prime}}\right) \subset T_{Q} \subset T_{\Delta}$.

Now suppose that $3 r / 4 \leq \delta(X)<5 r / 4$. Then $X \in I$ with $\ell(I) \approx r$, and $\operatorname{dist}\left(I, Q^{\prime}\right) \approx r$ for every $Q^{\prime}$ contained in any $Q \in \mathbb{D}^{\Delta}$, with $\ell\left(Q^{\prime}\right) \approx \ell(I)$. In that case, we have that $I \in \mathcal{W}_{Q^{\prime}}$, for each such $Q^{\prime}$, so that $X \in T_{Q}, \forall Q \in \mathbb{D}_{\Delta}$.

Lemma 3.61. - Suppose that $\Omega$ is a 1 -sided NTA domain with an ADR boundary. Then all of its Carleson boxes $T_{Q}$ and $T_{\Delta}$, and sawtooth regions $\Omega_{\mathcal{G}}$, and $\Omega_{\mathcal{F}, Q}$ are also 1-sided NTA domains with $A D R$ boundaries. If in addition $\partial \Omega$ is also $U R$, then so is the boundary of each Carleson box $T_{Q}$ and $T_{\Delta}$. In all cases, the implicit constants are uniform, and depend only on dimension and on the corresponding constants for $\Omega$.

We defer the proof until Appendix A.
We remark that it seems likely that one could show that the sawtooth regions also inherit the UR property, but in our case, the only sawtooths that we work with will enjoy an even stronger property, so we shall not bother to explore this issue here.

Lemma 3.62. - Suppose that $\Omega$ is a 1-sided NTA domain with an ADR boundary and that $\Omega$ also satisfies the qualitative exterior Corkscrew condition. Then all of its Carleson boxes $T_{Q}$ and $T_{\Delta}$, and sawtooth regions $\Omega_{\mathcal{F}}$, and $\Omega_{\mathcal{F}, Q}$ satisfy the qualitative exterior Corkscrew condition. In all cases, the implicit constants are uniform, and depend only on dimension and on the corresponding constants for $\Omega$.

The proof of this result is almost trivial. Consider for instance the domain $\Omega_{\mathcal{Y}, Q}$, and let $x \in \partial \Omega_{\mathcal{G}, Q}$ and $r \leq 2^{-N}$, with $N$ corresponding to the qualitative exterior Corkscrew condition assumed on $\Omega$. If either $x \in \partial \Omega$ or $x \in \Omega$ with $\delta(x)<r / 2$, there exists $y \in \partial \Omega$ such that $B(y, r / 2) \subset B(x, r)$. Then the exterior Corkscrew point relative to $\Delta(y, r / 2)$ is also a Corkscrew point relative to $B(x, r) \cap \Omega_{\mathscr{G}, Q}$. The case $x \in \Omega$ with $\delta(x) \geq r / 2$ is as follows. There exists a Whitney box $I$, with $\ell(I) \approx \delta(X)$, such that $x \in \partial I^{*}$ and $\operatorname{int}\left(I^{*}\right) \subset \Omega_{\mathcal{F}, Q}$. Note that $\partial I^{*}$ can be covered by Whitney boxes $J$ that meet $I$ by (3.50). Since $x$ is a boundary point of $\Omega_{\mathcal{F}, Q}$, there is a $J \ni x$, with $J \notin \mathcal{W}_{Q^{\prime}}^{*}$ for any $Q^{\prime} \in \mathbb{D}_{\mathcal{G}, Q}$. Consequently, $B(x, r)$ has an "ample" intersection with $J \backslash \Omega_{\mathcal{G}, Q}$, wherein we may find the required Corkscrew point. Further details are left to the reader.

We are now ready to state our comparison principle for the Green function. The result is already known [1], but we include the proof here for the sake of self-containment. Given a
surface ball $\Delta:=\Delta(x, r)$, let $B_{\Delta}:=B(x, r)$, so that $\Delta=B_{\Delta} \cap \partial \Omega$. We fix $\kappa_{0}$ large enough that

$$
\begin{equation*}
\overline{T_{\Delta}} \subset \kappa_{0} B_{\Delta} \cap \bar{\Omega} \tag{3.63}
\end{equation*}
$$

Lemma 3.64. - Suppose that $\Omega$ is a 1 -sided NTA domain with n-dimensional ADR boundary and that it also satisfies the qualitative exterior Corkscrew condition. Then there is a uniform constant $C$ such that for each surface ball $\Delta$, and for every $X, Y \in \Omega \backslash 2 \kappa_{0} B_{\Delta}$, and $Z \in B_{\Delta} \cap \Omega$, we have

$$
\frac{1}{C} \frac{G(Z, X)}{G(Z, Y)} \leq \frac{G\left(X_{\Delta}, X\right)}{G\left(X_{\Delta}, Y\right)} \leq C \frac{G(Z, X)}{G(Z, Y)}
$$

Remark 3.65. - By Lemma 3.61 and Lemma 3.62, every Carleson box $T_{\Delta} \subset \Omega$ is a 1 -sided NTA domain with $n$-dimensional ADR boundary and also satisfies the qualitative exterior Corkscrew condition.

Proof. - We follow [31, Lemma 1.3.7]. Given a surface ball $\Delta$, fix $X, Y \in \Omega \backslash 2 \kappa_{0} B_{\Delta}$, and let $\omega_{\Delta}^{Z}$ denote harmonic measure for the sub-domain $T_{\Delta}$. Set

$$
S_{1}:=\partial T_{\Delta} \cap\{Z \in \Omega: \delta(Z)>r / 2\}
$$

where $r$ is the radius of $B_{\Delta}$. By Remark 3.65, Corollary 3.36 applies in $T_{\Delta}$, whence by (3.60),

$$
\begin{equation*}
\omega_{\Delta}^{Z}\left(\partial T_{\Delta} \cap \Omega\right) \leq C \omega_{\Delta}^{Z}\left(S_{1}\right), \quad \forall Z \in B_{\Delta} \cap \Omega \tag{3.66}
\end{equation*}
$$

Now set $B^{*}:=\kappa_{0} B_{\Delta}$ and $\Delta^{*}:=B^{*} \cap \partial \Omega$. By Lemma 3.37 and the Harnack chain condition we have

$$
G(Z, X) \leq C G\left(X_{\Delta^{*}}, X\right) \lesssim G\left(X_{\Delta}, X\right), \quad \forall Z \in B^{*} \cap \Omega
$$

By the maximum principle and (3.63), we then have

$$
\begin{equation*}
G(Z, X) \leq C G\left(X_{\Delta}, X\right) \omega_{\Delta}^{Z}\left(\partial T_{\Delta} \cap \Omega\right), \quad \forall Z \in T_{\Delta} \tag{3.67}
\end{equation*}
$$

On the other hand, by the Harnack chain condition,

$$
G(Z, Y) \geq C^{-1} G\left(X_{\Delta}, Y\right), \quad \forall Z \in S_{1}
$$

and therefore by the maximum principle we have

$$
\begin{equation*}
G(Z, Y) \geq C^{-1} G\left(X_{\Delta}, Y\right) \omega_{\Delta}^{Z}\left(S_{1}\right), \quad \forall Z \in T_{\Delta} \tag{3.68}
\end{equation*}
$$

Combining (3.66), (3.67) and (3.68), we obtain

$$
\frac{G(Z, X)}{G\left(X_{\Delta}, X\right)} \lesssim \frac{G(Z, Y)}{G\left(X_{\Delta}, Y\right)}, \quad \forall Z \in B_{\Delta} \cap \Omega
$$

The opposite inequality follows by interchanging the roles of $X$ and $Y$.
Corollary 3.69. - Given the same hypotheses as in Lemma 3.64, there is a uniform constant $C$ such that for every pair of surface balls $\Delta:=B \cap \partial \Omega$, and $\Delta^{\prime}:=B^{\prime} \cap \partial \Omega$, with $B^{\prime} \subseteq B$, and for every $X \in \Omega \backslash 2 \kappa_{0} B$, where $\kappa_{0}$ is the constant in (3.63), we have

$$
\frac{1}{C} \omega^{X_{\Delta}}\left(\Delta^{\prime}\right) \leq \frac{\omega^{X}\left(\Delta^{\prime}\right)}{\omega^{X}(\Delta)} \leq C \omega^{X_{\Delta}}\left(\Delta^{\prime}\right)
$$

Proof. - We follow [31, Corollary 1.3.8]. Fix $\Delta^{\prime}, \Delta, B^{\prime}, B$ and $X$ as in the statement of the present corollary. Set $B^{* *}=\kappa_{1} B$ and $\Delta^{* *}:=B^{* *} \cap \partial \Omega$, where we may choose $\kappa_{1}$ large enough, depending only on $\kappa_{0}$ and on the constants in the Corkscrew condition, such that $X_{\Delta^{* *}} \in \Omega \backslash 2 \kappa_{0} B$. Let $r^{\prime}$ and $r$ denote the respective radii of $B^{\prime}$ and $B$. By Lemma 3.30, we have

$$
\begin{aligned}
\omega^{X}\left(\Delta^{\prime}\right) & \approx\left(r^{\prime}\right)^{n-1} G\left(X_{\Delta^{\prime}}, X\right) \\
\omega^{X}(\Delta) & \approx r^{n-1} G\left(X_{\Delta}, X\right) \\
\omega^{X_{\Delta}}\left(\Delta^{\prime}\right) & \approx \omega^{X_{\Delta^{* *}}}\left(\Delta^{\prime}\right)\left(r^{\prime}\right)^{n-1} G\left(X_{\Delta^{\prime}}, X_{\Delta^{* *}}\right)
\end{aligned}
$$

where in the third line we have also used the Harnack chain condition. Moreover, by Lemma 3.11 and the Harnack chain condition, we have

$$
r^{n-1} G\left(X_{\Delta}, X_{\Delta^{* *}}\right) \approx 1
$$

Note that $X_{\Delta^{\prime}} \subset B^{\prime} \subset B$. Thus, by Lemma 3.64, we have

$$
\frac{G\left(X_{\Delta^{\prime}}, X\right)}{G\left(X_{\Delta^{\prime}}, X_{\Delta^{* *}}\right)} \approx \frac{G\left(X_{\Delta}, X\right)}{G\left(X_{\Delta}, X_{\Delta^{* *}}\right)},
$$

and the conclusion of the corollary follows.

## 4. Harnack chains imply a Poincaré inequality

In this section we prove that a certain Poincaré inequality holds in any domain $\Omega$ satisfying the ADR, Corkscrew and Harnack chain properties. We therefore impose those three hypotheses throughout this section. It will be convenient to set some additional notation. As above, we let $W$ denote the collection of Whitney cubes of $\Omega$, and we recall that these have been constructed so that for $I \in \mathcal{W}$, we have $\operatorname{dist}(4 I, \partial \Omega) \approx \ell(I)$ (cf. (3.42)). Given a pairwise disjoint family $\mathcal{F} \in \mathbb{D}$, and a constant $\rho>0$, we derive from $\mathcal{F}$ another family $\mathcal{F}(\rho) \subset \mathbb{D}$, as follows. We augment $\mathcal{F}$ by adjoining to it all those $Q \in \mathbb{D}$ of side length $\ell(Q) \leq \rho$, and we denote this augmented collection by $\mathfrak{C}(\mathcal{F}, \rho)$. We then let $\mathcal{F}(\rho)$ denote the collection of the maximal cubes of $\mathfrak{C}(\mathcal{F}, \rho)$. Thus, the corresponding discrete sawtooth $\mathbb{D}_{\mathscr{G}(\rho)}$ consists precisely of those $Q \in \mathbb{D}_{\mathcal{G}}$ such that $\ell(Q)>\rho$.

Having constructed the family $\mathcal{F}(\rho)$, and given $Q \in \mathbb{D}$ with $\ell(Q)>\rho$, we may then define local discrete and geometric sawtooth regions $\mathbb{D}_{\mathcal{G}(\rho), Q}$ and $\Omega_{\mathscr{G}(\rho), Q}$ with respect to this family as in (3.41)-(3.47) and (3.54).
$4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - No 3

We shall also find it useful to consider certain "fattened" versions of the sawtooth regions, as follows. Bearing in mind (3.42), we set

$$
\begin{align*}
U_{Q}^{\mathrm{fat}} & :=\bigcup_{W_{Q}^{*}} 4 I,  \tag{4.1}\\
T_{Q}^{\mathrm{fat}} & :=\operatorname{int}\left(\bigcup_{Q^{\prime} \in \mathbb{D}_{Q}} U_{Q^{\prime}}^{\mathrm{fat}}\right),  \tag{4.2}\\
\Omega_{G, Q}^{\mathrm{fat}} & :=\operatorname{int}\left(\bigcup_{Q^{\prime} \in \mathbb{D}_{g, Q}} U_{Q^{\prime}}^{\mathrm{fat}}\right) \tag{4.3}
\end{align*}
$$

(compare to (3.47), (3.52) and (3.54)). We note that, by construction,

$$
\begin{array}{ll}
\delta(X) \gtrsim \rho, & \text { if } X \in \Omega_{\mathscr{G}(\rho), Q}^{\mathrm{fat}} \\
\delta(X) \lesssim \rho, & \text { if } X \in \Omega_{\mathscr{G}, Q} \backslash \Omega_{\mathscr{G}(\rho), Q} \tag{4.5}
\end{array}
$$

Given a pairwise disjoint family $\mathcal{F} \in \mathbb{D}$, and a cube $Q \in \mathbb{D}_{\mathscr{F}}$, we define

$$
\begin{equation*}
\mathcal{W}_{\mathcal{F}}:=\bigcup_{Q^{\prime} \in \mathbb{D}_{\mathcal{G}}} W_{Q^{\prime}}^{*}, \quad \mathcal{W}_{\mathcal{F}, Q}:=\bigcup_{Q^{\prime} \in \mathbb{D}_{\mathcal{G}, Q}} W_{Q^{\prime}}^{*} \tag{4.6}
\end{equation*}
$$

so that in particular, we may write

$$
\begin{equation*}
\Omega_{\mathcal{G}, Q}=\operatorname{int}\left(\bigcup_{I \in \mathcal{W}_{\mathcal{G}, Q}} I^{*}\right), \quad \Omega_{\mathcal{F}, Q}^{\mathrm{fat}}=\operatorname{int}\left(\bigcup_{I \in \mathcal{W}_{\mathcal{G}, Q}} 4 I\right), \tag{4.7}
\end{equation*}
$$

where we recall that $I^{*}:=(1+\lambda) I$.
Suppose now that $Q \in \mathbb{D}_{\mathcal{G}}$, let $r \approx \ell(Q)$ and fix a small $\varepsilon>0$. Then for $I, J \in \mathcal{W}_{\mathcal{G}(\varepsilon r), Q}$, we have $\ell(I) \approx \ell(J)$ and $\operatorname{dist}(I, J) \lesssim \ell(I)$ (where the implicit constants depend upon $\varepsilon$ ). By Lemma 3.61 there is a chain $\left\{I_{1}, I_{2}, \ldots, I_{N}\right\} \subset \mathcal{W}_{\mathscr{G}(\varepsilon r), Q}$, of bounded cardinality $N$ depending only on dimension, the Harnack chain constants, and $\varepsilon$, such that $I_{1}=J, I_{N}=I$, $\ell\left(I_{j}\right) \approx \ell(I)$ for each $j$ (again the implicit constants depend upon $\varepsilon$ ), and for which $\bigcup_{j=1}^{N} I_{j}^{*}$ contains a Harnack chain which connects the centers of $I$ and $J$. Moreover, for $\lambda$ chosen small enough, the chain may be constructed so that for each $j, 1 \leq j \leq N-1$, either $I_{j}^{*} \subset 4 I_{j+1}$, or $I_{j+1}^{*} \subset 4 I_{j}$. In the sequel, we shall refer to such a chain $\left\{I_{1}, I_{2}, \ldots, I_{N}\right\}$ as a "Harnack chain of Whitney cubes connecting $J$ to $I$ " (we beg the reader's indulgence for this mild abuse of terminology: it is of course really the dilates $\left\{I_{j}^{*}\right\}$ which form a Harnack chain).

Lemma 4.8. - Suppose that $\Omega$ is a 1-sided NTA domain with ADR boundary. Fix $Q_{0} \in \mathbb{D}$, and a pairwise disjoint family $\mathcal{F} \subset \mathbb{D}_{Q_{0}}$, and let $Q \in \mathbb{D}_{\mathcal{F}, Q_{0}}$. If $r \approx \ell(Q)$, then for every $p$, $1 \leq p<\infty$, and for every small $\varepsilon>0$, there is a constant $C_{\varepsilon, p}$ such that

$$
\iint_{\Omega_{\mathcal{G}(\varepsilon r), Q}}\left|f-c_{Q, \varepsilon}\right|^{p} \leq C_{\varepsilon, p} r^{p} \iint_{\Omega_{\mathcal{G}(e r), Q}^{\mathrm{fat}}}|\nabla f|^{p},
$$

where $c_{Q, \varepsilon}:=\left|\Omega_{\mathcal{G}(\varepsilon r), Q}\right|^{-1} \iint_{\Omega_{\mathcal{G}(\varepsilon r), Q}} f$.

Proof. - Let $X \in \Omega_{\mathcal{G}(\varepsilon r), Q}$, so that in particular, $X \in I_{X}^{*}$ where $I_{X} \in \mathcal{W}_{\mathcal{G}(\varepsilon r), Q}$, and observe that

$$
\begin{aligned}
\mid f(X)- & \frac{1}{\left|\Omega_{\mathcal{G}(\varepsilon r), Q}\right|} \iint_{\Omega_{\mathcal{G}(\varepsilon r), Q}} f\left|=\left|\frac{1}{\left|\Omega_{\mathcal{G}(\varepsilon r), Q}\right|} \iint_{\Omega_{\mathcal{G}(\varepsilon r), Q}}(f(X)-f(Y)) d Y\right|\right. \\
& \leq \frac{1}{\left|\Omega_{\mathcal{G}(\varepsilon r), Q}\right|} \sum_{I \in W_{\mathcal{G}(e r), Q}} \iint_{I^{*}}\left|f(X)-f_{I_{1}^{*}}+f_{I_{1}^{*}}-\cdots+f_{I_{N}^{*}}-f(Y)\right| d Y,
\end{aligned}
$$

where $I_{1}, \ldots, I_{N}$ is a Harnack chain of Whitney cubes connecting $I_{X}$ to $I$ and $f_{I_{j}^{*}}$ := $\left|I_{j}^{*}\right|^{-1} \iint_{I_{j}^{*}} f$. Of course, $N$ depends upon $\varepsilon$. It also depends upon $I$ in the sum, but in a uniformly bounded manner for $\varepsilon$ fixed. Consequently, it is enough to consider

$$
\iint_{J^{*}}\left(\frac{1}{\left|\Omega_{\mathcal{G}(\varepsilon r), Q}\right|} \iint_{I^{*}}\left|f(X)-f_{I_{1}^{*}}+f_{I_{1}^{*}}-\cdots+f_{I_{N}^{*}}-f(Y)\right| d Y\right)^{p} d X
$$

where $J, I \in \mathcal{W}_{\mathscr{G}(\varepsilon r), Q}$ and are connected by the chain of cubes $I_{1}, I_{2}, \ldots, I_{N}$. The desired bound may now be obtained from the standard Poincaré inequality as follows. First,

$$
\begin{aligned}
& \iint_{J^{*}}\left(\frac{1}{\left|\Omega_{\mathcal{G}(\varepsilon r), Q}\right|} \iint_{I^{*}}\left|f(X)-f_{I_{1}^{*}}\right| d Y\right)^{p} d X \\
& \leq \iint_{J^{*}}\left|f(X)-f_{I_{1}^{*}}\right|^{p} d X \leq C_{p} \ell(J)^{p} \iint_{J^{*}}|\nabla f|^{p},
\end{aligned}
$$

since $I_{1}=J$. Similarly, the contribution of $f_{I_{N}^{*}}-f(Y)$ is bounded by

$$
\iint_{I^{*}}\left|f_{I_{N}^{*}}-f(Y)\right|^{p} d Y \leq C_{p} \ell(I)^{p} \iint_{I^{*}}|\nabla f|^{p},
$$

since $I_{N}=I$. Finally, to handle the contribution of any term $f_{I_{j}^{*}}-f_{I_{j+1}^{*}}$, we observe that

$$
\left|f_{I_{j}^{*}}-f_{I_{j+1}^{*}}\right| \leq\left|f_{I_{j}^{*}}-f_{I_{j}^{* *}}\right|+\left|f_{I_{j}^{* *}}-f_{I_{j+1}^{*}}\right|,
$$

where $I_{j}^{* *}:=4 I_{j}$ or $4 I_{j+1}$, whichever has the larger diameter. Then for example,

$$
\left|f_{I_{j}^{*}}-f_{I_{j}^{* *}}\right| \leq \frac{1}{\left|I_{j}^{*}\right|} \iint_{I_{j}^{*}}\left|f-f_{I_{j}^{* *}}\right| \lesssim \ell\left(I_{j}\right) \frac{1}{\left|I_{j}\right|} \iint_{I_{j}^{* *}}|\nabla f|,
$$

and similarly for the term $\left|f_{I_{j}^{* *}}-f_{I_{j+1}^{*}}\right|$ since, as noted above, $I_{j}^{* *}$ contains both $I_{j}^{*}$ and $I_{j+1}^{*}$. The Poincaré inequality now follows, since each $I \in \mathcal{W}_{\mathcal{G}(\varepsilon r), Q}$, and thus every $I_{j}$ in any of the chains, has side length proportional to $r$, depending on $\varepsilon$.

## 5. A criterion for exterior Corkscrew points

We present a criterion for the existence of Corkscrew points in the domain exterior to a sawtooth region. We begin with a series of lemmas in which we establish some local estimates for the single layer potential operator $\&$ defined in (1.11), and also prove some geometric properties of sawtooth regions and of domains with ADR boundaries.

Lemma 5.1. - Suppose that $E \subset \mathbb{R}^{n+1}$ is $n$-dimensional $A D R$, and let $\kappa>1$. If $1 \leq q<(n+1) / n$, there is a constant $C_{q, \kappa}$ depending only on $n, q, \kappa$ and the ADR constants such that for every $x \in E, B:=B(x, r)$ and $\kappa \Delta:=\kappa B \cap E$, we have

$$
\begin{equation*}
\iint_{B}\left|\nabla \& 1_{\kappa \Delta}(X)\right|^{q} d X \leq C_{q, \kappa} r^{n+1} \tag{5.2}
\end{equation*}
$$

Proof. - The left hand side of (5.2) is crudely dominated by a constant times

$$
\begin{aligned}
\iint_{B} \mid \int_{\kappa \Delta} & \left.\frac{1}{|X-y|^{n}} d H^{n}(y)\right|^{q} d X \leq H^{n}(\kappa \Delta)^{q-1} \iint_{B}\left(\int_{\kappa \Delta} \frac{1}{|X-y|^{n q}} d H^{n}(y)\right) d X \\
& \lesssim r^{n(q-1)} \int_{\kappa \Delta}\left(\iint_{|X-y|<(\kappa+1) r} \frac{1}{|X-y|^{n q}} d X\right) d H^{n}(y) \\
& \approx r^{n(q-1)} r^{n} r^{n+1-n q}=r^{n+1}
\end{aligned}
$$

where of course the implicit constants depend on $\kappa$ and $q$.
Lemma 5.3. - Suppose that $E \subset \mathbb{R}^{n+1}$ is $n$-dimensional $A D R$. For $\rho>0$, define the "boundary strip" $\Sigma_{\rho}:=\left\{X \in \mathbb{R}^{n+1} \backslash E: \operatorname{dist}(X, E)<\rho\right\}$. Then there is a uniform constant $C$ such that for every ball $B:=B(x, r)$ centered on $E$, and for $\rho \leq r$, we have

$$
\begin{equation*}
\left|\Sigma_{\rho} \cap B\right| \leq C \rho r^{n} . \tag{5.4}
\end{equation*}
$$

Proof. - Let $W_{E}$ denote the collection of cubes in the Whitney decomposition of $\mathbb{R}^{n+1} \backslash E$, and for each $k \in \mathbb{Z}$, set $\mathcal{W}_{k}:=\left\{I \in \mathcal{W}_{E}: \ell(I)=2^{-k}\right\}$. For each $I \in \mathcal{W}_{E}$, choose $Q_{I} \in \mathbb{D}(E)$ such that $\ell\left(Q_{I}\right)=\ell(I)$, and $\operatorname{dist}\left(I, Q_{I}\right)=\operatorname{dist}(I, E) \approx \ell(I)$. By ADR, for each $I$ there are at most a bounded number of $Q \in \mathbb{D}(E)$ having these properties, and we just pick one. We note that if $I \cap B$ is non-empty, and if $\ell(I) \lesssim \rho \leq r$, then $Q_{I} \subset \kappa_{1} B$ for some uniform constant $\kappa_{1}$. Moreover, the collection $\left\{Q_{I}\right\}_{I \in W_{k}}$ has bounded overlaps for each fixed $k$. We then have

$$
\begin{align*}
\left|\Sigma_{\rho} \cap B\right| & \leq \sum_{k: 2^{-k} \lesssim \rho} \sum_{I \in W_{k}}|I \cap B| \\
& \approx \sum_{k: 2^{-k} \leq \rho} \sum_{I \in W_{k}}|I \cap B| \ell(I)^{-n} H^{n}\left(Q_{I}\right) \quad(\text { by ADR) }  \tag{5.5}\\
& \lesssim \sum_{k: 2^{-k} \lesssim \rho} 2^{-k} \sum_{I \in W_{k}: Q_{I} \subset \kappa_{1} B} H^{n}\left(Q_{I}\right) \lesssim \rho r^{n},
\end{align*}
$$

where in the last step we have used the bounded overlap property of the $Q_{I}$ 's.
Corollary 5.6. - Let $0<\gamma<1 /(n+1)$, and suppose that $E \subset \mathbb{R}^{n+1}$ is $n$-dimensional $A D R$. Then there is a uniform constant $C_{\gamma, \kappa}$ such that for every ball $B:=B(x, r)$ centered on $E, \kappa \Delta:=\kappa B \cap E$, and every $\rho$ with $0<\rho \leq r$, we have

$$
\iint_{\Sigma_{\rho} \cap B}\left|\nabla \delta 1_{\kappa \Delta}(X)\right| d X \leq C_{\gamma, \kappa} \rho^{\gamma} r^{n+1-\gamma} .
$$

Proof. - The corollary follows immediately from Hölder's inequality and the previous two lemmata. We omit the routine details.

Lemma 5.7. - Suppose that $\partial \Omega$ is $A D R$, and let $B:=B(x, r), \Delta:=B \cap \partial \Omega$, with $r \leq \operatorname{diam}(\partial \Omega)$, and $x \in \partial \Omega$. If

$$
\begin{equation*}
\left|B \cap\left(\mathbb{R}^{n+1} \backslash \bar{\Omega}\right)\right| \geq a r^{n+1} \tag{5.8}
\end{equation*}
$$

for some $a>0$, then there is a point $X_{\Delta}^{-} \in B \backslash \bar{\Omega}$, and a constant $c_{1}$ depending only on $a, n$ and the $A D R$ constants such that

$$
B\left(X_{\Delta}^{-}, c_{1} r\right) \subset \mathbb{R}^{n+1} \backslash \bar{\Omega}
$$

Proof. - For notational convenience, we set

$$
B^{-}:=B \cap\left(\mathbb{R}^{n+1} \backslash \bar{\Omega}\right)
$$

We apply Lemma 5.3 , with $E=\partial \Omega$, and $\rho=\operatorname{ar} /(2 C)$ (notice that without loss of generality we may assume that $a \leq 1, C \geq 1$ ), and (5.8) to deduce that

$$
\left|B^{-} \backslash \Sigma_{a r /(2 C)}\right| \geq \frac{1}{2} a r^{n+1}
$$

In particular $B^{-} \backslash \Sigma_{a r /(2 C)}$ is non-empty. Moreover, by definition of $\Sigma_{\rho}$, we have that $\operatorname{dist}(X, \partial \Omega) \geq a r /(2 C)$, for every $X \in B^{-} \backslash \Sigma_{a r /(2 C)}$. Therefore, any such $X$ may be taken as the point $X_{\Delta}^{-}$, with $c_{1}:=a /(4 C)$.

Remark. - Given a domain $\Omega$, we shall henceforth refer to a Corkscrew point for the domain $\mathbb{R}^{n+1} \backslash \bar{\Omega}$, such as the point $X_{\Delta}^{-}$in the lemma, as an "exterior Corkscrew point".

Lemma 5.9. - Suppose that $\Omega$ is a 1-sided NTA domain with ADR boundary. Let $\mathcal{F} \subset \mathbb{D}$ be a pairwise disjoint family. Then for every $Q \subseteq Q_{j} \in \mathcal{F}$, there is a ball $B^{\prime} \subset \mathbb{R}^{n+1} \backslash \Omega_{\mathcal{F}}$, centered at $\partial \Omega$, with radius $r^{\prime} \approx \ell(Q) / K_{0}$, and $\Delta^{\prime}:=B^{\prime} \cap \partial \Omega \subset Q$.

Proof. - Recall that there exist $B_{Q}:=B\left(x_{Q}, r\right)$ and $\Delta_{Q}:=B_{Q} \cap \partial \Omega \subset Q$, as defined in (1.16) and (1.17), where $r \approx \ell(Q)$. We now set

$$
B^{\prime}=B\left(x_{Q},\left(M K_{0}\right)^{-1} r\right)
$$

where $M$ is a sufficiently large number to be chosen momentarily. We need only verify that $B^{\prime} \cap \Omega_{\mathscr{F}}=\varnothing$. Suppose not. Then by definition of $\Omega_{\mathscr{F}}$, there is a Whitney cube $I \in \mathcal{W}_{\mathscr{G}}$ (cf. (4.6)) such that $I^{*}$ meets $B^{\prime}$. Since $I^{*}$ meets $B^{\prime}$, there is a point $Y_{I} \in I^{*}$ such that

$$
\ell(I) \approx \operatorname{dist}\left(I^{*}, \partial \Omega\right) \leq\left|Y_{I}-x_{Q}\right| \leq r /\left(M K_{0}\right) \approx \ell(Q) /\left(M K_{0}\right)
$$

On the other hand, since $I \in \mathcal{W}_{\mathcal{F}}$, there is a $Q_{I} \in \mathbb{D}_{\mathcal{F}}$ (hence $Q_{I}$ is not contained in $Q_{j}$ ) with $\ell(I) \approx \ell\left(Q_{I}\right)$, and $\operatorname{dist}\left(Q_{I}, Y_{I}\right) \approx \operatorname{dist}\left(Q_{I}, I\right) \lesssim K_{0} \ell(I) \lesssim \ell(Q) / M$. Then by the triangle inequality,

$$
\left|y-x_{Q}\right| \lesssim \ell(Q) / M, \quad \forall y \in Q_{I}
$$

Thus, if $M$ is chosen large enough, $Q_{I} \subset \Delta_{Q} \subset Q \subset Q_{j}$, a contradiction.
We now come to the main lemma of this section.
$4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - No 3

Lemma 5.10. - Suppose that $\Omega$ is a 1-sided NTA domain with ADR boundary. Fix $Q_{0} \in \mathbb{D}$, and a pairwise disjoint family $\mathcal{F} \subset \mathbb{D}_{Q_{0}}$, and let $\Omega_{\mathcal{G}, Q_{0}}$ be the corresponding sawtooth domain. Suppose also that for some $\eta>0$, we have

$$
\begin{equation*}
\sup _{Q \in \mathbb{D}_{Q_{0}}} \frac{1}{\sigma(Q)} \iint_{\Omega_{\xi, Q}^{\mathrm{f}} \mathrm{tat}}\left|\nabla^{2} \& 1(X)\right|^{2} \delta(X) d X \leq \eta \tag{5.11}
\end{equation*}
$$

If $\eta \leq \eta_{0}$ with $\eta_{0}$ small enough, depending only on $n, K_{0}$, and the Corkscrew, Harnack chain and $A D R$ constants for $\Omega$, then for every $B:=B(x, r)$ and $\Delta_{\star}:=B \cap \partial \Omega_{\mathscr{G}, Q_{0}}$, with $x \in \partial \Omega_{\mathcal{G}, Q_{0}}$ and $r \leq \operatorname{diam}\left(Q_{0}\right)$, there is an exterior Corkscrew point $X_{\Delta_{*}} \in B \cap\left(\mathbb{R}^{n+1} \backslash \Omega_{\mathcal{G}, Q_{0}}\right)$. Moreover, the exterior Corkscrew constants depend only upon $\eta_{0}, K_{0}$, and the other parameters stated above.

To avoid confusion, we note that, as usual, $\delta(X):=\operatorname{dist}(X, \partial \Omega)$, and $\Delta=B \cap \Omega$ denotes a surface ball on $\partial \Omega$; we shall use the notation $\delta_{\star}(X):=\operatorname{dist}\left(X, \partial \Omega_{\mathcal{F}, Q_{0}}\right)$, and $\Delta_{\star}:=B \cap \partial \Omega_{\mathcal{G}, Q_{0}}$.

Proof. - We fix $x \in \partial \Omega_{\mathscr{G}, Q_{0}}$, and consider two separate cases. Let $M$ be a sufficiently large constant, to be chosen, whose value will remain fixed throughout the proof of the present lemma.

CASE 1: $\operatorname{dist}(x, \partial \Omega)>r /\left(M K_{0}\right)$.
In this case, $x \in \partial I^{*} \cap J$, where as usual $I^{*}=(1+\lambda) I$ and $I \in \mathcal{W}_{\mathcal{F}, Q_{0}}$ (cf. (4.6)), and where $J \in \mathcal{W}$ with $\tau J \subset \mathbb{R}^{n+1} \backslash \Omega_{\mathcal{G}, Q_{0}}$, for some $\tau \in(1 / 2,1)$ (cf. (3.51)). By the nature of Whitney cubes, we have $\ell(I) \approx \ell(J) \approx \operatorname{dist}(x, \partial \Omega)>r /\left(M K_{0}\right)$. In this case, it is evident that there is a Corkscrew point in $J$, with $c \approx\left(M K_{0}\right)^{-1}$.

CASE 2: $\operatorname{dist}(x, \partial \Omega) \leq r /\left(M K_{0}\right)$.
In this case, either $x \in \partial \Omega \cap \partial \Omega_{\mathcal{G}, Q_{0}}$, or else $x$ lies on a face of $I^{*}$ for some Whitney cube $I \in \mathcal{W}_{\mathcal{G}, Q_{0}}$, with $\ell(I) \lesssim r /\left(M K_{0}\right)$. In the former scenario, by Proposition 6.1 below, we may choose $Q \in \mathbb{D}_{Q_{0}}$, with $x \in \bar{Q} \subset B$, and $\ell(Q) \approx r$. If $Q \subseteq Q_{j}$, for some $Q_{j} \in \mathcal{F}$ (which might happen if $x \in \partial Q_{j}$ ), then by Lemma 5.9 we immediately obtain the existence of the desired exterior Corkscrew point for $\Omega_{\mathcal{G}, Q_{0}}$, at the scale $r$. Thus, in this scenario, it is enough to suppose that $Q$ is not contained in any $Q_{j} \in \mathcal{F}$.

Otherwise, if $x \in \partial I^{*}$ for some $I \in \mathcal{W}_{\mathscr{G}, Q_{0}}$, with $\ell(I) \lesssim r /\left(M K_{0}\right)$, then there is a $Q_{I} \in \mathbb{D}_{\mathcal{G}, Q_{0}}$ such that $\ell\left(Q_{I}\right) \approx \ell(I)$, and $\operatorname{dist}\left(Q_{I}, I\right) \lesssim K_{0} \ell(I) \lesssim r / M$. Consequently, we have $\operatorname{dist}(I, Q) \lesssim r / M$ for any $Q \in \mathbb{D}$ with $Q_{I} \subseteq Q \subseteq Q_{0}$. Choosing $M$ large enough, we may then fix such a $Q$ with $\ell(Q) \approx r$, and $Q \subset B$. If $Q$ is contained in some $Q_{j} \in \mathcal{F}$, then by Lemma 5.9, we again obtain the existence of an exterior Corkscrew point exactly as before.

Therefore, in either scenario, we have reduced matters to the following situation: there is a $Q \in \mathbb{D}_{\mathcal{G}, Q_{0}}$ (i.e., not contained in any $Q_{j} \in \mathcal{F}$ ), with $\ell(Q) \approx r$, and $Q \subset B$. Having fixed this $Q$, we recall that, by Lemma 3.55, there is a ball $B_{Q}^{\prime}:=B\left(x_{Q}, s\right)$, with radius $s \approx\left(K_{0}\right)^{-1} \ell(Q)$, such that (3.57) holds.

By Lemma 3.61, the sawtooth domain $\Omega_{\mathcal{F}, Q_{0}}$ inherits the 1 -sided NTA (i.e., interior Corkscrew and Harnack chain) and ADR properties from $\Omega$. Thus, by Lemma 5.7, applied with $\Omega_{\mathcal{G}, Q_{0}}$ in place of $\Omega$, and $B_{Q}^{\prime}$ in place of $B$, it is enough to establish the analogue of (5.8) with $a$ depending only on the allowable parameters.

To this end, we proceed by a variant of the argument in [19, pp. 254-256]. We remind the reader of the definition of the family $\mathcal{F}(\rho)$ (see the discussion at the beginning of Section 4), and we also note that, by construction, there is a purely dimensional constant $C_{n}$ such that

$$
\begin{equation*}
T_{Q}^{\mathrm{fat}} \subset B\left(x_{Q}, C_{n} K_{0} \ell(Q)\right)=: B_{Q}^{*} \tag{5.12}
\end{equation*}
$$

Set $\Delta_{Q}^{*}:=B_{Q}^{*} \cap \partial \Omega$ and $\Delta_{Q}^{\prime}:=B_{Q}^{\prime} \cap \partial \Omega$, and let $\Phi \in C_{0}^{\infty}\left(B_{Q}^{\prime}\right)$, with $\Phi \equiv 1$ on $B\left(x_{Q}, s / 2\right)$, $0 \leq \Phi \leq 1$, and $\|\nabla \Phi\|_{\infty} \lesssim s^{-1}$. Let $\mathscr{L}:=\nabla \cdot \nabla$ denote the usual Laplacian in $\mathbb{R}^{n+1}$. By the ADR property, and the fact that $s \approx \ell(Q) \approx r$, we have

$$
\begin{align*}
r^{n+1} & \approx r \sigma\left(\frac{1}{2} \Delta_{Q}^{\prime}\right) \leq r \int_{\partial \Omega} \Phi d \sigma=r\langle-\mathcal{L} \& 1, \Phi\rangle  \tag{5.13}\\
& =r \iint_{\mathbb{R}^{n+1}}\left(\nabla \& 1(X)-\nabla \& 1_{\left(2 \Delta_{Q}^{*}\right)^{c}}\left(x_{Q}\right)-\vec{\alpha}\right) \cdot \nabla \Phi(X) d X \\
& \lesssim \iint_{B_{Q}^{\prime}} \mid \nabla \& 1(X)-\nabla \& 1_{\left(2 \Delta_{Q}^{*}\right)^{c}\left(x_{Q}\right)-\vec{\alpha} \mid d X=\iint_{\Omega \cap B_{Q}^{\prime}}+\iint_{\Omega_{\mathrm{ext}} \cap B_{Q}^{\prime}}} \quad=\iint_{\Omega_{\mathcal{G}(\varepsilon r), Q} \cap B_{Q}^{\prime}}+\iint_{\left(\Omega_{\left.\mathscr{G}, Q_{0} \backslash \Omega_{\mathscr{G}(\varepsilon r), Q}\right) \cap B_{Q}^{\prime}}+\iint_{\left(\Omega \backslash \Omega_{\mathcal{G}, Q_{0}}\right) \cap B_{Q}^{\prime}}+\iint_{\Omega_{\mathrm{ext}} \cap B_{Q}^{\prime}}\right.} \\
& =I+I I+I I I+I V
\end{align*}
$$

where $\vec{\alpha}$ is a constant vector at out disposal, $\varepsilon>0$ is a small number to be determined, and where as above, $\Omega_{\mathrm{ext}}:=\mathbb{R}^{n+1} \backslash \bar{\Omega}$.

We now set

$$
\vec{\alpha}:=\frac{1}{\left|\Omega_{\mathcal{F}(\varepsilon r), Q}\right|} \iint_{\Omega_{\mathcal{G}(\varepsilon r), Q}}\left(\nabla \& 1(X)-\nabla \& 1_{\left(2 \Delta_{Q}^{*}\right)^{c}}\left(x_{Q}\right)\right) d X
$$

We note for future reference that by standard Calderón-Zygmund estimates,

$$
\begin{equation*}
\left|\nabla \& 1_{\left(2 \Delta_{Q}^{*}\right)^{c}}(X)-\nabla \& 1_{\left(2 \Delta_{Q}^{*}\right)^{c}}\left(x_{Q}\right)\right| \leq C, \quad \forall X \in B_{Q}^{*} \tag{5.14}
\end{equation*}
$$

We also note that by Lemma 3.61, the sawtooth domain $\Omega_{\mathcal{G}(\varepsilon r), Q}$, if non-empty, must contain a Corkscrew point at the scale of $\ell(Q) \approx r$, so that, in particular,

$$
r^{n+1} \lesssim\left|\Omega_{\mathcal{F}(\varepsilon r), Q}\right|
$$

Consequently, by (5.12) and the fact that $\Omega_{\mathscr{F}, Q} \subset T_{Q} \subset T_{Q}^{\text {fat }}$ for any pairwise disjoint family $\mathcal{F}$ and every $Q \in \mathbb{D}$, we have

$$
\begin{equation*}
|\vec{\alpha}| \leq \frac{C_{K_{0}}}{\left|B_{Q}^{*}\right|} \iint_{B_{Q}^{*}}\left|\nabla \& 1_{2 \Delta_{Q}^{*}}(X)\right| d X+C \leq C_{K_{0}} \tag{5.15}
\end{equation*}
$$

where in the last step we have used Lemma 5.1.

By the Poincaré inequality (Lemma 4.8), (5.12), and (4.4) with $\rho=\varepsilon r$, we obtain

$$
\begin{aligned}
I & \leq C_{\varepsilon} r \iint_{\Omega_{G(e r), Q}^{\mathrm{fat}}}\left|\nabla^{2} \& 1(X)\right| d X \\
& \leq C_{\varepsilon, K_{0}} r^{(n+3) / 2}\left(\iint_{\Omega_{G(\varepsilon r), Q}^{\mathrm{fat}}}\left|\nabla^{2} \& 1(X)\right|^{2} d X\right)^{1 / 2} \\
& \leq C_{\varepsilon, K_{0}} r^{(n+2) / 2}\left(\iint_{\Omega_{g(e r), Q}^{\mathrm{fat}}}\left|\nabla^{2} \& 1(X)\right|^{2} \delta(X) d X\right)^{1 / 2} \\
& \leq C_{\varepsilon, K_{0}} \sqrt{\eta} r^{n+1},
\end{aligned}
$$

by Hypothesis (5.11), since $\ell(Q) \approx r$, and $\Omega_{\mathcal{G}(\varepsilon r), Q}^{\mathrm{fat}} \subset \Omega_{\mathcal{G}, Q}^{\mathrm{fat}}$.
Next, we claim that, for each $\gamma \in(0,1 /(n+1))$, we have

$$
\begin{equation*}
I I \leq C_{\gamma, K_{0}} \varepsilon^{\gamma} r^{n+1} \tag{5.16}
\end{equation*}
$$

We defer the proof of this claim momentarily, and observe that

$$
I I I+I V=\iint_{\left(\mathbb{R}^{n+1} \backslash \overline{\Omega_{G, Q_{0}}}\right) \cap B_{Q}^{\prime}}\left|\nabla \& 1(X)-\nabla \& 1_{\left(2 \Delta_{Q}^{*}\right)^{c}}\left(x_{Q}\right)-\vec{\alpha}\right| d X
$$

(to avoid possible confusion, we point out that the boundaries of $\Omega$ and all of its sub-domains that we consider here, have $(n+1)$-dimensional Lebesgue measure equal to zero). Then by (5.14), (5.15), Hölder's inequality and Lemma 5.1, we deduce that for any $q \in(1,(n+1) / n)$,

$$
I I I+I V \leq C_{K_{0}}\left(\left|\left(\mathbb{R}^{n+1} \backslash \overline{\Omega_{\mathcal{F}, Q_{0}}}\right) \cap B_{Q}^{\prime}\right|^{1 / q^{\prime}} r^{(n+1) / q}+\left|\left(\mathbb{R}^{n+1} \backslash \overline{\Omega_{\mathcal{F}, Q_{0}}}\right) \cap B_{Q}^{\prime}\right|\right) .
$$

Now, choosing first $\varepsilon$, and then $\eta$ sufficiently small, we can hide $I+I I$ on the left hand side of (5.13). Our estimate for $I I I+I V$ then implies that

$$
r^{n+1} \leq C_{K_{0}}\left|\left(\mathbb{R}^{n+1} \backslash \overline{\Omega_{\mathcal{F}, Q_{0}}}\right) \cap B_{Q}^{\prime}\right|
$$

As noted above, the existence of an exterior Corkscrew point now follows by applying Lemma 5.7, with $B_{Q}^{\prime}$ in place of $B$, and $\Omega_{\mathcal{F}, Q_{0}}$ in place of $\Omega$.

To complete the proof of Lemma 5.10, it remains only to prove the claimed estimate (5.16). By (3.57), we may replace $\Omega_{\mathscr{G}, Q_{0}}$ by $\Omega_{\mathcal{G}, Q}$ in the domain of integration which defines $I I$. Consequently, by (4.5) with $\rho=\varepsilon r$, we have that

$$
I I \leq \int_{\Sigma_{C \varepsilon r} \cap B_{Q}^{\prime}}\left|\nabla \& 1(X)-\nabla \& 1_{\left(2 \Delta_{Q}^{*}\right)^{c}}\left(x_{Q}\right)-\vec{\alpha}\right| d X
$$

where $\Sigma_{\rho}:=\left\{X \in \mathbb{R}^{n+1}: \delta(X)<\rho\right\}$. The desired bound now follows readily from (5.14), (5.15), Lemma 5.3 and Corollary 5.6. We omit the routine details.

We conclude this section with an estimate for harmonic measure in "good" sawtooth regions (that is, those for which (5.11) holds for sufficiently small $\eta$ ). Given a subdomain $\Omega^{\prime} \subset \Omega$, we shall use the notational convention that $\omega_{\star}^{X}$ denotes harmonic measure for $\Omega^{\prime}$ with pole at $X$, when there is no chance for confusion.

Corollary 5.17. - Suppose that $\Omega$ is a 1 -sided NTA domain with ADR boundary. Suppose also that (5.11) holds for some $Q_{0} \in \mathbb{D}$, and some pairwise disjoint family $\mathcal{F} \subset \mathbb{D}_{Q_{0}}$, with $\eta \leq \eta_{0}$ (cf. Lemma 5.10). Let $\omega_{\star}^{X}$ denote harmonic measure for $\Omega_{\mathscr{F}, Q_{0}}$ with pole at $X$. Then, for every $x \in \partial \Omega_{\mathscr{G}, Q_{0}}$, every $r \leq \operatorname{diam}\left(Q_{0}\right)$, and every surface ball $\Delta_{\star}=\Delta_{\star}(x, r)$, the harmonic measure $\omega_{\star}^{X \Delta_{\star}}$ belongs to $A_{\infty}\left(\Delta_{\star}\right)$ (cf. Definition 1.19), with uniform $A_{\infty}$ constants depending only upon dimension and the $A D R$, Harnack chain and Corkscrew constants, including $K_{0}$.

Proof. - By Lemma 3.61, $\Omega_{\mathcal{F}, Q_{0}}$ is a 1 -sided NTA domain with ADR boundary. Moreover, by Lemma 5.10, it also satisfies an exterior Corkscrew condition. The conclusion of the corollary now follows immediately by [17, Theorem 2].

## 6. F-Projections and a Dahlberg-Jerison-Kenig "sawtooth lemma"

In this section, we present a dyadic version of the main lemma of [16]. Our approach here is modeled on an analogous result in the Euclidean case which appeared in our previous work [27] (see also [26]). As in [27], we shall utilize a certain projection operator adapted to a pairwise disjoint family $\mathcal{F}$.

Consider now such a family $\mathcal{F}=\left\{Q_{j}\right\} \subset \mathbb{D}$. The projection operator $\mathscr{\mathscr { G }}_{\mathscr{G}}$ associated to $\mathcal{F}$ (the " $\mathcal{F}$-projection operator") is defined by:

$$
\mathscr{P}_{\mathcal{G}} f(x):=f(x) 1_{\partial \Omega \backslash\left(\cup_{\mathscr{G}} Q_{j}\right)}(x)+\sum_{\mathcal{G}}\left(f_{Q_{j}} f d \sigma\right) 1_{Q_{j}}(x) .
$$

We may naturally extend $\mathscr{P}_{\mathcal{G}}$ to act on non-negative Borel measures on $\partial \Omega$. Suppose that $\mu$ is such a measure, and let $A \subset \partial \Omega$. We then define the measure $\mathscr{\mathscr { F }}_{\mathscr{F}} \mu$ as follows:

$$
\mathscr{P}_{\mathscr{F}} \mu(A):=\int_{\partial \Omega} \mathscr{P}_{\mathscr{G}}\left(1_{A}\right) d \mu=\mu\left(A \backslash \bigcup_{\mathscr{F}} Q_{j}\right)+\sum_{\mathcal{F}} \frac{\sigma\left(A \cap Q_{j}\right)}{\sigma\left(Q_{j}\right)} \mu\left(Q_{j}\right)
$$

In particular, we have that $\mathcal{P}_{\mathscr{G}} \mu(Q)=\mu(Q)$, for every $Q \in \mathbb{D}_{\mathcal{G}}$ (i.e., for $Q$ not contained in any $\left.Q_{j} \in \mathscr{F}\right)$, and also that $\mathscr{P}_{\mathscr{F}} \mu\left(Q_{j}\right)=\mu\left(Q_{j}\right)$ for every $Q_{j} \in \mathcal{F}$.

We shall prove a version of the main lemma in [16] which is valid for $\mathscr{F}$-projections of harmonic measure. Our proof follows the idea of the argument in [16], but is technically simpler (given certain geometric preliminaries), owing to the dyadic setting in which we work here. In more precise detail, we follow our earlier Euclidean version of this lemma, which appears in [27, Lemma A.1].

Let us set a bit of notation: given $Q_{0} \in \mathbb{D}$, a pairwise disjoint family $\mathcal{F} \subset \mathbb{D}$, and the corresponding sawtooth domain $\Omega_{\mathcal{F}, Q_{0}}$ (cf. (3.39)-(3.47) and (3.54); also (4.6) and (4.7)), we let $\Delta_{\star}, \delta_{\star}$, and $\omega_{\star}^{X}$ denote, respectively, a surface ball on $\partial \Omega_{\mathcal{G}, Q_{0}}$, the distance to the boundary of $\partial \Omega_{\mathcal{G}, Q_{0}}$, and harmonic measure for the domain $\Omega_{\mathcal{G}, Q_{0}}$ with pole at $X$; i.e., for $x \in \partial \Omega_{\mathcal{G}, Q_{0}}, \Delta_{\star}(x, r)=B(x, r) \cap \partial \Omega_{\mathscr{F}, Q_{0}}$, and for $X \in \Omega_{\mathscr{G}, Q_{0}}$, $\delta_{\star}(X):=\operatorname{dist}\left(X, \partial \Omega_{\mathcal{G}, Q_{0}}\right)$. We continue to use $\Delta=\Delta(x, r), \delta(X)$ and $\omega^{X}$ to denote the analogous objects in reference to the original domain $\Omega$ and its boundary.

Before stating our sawtooth lemma, let us record some useful geometric observations. We recall that by Lemma 3.61, the sawtooth domain $\Omega_{\mathcal{F}, Q_{0}}$ inherits the 1 -sided NTA (i.e., interior Corkscrew and Harnack chain) and ADR properties from $\Omega$. We begin with the following.

Proposition 6.1. - Suppose that $\Omega$ is a 1-sided NTA domain with ADR boundary. Fix $Q_{0} \in \mathbb{D}$, and let $\mathcal{F} \subset \mathbb{D}_{Q_{0}}$ be a disjoint family. Then

$$
\begin{equation*}
Q_{0} \backslash\left(\bigcup_{\mathcal{F}} Q_{j}\right) \subset \partial \Omega \cap \partial \Omega_{\mathcal{G}, Q_{0}} \subset \overline{Q_{0}} \backslash\left(\bigcup_{\mathcal{F}} \operatorname{int}\left(Q_{j}\right)\right) . \tag{6.2}
\end{equation*}
$$

Proof. - We first prove the right hand containment. Suppose that $x \in \partial \Omega \cap \partial \Omega_{\mathscr{G}, Q_{0}}$. Then there is a sequence $X^{k} \in \Omega_{\mathcal{Y}, Q_{0}}$, with $X^{k} \rightarrow x$. By definition of $\Omega_{\mathcal{Y}, Q_{0}}$, each $X^{k}$ is contained in $I_{k}^{*}$ for some $I_{k} \in \mathcal{W}_{\mathcal{G}, Q_{0}}$ (cf. (4.6)-(4.7)), so that $\ell\left(I_{k}\right) \approx \delta\left(X^{k}\right) \rightarrow 0$. Moreover, again by definition, each $I_{k}$ belongs to some $\mathcal{W}_{Q^{k}}^{*}, Q^{k} \in \mathbb{D}_{\mathcal{G}}$ so that,

$$
\operatorname{dist}\left(Q^{k}, I_{k}\right) \lesssim K_{0} \ell\left(Q^{k}\right) \approx K_{0} \ell\left(I_{k}\right) \rightarrow 0
$$

Consequently, $\operatorname{dist}\left(Q^{k}, x\right) \rightarrow 0$. Since each $Q^{k} \subset Q_{0}$, we have $x \in \overline{Q_{0}}$. On the other hand, if $x \in \operatorname{int}\left(Q_{j}\right)$, for some $Q_{j} \in \mathcal{F}$, then there is an $\varepsilon>0$ such that $\operatorname{dist}(x, Q)>\varepsilon$ for every $Q \in \mathbb{D}_{\mathcal{G}, Q_{0}}$ with $\ell(Q) \ll \varepsilon$, because no $Q \in \mathbb{D}_{\mathcal{G}, Q_{0}}$ can be contained in any $Q_{j}$. Since this cannot happen if $\ell\left(Q^{k}\right)+\operatorname{dist}\left(Q^{k}, x\right) \rightarrow 0$, the right hand containment is established.

Now suppose that $x \in Q_{0} \backslash\left(\bigcup_{\mathcal{G}} Q_{j}\right)$. By definition, if $x \in Q \in \mathbb{D}_{Q_{0}}$, then $Q \in \mathbb{D}_{\mathcal{G}, Q_{0}}$. Therefore, we may choose a sequence $\left\{Q^{k}\right\} \subset \mathbb{D}_{\mathcal{G}, Q_{0}}$ shrinking to $x$, whence there exist $I_{k} \in \mathcal{W}_{Q^{k}}^{*} \subset \mathcal{W}_{\mathcal{G}, Q_{0}}$ with $\operatorname{dist}\left(I_{k}, x\right) \rightarrow 0$. The left hand containment now follows.

Proposition 6.3. - Suppose that $\partial \Omega$ is $A D R$, and that $\mu$ is a doubling measure on $\partial \Omega$; i.e., there is a uniform constant $M_{0}$ such that $\mu(2 \Delta) \leq M_{0} \mu(\Delta)$ for every surface ball $\Delta$. Then $\partial Q:=\bar{Q} \backslash \operatorname{int}(Q)$ has $\mu$-measure 0 , for every $Q \in \mathbb{D}$. In particular, the sets in (6.2) have the same $\mu$ measure.

Proof. - The argument is a refinement of that in [21, p. 403], where the Euclidean case was treated. Fix an integer $k$, a cube $Q \in \mathbb{D}_{k}$, and a positive integer $m$ to be chosen. We set

$$
\left\{Q_{i}^{1}\right\}:=\mathbb{D}^{1}:=\mathbb{D}_{Q} \cap \mathbb{D}_{k+m},
$$

and make the disjoint decomposition $Q=\cup Q_{i}^{1}$. We then split $\mathbb{D}^{1}=\mathbb{D}^{1,1} \cup \mathbb{D}^{1,2}$, where $Q_{i}^{1} \in \mathbb{D}^{1,1}$ if $\partial Q_{i}^{1}$ meets $\partial Q$, and $Q_{i}^{1} \in \mathbb{D}^{1,2}$ otherwise. We then write $\bar{Q}=R^{1,1} \cup R^{1,2}$, where

$$
R^{1,1}:=\bigcup_{\mathbb{D}^{1,1}} \widehat{Q}_{i}^{1}, \quad R^{1,2}:=\bigcup_{\mathbb{D}^{1,2}} Q_{i}^{1},
$$

and for each cube $Q_{i}^{1} \in \mathbb{D}^{1,1}$, we construct $\widehat{Q}_{i}^{1}$ as follows. We enumerate the elements in $\mathbb{D}^{1,1}$ as $Q_{i_{1}}^{1}, Q_{i_{2}}^{1}, \ldots, Q_{i_{N}}^{1}$, and then set $\left(Q_{i}^{1}\right)^{*}=Q_{i}^{1} \cup\left(\partial Q_{i}^{1} \cap \partial Q\right)$ and

$$
\widehat{Q}_{i_{1}}^{1}:=\left(Q_{i_{1}}^{1}\right)^{*}, \quad \widehat{Q}_{i_{2}}^{1}:=\left(Q_{i_{2}}^{1}\right)^{*} \backslash\left(Q_{i_{1}}^{1}\right)^{*}, \quad \widehat{Q}_{i_{3}}^{1}:=\left(Q_{i_{3}}^{1}\right)^{*} \backslash\left(\left(Q_{i_{1}}^{1}\right)^{*} \cup\left(Q_{i_{2}}^{1}\right)^{*}\right), \ldots
$$

so that $R^{1,1}$ covers $\partial Q$ and the modified cubes $\widehat{Q}_{i}^{1}$ are pairwise disjoint.
We recall the surface ball $\Delta_{Q}=\Delta\left(x_{Q}, r\right) \subset Q$, with $r \approx \ell(Q)$ as in (1.16)-(1.17). Then

$$
\operatorname{dist}\left(\Delta\left(x_{Q}, r / 2\right), \partial Q\right) \geq \frac{r}{2} \geq c_{0} \ell(Q)=c_{0} 2^{-k},
$$

for some uniform constant $c_{0}$. By Lemma ??, there is a uniform constant $C_{1}$ such that $\operatorname{diam}\left(Q^{\prime}\right) \leq C_{1} \ell\left(Q^{\prime}\right)$, for every $Q^{\prime} \in \mathbb{D}$. We may therefore choose $m$ depending only on the ADR constants and dimension so that $2^{-m}<c_{0} / C_{1}$, whence

$$
\operatorname{diam}\left(Q_{i}^{1}\right) \leq C_{1} 2^{-k-m}<c_{0} 2^{-k}
$$

Consequently, $R^{1,1}$ misses $\Delta\left(x_{Q}, r / 2\right)$, so that by the doubling property,

$$
\mu(\bar{Q}) \leq C_{M_{0}} \mu\left(\Delta\left(x_{Q}, r / 2\right) \leq C_{M_{0}} \mu\left(R^{1,2}\right)\right.
$$

Since $R^{1,1}$ and $R^{1,2}$ are disjoint, the latter estimate yields

$$
\mu\left(R^{1,1}\right) \leq\left(1-\frac{1}{C_{M_{0}}}\right) \mu(\bar{Q})=: \theta \mu(\bar{Q})
$$

where we note that $\theta<1$.
Let us now repeat this procedure, decomposing $\widehat{Q}_{i}^{1}$ for each $Q_{i}^{1} \in \mathbb{D}^{1,1}$. We set $\mathbb{D}^{2}\left(Q_{i}^{1}\right)=\mathbb{D}_{Q_{i}^{1}} \cap \mathbb{D}_{k+2 m}$ and split it into $\mathbb{D}^{2,1}\left(Q_{i}^{1}\right)$ and $\mathbb{D}^{2,2}\left(Q_{i}^{1}\right)$ where $Q^{\prime} \in \mathbb{D}^{2,1}\left(Q_{i}^{1}\right)$ if $\partial Q^{\prime}$ meets $\partial Q \cap \widehat{Q}_{i}^{1}$ (this set plays the role of $\partial Q$ in the previous step). Associated to any $Q^{\prime} \in \mathbb{D}^{2,1}\left(Q_{i}^{1}\right)$ we set $\left(Q^{\prime}\right)^{*}=\left(Q^{\prime} \cap \widehat{Q}_{i}^{1}\right) \cup\left(\partial Q^{\prime} \cap\left(\partial Q \cap \widehat{Q}_{i}^{1}\right)\right)$. Then we make these sets disjoint as before and we have that $R^{2,1}\left(Q_{i}^{1}\right)$ is defined as the disjoint union of the corresponding $\widehat{Q^{\prime}}$. Note that $\widehat{Q}_{i}^{1}=R^{2,1}\left(Q_{i}^{1}\right) \cup R^{2,2}\left(Q_{i}^{1}\right)$ and this a disjoint union. As before, $R^{2,1}\left(Q_{i}^{1}\right)$ misses $(1 / 2) \Delta_{Q_{i}^{1}}$ so that by the doubling property

$$
\mu\left(\widehat{Q}_{i}^{1}\right) \leq C_{M_{0}} \mu\left(\frac{1}{2} \Delta_{Q_{i}^{1}}\right) \leq C_{M_{0}} \mu\left(R^{2,2}\left(Q_{i}^{1}\right)\right)
$$

and then $\mu\left(R^{2,1}\right) \leq \theta \mu\left(\widehat{Q}_{i}^{1}\right)$. Next we set $R^{2,1}$ and $R^{2,2}$ as the union of the corresponding $R^{2,1}\left(Q_{i}^{1}\right)$ and $R^{2,2}\left(Q_{i}^{1}\right)$ with $Q_{i}^{1} \in \mathbb{D}^{1,1}$. Then,

$$
\begin{aligned}
\mu\left(R^{2,1}\right) & :=\mu\left(\bigcup_{Q_{i}^{1} \in \mathbb{D}^{1,1}} R^{2,1}\left(Q_{i}^{1}\right)\right)=\sum_{Q_{i}^{1} \in \mathbb{D}^{1,1}} \mu\left(R^{2,1}\left(Q_{i}^{1}\right)\right) \\
& \leq \theta \sum_{Q_{i}^{1} \in \mathbb{D}^{1,1}} \mu\left(\widehat{Q}_{i}^{1}\right)=\theta \mu\left(R^{1,1}\right) \leq \theta^{2} \mu(\bar{Q})
\end{aligned}
$$

A straightforward iteration argument now yields that $\mu(\partial Q)=0$. We omit the details.
Proposition 6.4. - Suppose that $\Omega$ is a 1 -sided NTA domain with ADR boundary. Fix $Q_{0} \in \mathbb{D}$, and let $\mathcal{F} \subset \mathbb{D}_{Q_{0}}$ be a disjoint family. Then for each $Q \in \mathbb{D}_{\mathcal{F}, Q_{0}}$, there is a radius $r_{Q} \approx K_{0} \ell(Q)$, and a point $A_{Q} \in \Omega_{\mathcal{G}, Q_{0}}$ which serves as a Corkscrew point simultaneously for $\Omega_{\mathcal{F}, Q_{0}}$, with respect to the surface ball $\Delta_{\star}\left(y_{Q}, r_{Q}\right)$, for some $y_{Q} \in \partial \Omega_{\mathcal{F}, Q_{0}}$, and for $\Omega$, with respect to each surface ball $\Delta\left(x, r_{Q}\right)$, for every $x \in Q$.

Proof. - Let $Q \in \mathbb{D}_{\mathcal{F}, Q_{0}}$. Recall that by construction, $W_{Q}$ is non-empty. It follows that there is an $I$ for which $\ell(I) \approx \ell(Q)$ and $\operatorname{dist}(Q, I) \leq C_{0} \ell(Q)$. Furthermore, $I \subset \Omega_{\mathcal{G}, Q_{0}}$, and $\operatorname{dist}\left(I, \partial \Omega_{\mathcal{G}, Q_{0}}\right) \leq \operatorname{dist}(I, \partial \Omega) \approx \ell(I)$. We let $A_{Q}$ denote the center of this particular $I$, so that

$$
\begin{equation*}
\ell(Q) \approx \operatorname{dist}\left(A_{Q}, \partial \Omega_{\mathcal{F}, Q_{0}}\right) \leq \operatorname{dist}\left(A_{Q}, Q\right) \lesssim C_{0} \ell(Q) \tag{6.5}
\end{equation*}
$$

Fix $y_{Q} \in \partial \Omega_{\mathscr{F}, Q_{0}}$ so that $\operatorname{dist}\left(A_{Q}, \partial \Omega_{\mathcal{F}, Q_{0}}\right)=\left|A_{Q}-y_{Q}\right|$. Then $A_{Q}$ is the promised simultaneous Corkscrew point, for $r_{Q} \approx K_{0} \ell(Q) \geq C_{0} \ell(Q)$.

Corollary 6.6. - The point $A_{Q_{0}}$ is a Corkscrew point with respect to $\Delta_{\star}\left(x, r_{Q_{0}}\right)$, for all $x \in \partial \Omega_{\mathcal{F}, Q_{0}}$, and for $\Delta\left(x, r_{Q_{0}}\right)$, for all $x \in Q_{0}$, with $r_{Q_{0}} \approx K_{0} \ell\left(Q_{0}\right)$.

The proof is almost immediate, since $\operatorname{diam}\left(\Omega_{\mathcal{F}, Q_{0}}\right) \lesssim K_{0} \ell\left(Q_{0}\right)$, and we omit it.

Proposition 6.7. - Suppose that $\Omega$ is a 1 -sided NTA domain with ADR boundary. Fix $Q_{0} \in \mathbb{D}$, and let $\mathcal{F} \subset \mathbb{D}_{Q_{0}}$ be a disjoint family. Then for each $Q_{j} \in \mathcal{F}$, there is an $n$-dimensional cube $P_{j} \subset \partial \Omega_{\mathcal{G}, Q_{0}}$, which is contained in a face of $I^{*}$, for some $I \in \mathcal{W}$, and which satisfies

$$
\begin{equation*}
\ell\left(P_{j}\right) \approx \operatorname{dist}\left(P_{j}, Q_{j}\right) \approx \operatorname{dist}\left(P_{j}, \partial \Omega\right) \approx \ell(I) \approx \ell\left(Q_{j}\right) \tag{6.8}
\end{equation*}
$$

where the uniform implicit constants are allowed to depend upon $K_{0}$.
Proof. - Fix $Q_{j} \in \mathscr{F}$. It follows from Lemma 5.9 (with $Q=Q_{j}$ ) and the Corkscrew condition that there is an $I_{1} \in \mathcal{W}$ with $I_{1} \subset \Omega \backslash \Omega_{\mathcal{G}, Q_{0}}, \quad \ell\left(I_{1}\right) \approx \ell\left(Q_{j}\right) / K_{0}$, and $\operatorname{dist}\left(I_{1}, Q_{j}\right) \lesssim \ell\left(Q_{j}\right) / K_{0}$. On the other hand, the dyadic parent $\widetilde{Q}_{j}$ of $Q_{j}$ belongs to $\mathbb{D}_{\mathcal{G}, Q_{0}}$, so there is an $I_{2} \in W_{\widetilde{Q}_{j}}^{*}$ with $I_{2} \subset \Omega_{\mathcal{F}, Q_{0}}, \ell\left(I_{2}\right) \approx \ell\left(Q_{j}\right)$, and $\operatorname{dist}\left(Q_{j}, I_{2}\right) \lesssim K_{0} \ell\left(Q_{j}\right)$. The Harnack chain (in $\Omega$ ) connecting the centers of $I_{1}$ and $I_{2}$, then passes through $\partial \Omega_{\mathcal{G}, Q_{0}}$, and maintains a distance to $\partial \Omega$ on the order of $\ell\left(Q_{j}\right)$. Consequently, there is an interface between some pair $I, J \in \mathcal{W}$, with $\operatorname{int}\left(I^{*}\right) \subset \Omega_{\mathcal{F}, Q_{0}}$ and $J \notin \mathcal{W}_{Q}^{*}$, for any $Q \in \mathbb{D}_{\mathcal{F}, Q_{0}}$ (so that $\tau J \subset \Omega \backslash \Omega_{\mathcal{G}, Q_{0}}$ for some $\tau \in(1 / 2,1)$; cf. (3.51)), and

$$
\operatorname{dist}\left(I, Q_{j}\right) \approx \operatorname{dist}\left(J, Q_{j}\right) \approx \ell(I) \approx \ell(J) \approx \ell\left(Q_{j}\right)
$$

(here, some of the implicit constants may depend upon $K_{0}$ ). Of course, the interface between $I$ and $J$ is precisely one face of the smaller of these two cubes. Therefore, if $\lambda$ is chosen small enough, then $\partial I^{*} \cap J$ contains an $n$-dimensional cube $P_{j}$ with the stated properties.

Remark 6.9. - It follows from the proof that if $P_{j} \cap P_{k}$ then $\ell\left(Q_{j}\right) \approx \ell\left(Q_{k}\right)$ since two adjacent Whitney cubes have comparable side length. Thus, $\operatorname{dist}\left(Q_{j}, Q_{k}\right) \lesssim \ell\left(Q_{j}\right)$ and therefore we have the bounded overlap property

$$
\sum_{j} 1_{P_{j}}(x) \leq C,
$$

with $C$ depending on the ADR constants.
For future reference, we note that, under the assumptions of Proposition 6.7, if $x_{j}^{\star}$ denotes the center of $P_{j}$, then for an appropriate choice of $r_{j} \approx K_{0} \ell\left(Q_{j}\right)$, we have $P_{j} \subset \Delta_{\star}\left(x_{j}^{\star}, r_{j}\right)$ and

$$
\begin{equation*}
\overline{T_{Q_{j}}} \subset B\left(x_{j}^{\star}, r_{j}\right), \tag{6.10}
\end{equation*}
$$

since $\operatorname{diam}\left(T_{Q_{j}}\right) \lesssim K_{0} \ell\left(Q_{j}\right)$. Moreover, given $Q \in \mathbb{D}_{\mathcal{G}, Q_{0}}$ and $r_{Q} \approx K_{0} \ell(Q)$ from Proposition 6.4, by choosing $\widehat{r}_{Q} \approx r_{Q}$ (with implicit constants depending on $K_{0}$ ) we may suppose that

$$
\begin{equation*}
Q \cup\left(\bigcup_{Q_{j} \in \mathscr{F}: Q_{j} \subset Q} B\left(x_{j}^{\star}, r_{j}\right)\right) \subset B\left(y_{Q}, \widehat{r}_{Q}\right) . \tag{6.11}
\end{equation*}
$$

Here, $y_{Q}$ is the center of $\Delta_{\star}\left(y_{Q}, r_{Q}\right) \subset \partial \Omega_{\mathcal{F}, Q_{0}}$, appearing in Proposition 6.4. We omit the routine geometric argument.

We conclude this preamble with the following.

Proposition 6.12. - Suppose that $\Omega$ is a 1 -sided NTA domain with ADR boundary. Fix $Q_{0} \in \mathbb{D}$, and let $\mathcal{F} \subset \mathbb{D}_{Q_{0}}$ be a disjoint family. For $Q_{j} \in \mathcal{F}$, let $B\left(x_{j}^{\star}, r_{j}\right)$ be the ball, concentric with $P_{j}$, satisfying (6.10). Then for each $Q \in \mathbb{D}_{\mathcal{G}, Q_{0}}$, there is a surface ball

$$
\Delta_{\star}^{Q}:=\Delta_{\star}\left(x_{Q}^{\star}, t_{Q}\right) \subset\left(Q \cap \partial \Omega_{\mathcal{F}, Q_{0}}\right) \cup\left(\bigcup_{Q_{j} \in \mathscr{F}: Q_{j} \subset Q}\left(B\left(x_{j}^{\star}, r_{j}\right) \cap \partial \Omega_{\mathscr{F}, Q_{0}}\right)\right),
$$

with $t_{Q} \approx \ell(Q), x_{Q}^{\star} \in \partial \Omega_{\mathcal{F}, Q_{0}}$, and $\operatorname{dist}\left(Q, \Delta_{\star}^{Q}\right) \lesssim \ell(Q)$, where the implicit constants may depend upon $K_{0}$.

Proof. - Suppose first that there is some $Q_{j_{0}} \subset Q$, for which $\ell\left(Q_{j_{0}}\right) \geq \ell(Q) / M$, where $M$ is a sufficiently large number to be chosen. We then set $\Delta_{\star}^{Q}=\Delta_{\star}\left(x_{j_{0}}^{\star}, \ell\left(P_{j_{0}}\right) / 2\right)$, a surface ball contained in the cube $P_{j_{0}}$ whose existence was established in Proposition 6.7.

Now suppose that $\ell\left(Q_{j}\right)<\ell(Q) / M$, for every $Q_{j} \subset Q$. By Lemma 3.55, there is a ball $B_{Q}^{\prime}=B\left(x_{Q}, s\right)$ with $s \approx \ell(Q)$ and $B_{Q}^{\prime} \cap \Omega \subset T_{Q} \subseteq T_{Q_{0}}$. In particular, $B_{Q}^{\prime}$ misses $\partial T_{Q_{0}} \backslash Q_{0}$. Moreover, $\Delta_{Q}^{\prime}:=B_{Q}^{\prime} \cap \partial \Omega \subset Q$. Consider those $Q_{j} \subset Q$ which meet $\Delta\left(x_{Q}, s /(4 \sqrt{M})\right)$. If there are no such $Q_{j}$, then we set $\Delta_{\star}^{Q}=\Delta\left(x_{Q}, s /(4 \sqrt{M})\right)$, which in this case is contained in $Q \cap \partial \Omega_{\mathscr{F}, Q_{0}}$ by Proposition 6.1. On the other hand, suppose that there is some $Q_{j_{0}} \subset Q$ which meets $\Delta\left(x_{Q}, s /(4 \sqrt{M})\right)$. Then for $M$ large enough, depending on $K_{0}$, we have $P_{j_{0}} \subset B\left(x_{Q}, s /(2 \sqrt{M})\right)$, and thus also

$$
\begin{equation*}
\Delta_{\star}^{Q}:=\Delta_{\star}\left(x_{j_{0}}^{\star}, s /(2 \sqrt{M})\right) \subset B\left(x_{Q}, s / \sqrt{M}\right) \subset B_{Q}^{\prime}, \tag{6.13}
\end{equation*}
$$

by the triangle inequality. Consequently, $\Delta_{\star}^{Q}$ misses $\partial T_{Q_{0}} \backslash Q_{0}$, and $\Delta_{\star}^{Q} \cap \partial \Omega \subset Q$, by the properties of $B_{Q}^{\prime}$. Moreover, we claim that

$$
\begin{equation*}
\partial \Omega_{\mathscr{F}, Q_{0}} \subset\left(\partial T_{Q_{0}} \backslash Q_{0}\right) \cup\left(\partial \Omega_{\mathscr{F}, Q_{0}} \cap Q_{0}\right) \cup\left(\bigcup_{Q_{j} \in \mathscr{F}}\left(\partial \Omega_{\mathscr{F}, Q_{0}} \cap \overline{T_{Q_{j}}}\right)\right) \tag{6.14}
\end{equation*}
$$

Let us defer for the moment the proof of this claim. Given (6.14), by (6.10) and properties of dyadic cubes, it is enough to verify that if $\Delta_{\star}^{Q}$ meets $\overline{T_{Q_{j}}}$, for some $Q_{j} \in \mathcal{F}$, then $Q_{j}$ meets $Q$. Suppose now that $\Delta_{\star}^{Q}$ meets $\overline{T_{Q_{j}}}$. By (6.13) and the definition of $\overline{T_{Q_{j}}}$, this means that there is a $Q^{\prime} \subseteq Q_{j}$, and an $I \in \mathcal{W}_{Q^{\prime}}^{*}$ such that $I^{*}$ meets $B\left(x_{Q}, s / \sqrt{M}\right)$. It follows that

$$
\ell\left(Q^{\prime}\right) \approx \ell(I) \approx \operatorname{dist}\left(I^{*}, \partial \Omega\right) \leq \operatorname{dist}\left(I^{*}, x_{Q}\right) \leq s / \sqrt{M}
$$

Since $\operatorname{dist}\left(I^{*}, Q^{\prime}\right) \lesssim K_{0} \ell\left(Q^{\prime}\right)$, by the triangle inequality we have

$$
\left|y-x_{Q}\right| \lesssim K_{0} s / \sqrt{M}<s, \quad \forall y \in Q^{\prime}
$$

if $M \gg\left(K_{0}\right)^{2}$; i.e., $Q^{\prime} \subset \Delta_{Q}^{\prime}:=\Delta\left(x_{Q}, s\right) \subset Q$, whence $Q_{j}$ meets $Q$, as desired.
Finally, we establish (6.14). Let $X \in \partial \Omega_{\mathcal{F}, Q_{0}}$. There are two cases.
CASE 1: $\delta(X)=0$. If $X \in Q_{0}$ we are done. Otherwise, since

$$
\partial \Omega_{\mathcal{F}, Q_{0}} \subset \overline{\Omega_{\mathcal{G}, Q_{0}}} \subset \overline{T_{Q_{0}}}
$$

it suffices to show that $X \notin T_{Q_{0}}$. But this is trivial, since $T_{Q_{0}} \subset \Omega$, and for $X \in \Omega$, we have that $\delta(X)>0$.

CASE 2: $\delta(X)>0$. If $X \in \overline{T_{Q_{j}}}$ for some $j$, we are done, so suppose that this never happens. As in Case 1, it is enough to show that $X \notin T_{Q_{0}}$, so suppose by way of contradiction that
$X \in T_{Q_{0}}$. Since $T_{Q_{0}}$ is open, this means that there is a small number $\varepsilon_{0} \ll \delta(X)$ such that the ball $B(X, \varepsilon) \subset T_{Q_{0}}$, whenever $\varepsilon \leq \varepsilon_{0}$. By definition of $T_{Q_{0}}$, and properties of Whitney cubes, there exist a uniformly bounded number of Whitney cubes, say, $I_{1}, \ldots, I_{M}$, such that

$$
B\left(X, \varepsilon_{0}\right) \subset \bigcup_{k=1}^{M} I_{k}^{*},
$$

and for each $k \in[1, M]$, there is a $Q^{k} \in \mathbb{D}_{Q_{0}}$ with $I_{k} \in W_{Q^{k}}^{*}$. It is possible that for a smaller $\varepsilon$, there may be a smaller collection of $I_{k}$ 's required to cover $B(X, \varepsilon)$, but these $I_{k}$ 's are of course always chosen from the original collection (i.e., the one for $\varepsilon_{0}$.) Observe that since $B(X, \varepsilon)$ is open, if $B(X, \varepsilon)$ meets $I_{k}^{*}$, then it meets $\operatorname{int}\left(I_{k}^{*}\right)$. For a given $\varepsilon$, we may assume that the covering collection is "minimal" in the sense that $B(X, \varepsilon)$ meets $\operatorname{int}\left(I_{k}^{*}\right)$ for each $k$, i.e., we remove those $I_{k}^{*}$ which do not meet $B(X, \varepsilon)$.

We claim that there must be some $\varepsilon>0$ and a corresponding "minimal" collection, with the property that each $Q^{k} \in \mathbb{D}_{\mathcal{G}, Q_{0}}$. Indeed, if not, then there is a sequence $\varepsilon_{i} \rightarrow 0$, and for each $i$, a $Q_{j}(i) \in \mathcal{F}$, and a $k(i) \in[1, M]$, such that $Q^{k(i)} \in \mathbb{D}_{Q_{j}(i)}$. Since there were only a bounded number of $I_{k}$ 's and thus also $Q^{k}$ 's, to start with, there must be some subsequence, again call it $\varepsilon_{i}$, such that $k(i)=$ constant. But this means that there is a fixed $Q_{j} \in \mathcal{F}$, and a sequence $\varepsilon_{i}$ such that $B\left(X, \varepsilon_{i}\right)$ meets $T_{Q_{j}}$, which contradicts our assumption that $X \notin \overline{T_{Q_{j}}}$ for any $Q_{j} \in \mathcal{F}$. This proves the claim.

We now choose $\varepsilon$ as in the claim, and observe that we then have

$$
B(X, \varepsilon) \subset \bigcup_{Q \in \mathbb{D}_{\mathcal{G}, Q_{0}}} \bigcup_{W_{Q}^{*}} I^{*}
$$

i.e., that $X$ is an interior point for the set $\bigcup_{Q \in \mathbb{D}_{\mathcal{F}, Q_{0}}} \bigcup_{W_{Q}^{*}} I^{*}$. But by definition, this means that $X \in \Omega_{\mathcal{G}, Q_{0}}$, and since the latter set is open, this contradicts that $X \in \partial \Omega_{\mathcal{F}, Q_{0}}$.

Lemma 6.15 (Dyadic sawtooth lemma for projections). - Suppose that $\Omega$ is a 1 -sided NTA domain with ADR boundary and that it also satisfies the qualitative exterior Corkscrew condition. Fix $Q_{0} \in \mathbb{D}$, let $\mathcal{F}=\left\{Q_{j}\right\} \subset \mathbb{D}_{Q_{0}}$ be a family of pairwise disjoint dyadic cubes and let $\mathscr{P}_{\mathcal{G}}$ be the corresponding projection operator. We write $\omega=\omega^{X_{0}}$ and $\omega_{\star}=\omega_{\star}^{X_{0}}$ to denote the respective harmonic measures for the domains $\Omega$ and $\Omega_{\mathcal{G}, Q_{0}}$, with fixed pole at the Corkscrew point $X_{0}:=A_{Q_{0}}$ whose existence was noted in Proposition 6.4 and Corollary 6.6. Let $\nu=\nu^{X_{0}}$ be the measure defined by

$$
\begin{equation*}
\nu(F)=\omega_{\star}\left(F \backslash\left(\bigcup_{\mathcal{F}} Q_{j}\right)\right)+\sum_{Q_{j} \in \mathcal{F}} \frac{\omega\left(F \cap Q_{j}\right)}{\omega\left(Q_{j}\right)} \omega_{\star}\left(P_{j}\right), \quad F \subset Q_{0}, \tag{6.16}
\end{equation*}
$$

where $P_{j}$ is the $n$-dimensional cube produced by Proposition 6.7. Then $\mathscr{P}_{g} \nu$ depends only on $\omega_{\star}$ and not on $\omega$. More precisely,

$$
\begin{equation*}
\mathscr{P}_{g} \nu(F)=\omega_{\star}\left(F \backslash\left(\bigcup_{\mathcal{G}} Q_{j}\right)\right)+\sum_{Q_{j} \in \mathcal{G}} \frac{\sigma\left(F \cap Q_{j}\right)}{\sigma\left(Q_{j}\right)} \omega_{\star}\left(P_{j}\right), \quad F \subset Q_{0} \tag{6.17}
\end{equation*}
$$

Moreover, there exists $\theta>0$ such that for all $Q \in \mathbb{D}_{Q_{0}}$ and $F \subset Q$, we have

$$
\begin{equation*}
\left(\frac{\mathscr{P}_{g} \omega(F)}{\mathcal{P}_{\mathscr{G}} \omega(Q)}\right)^{\theta} \lesssim \frac{\mathscr{P}_{g} \nu(F)}{\mathcal{P}_{\mathcal{G}} \nu(Q)} \lesssim \frac{\mathscr{P}_{g} \omega(F)}{\mathcal{P}_{\mathscr{G}} \omega(Q)} \tag{6.18}
\end{equation*}
$$

Proof. - We observe that (6.17) follows immediately from the definitions of $\mathscr{P}_{\mathscr{G}}$ and $\nu$, as the reader may readily verify. We omit the details.

Our first main task is to establish the right hand side inequality in (6.18). Let us fix $Q \in \mathbb{D}_{Q_{0}}, F \subset Q$.

Case 1: There exists $Q_{j} \in \mathcal{F}$ such that $Q \subset Q_{j}$. Note that by (6.17) we have

$$
\frac{\mathscr{P}_{g} \nu(F)}{\mathcal{P}_{g} \nu(Q)}=\frac{\frac{\sigma\left(F \cap Q_{j}\right)}{\sigma\left(Q_{j}\right)} \omega_{\star}\left(P_{j}\right)}{\frac{\sigma\left(Q \cap Q_{j}\right)}{\sigma\left(Q_{j}\right)} \omega_{\star}\left(P_{j}\right)}=\frac{\sigma(F)}{\sigma(Q)}=\frac{\frac{\sigma\left(F \cap Q_{j}\right)}{\sigma\left(Q_{j}\right)} \omega\left(Q_{j}\right)}{\frac{\sigma\left(Q \cap Q_{j}\right)}{\sigma\left(Q_{j}\right)} \omega\left(Q_{j}\right)}=\frac{\mathscr{P}_{g} \omega(F)}{\mathscr{P}_{g} \omega(Q)}
$$

Case 2: $Q$ is not contained in any $Q_{j} \in \mathscr{F}$ (i.e., $Q \in \mathbb{D}_{\mathcal{F}, Q_{0}}$ ). Notice that if $Q_{j} \in \mathscr{F}$ with $Q_{j} \cap Q \neq \varnothing$, then $Q_{j}$ is strictly contained in $Q$. Let us note also that $\omega_{\star}$ satisfies the doubling property, by Lemma 3.61, Lemma 3.62 and Corollary 3.36. Set $E_{0}=Q_{0} \backslash\left(\bigcup_{\mathcal{G}} Q_{j}\right)$. Using (6.17) we observe that

$$
\begin{align*}
\mathscr{P}_{\mathscr{G}} \nu(Q) & =\omega_{\star}\left(Q \cap E_{0}\right)+\sum_{Q_{j} \in \mathcal{F}, Q_{j} \subsetneq Q} \frac{\sigma\left(Q \cap Q_{j}\right)}{\sigma\left(Q_{j}\right)} \omega_{\star}\left(P_{j}\right) \\
& =\omega_{\star}\left(Q \cap E_{0}\right)+\sum_{Q_{j} \in \mathcal{F}, Q_{j} \subsetneq Q} \omega_{\star}\left(P_{j}\right)  \tag{6.19}\\
& \gtrsim \omega_{\star}\left(Q \cap E_{0}\right)+\sum_{Q_{j} \in \mathcal{F}, Q_{j} \subsetneq Q} \omega_{\star}\left(B\left(x_{j}^{\star}, r_{j}\right) \cap \partial \Omega_{\mathcal{G}, Q_{0}}\right) \\
& \geq \omega_{\star}\left(\Delta_{\star}^{Q}\right),
\end{align*}
$$

where in the third line we have used the doubling property of $\omega_{\star}$ (plus a subdivision and Harnack chain argument if $\ell\left(Q_{j}\right) \approx \ell\left(Q_{0}\right)$ ), and in the last line we have used Proposition 6.12, along with Propositions 6.1 and 6.3 and the doubling property to ignore the difference between $Q \backslash\left(\bigcup_{\mathscr{F}} Q_{j}\right)$ and $Q \cap \partial \Omega_{\mathcal{F}, Q_{0}}$.

Let $A_{Q}$ be as in Proposition 6.4. Then by Corollary 3.69 plus the doubling property and Harnack chain condition, and a differentiation argument, we have that for any Borel set $H \subset Q$,

$$
\begin{equation*}
\omega^{A_{Q}}(H) \approx \frac{\omega^{X_{0}}(H)}{\omega^{X_{0}}(Q)}=\frac{\omega(H)}{\omega(Q)} \tag{6.20}
\end{equation*}
$$

The same occurs for $\omega_{\star}$ and $\omega_{\star}^{A_{Q}}$ and for any $H_{\star} \subset \Delta_{\star}\left(y_{Q}, \hat{r}_{Q}\right)$, (see (6.11) and Proposition 6.4). More precisely,

$$
\begin{equation*}
\omega_{\star}^{A_{Q}}\left(H_{\star}\right) \approx \frac{\omega_{\star}^{X_{0}}\left(H_{\star}\right)}{\omega_{\star}^{X_{0}}\left(\Delta_{\star}\left(y_{Q}, \widehat{r}_{Q}\right)\right)}=\frac{\omega_{\star}\left(H_{\star}\right)}{\omega_{\star}\left(\Delta_{\star}\left(y_{Q}, \widehat{r}_{Q}\right)\right)} \approx \frac{\omega_{\star}\left(H_{\star}\right)}{\omega_{\star}\left(\Delta_{\star}^{Q}\right)}, \tag{6.21}
\end{equation*}
$$

where $\Delta_{\star}^{Q}$ is the surface ball in Proposition 6.12, and where the last step follows by the doubling property of $\omega_{\star}$, since $\operatorname{dist}\left(\Delta_{\star}^{Q}, \Delta_{\star}\left(y_{Q}, \widehat{r}_{Q}\right)\right) \lesssim \ell(Q)$, and the radius of each surface ball is comparable to $\ell(Q)$.

Using (6.19) and (6.21) (and (6.11)), we obtain

$$
\begin{align*}
\frac{\mathscr{P}_{\mathcal{G}} \nu(F)}{\mathscr{P}_{\mathcal{F}} \nu(Q)} & \lesssim \frac{\omega_{\star}\left(F \cap E_{0}\right)}{\omega_{\star}\left(\Delta_{\star}^{Q}\right)}+\sum_{Q_{j} \in \mathcal{F}, Q_{j} \subsetneq Q} \frac{\sigma\left(F \cap Q_{j}\right)}{\sigma\left(Q_{j}\right)} \frac{\omega_{\star}\left(P_{j}\right)}{\omega_{\star}\left(\Delta_{\star}^{Q}\right)}  \tag{6.22}\\
& \approx \omega_{\star}^{A_{Q}}\left(F \cap E_{0}\right)+\sum_{Q_{j} \in \mathcal{F}, Q_{j} \subsetneq Q} \frac{\sigma\left(F \cap Q_{j}\right)}{\sigma\left(Q_{j}\right)} \omega_{\star}^{A_{Q}}\left(P_{j}\right)
\end{align*}
$$

We claim that the following estimates hold:

$$
\begin{equation*}
\omega_{\star}^{A_{Q}}\left(F \cap E_{0}\right) \lesssim \omega^{A_{Q}}\left(F \cap E_{0}\right), \quad \omega_{\star}^{A_{Q}}\left(P_{j}\right) \lesssim \omega^{A_{Q}}\left(Q_{j}\right) \tag{6.23}
\end{equation*}
$$

Indeed, the first estimate follows immediately from the maximum principle, since $\Omega_{\mathcal{F}, Q_{0}} \subset \Omega$, and $E_{0} \subset \partial \Omega \cap \partial \Omega_{\mathcal{F}, Q_{0}}$, by Proposition 6.1. To prove the second estimate, we observe that, again by the maximum principle, it suffices to show that $\omega^{X}\left(Q_{j}\right) \gtrsim 1$, for $X \in P_{j}$. But the latter bound follows immediately from (3.8) with $\Delta=\Delta_{Q_{j}}$ (cf. (1.16)-(1.17)), the Harnack chain condition and (6.8).

The bounds in (6.22), (6.23) and (6.20) imply

$$
\begin{aligned}
\frac{\mathscr{P}_{\mathscr{G}} \nu(F)}{\mathscr{P}_{\mathscr{F}} \nu(Q)} & \lesssim \omega^{A_{Q}}\left(F \cap E_{0}\right)+\sum_{Q_{j} \in \mathscr{F}, Q_{j} \subsetneq Q} \frac{\sigma\left(F \cap Q_{j}\right)}{\sigma\left(Q_{j}\right)} \omega^{A_{Q}}\left(Q_{j}\right) \\
& \approx \frac{\omega\left(F \cap E_{0}\right)}{\omega(Q)}+\sum_{Q_{j} \in \mathcal{F}, Q_{j} \subsetneq Q} \frac{\sigma\left(F \cap Q_{j}\right)}{\sigma\left(Q_{j}\right)} \frac{\omega\left(Q_{j}\right)}{\omega(Q)}=\frac{\mathscr{P}_{\mathscr{G}} \omega(F)}{\omega(Q)}=\frac{\mathscr{P}_{\mathscr{G}} \omega(F)}{\mathscr{P}_{\mathscr{G}} \omega(Q)}
\end{aligned}
$$

where in the last equality we have used that $\mathscr{P}_{\mathscr{F}} \omega(Q)=\omega(Q)$. Thus, we have established the right hand inequality in (6.18). We may now obtain the left hand side of (6.18) by a direct application of Lemma B. 7 (see Appendix B below), using the fact that $\mathscr{P}_{\mathcal{F}} \omega$ and $\mathcal{P}_{\mathscr{F}} \nu$ are dyadically doubling by Lemmas B. 1 and B.2.

## 7. A discrete Corona decomposition

In this section we present a discretized version of the stopping time decomposition of a Carleson region appearing in [12, 3], [4, 27] and [26] (cf. [11, 36], [25]). We suppose that $\left\{\alpha_{Q}\right\}_{Q \in \mathbb{D}}$ is a sequence of non-negative numbers indexed on the dyadic "cubes", and for any collection $\mathbb{D}^{\prime} \subset \mathbb{D}$, we define

$$
\mathfrak{m}\left(\mathbb{D}^{\prime}\right):=\sum_{Q \in \mathbb{D}^{\prime}} \alpha_{Q}
$$

For a fixed $Q_{0} \in \mathbb{D}$, we say that $\mathfrak{m}$ is a "Carleson measure" on $\mathbb{D}_{Q_{0}}$ (with respect to $\sigma$ ), and we write $\mathfrak{m} \in \mathscr{C}\left(Q_{0}\right)$, if

$$
\|\mathfrak{m}\|_{\mathscr{C}\left(Q_{0}\right)}:=\sup _{Q \in \mathbb{D}_{Q_{0}}} \frac{\mathfrak{m}\left(\mathbb{D}_{Q}\right)}{\sigma(Q)}<\infty
$$

We also write

$$
\begin{equation*}
\|\mathfrak{m}\|_{\mathscr{C}}:=\sup _{Q \in \mathbb{D}} \frac{\mathfrak{m}\left(\mathbb{D}_{Q}\right)}{\sigma(Q)}<\infty \tag{7.1}
\end{equation*}
$$



Figure 1. "Tree-graph" with vertex $Q$ and its random walk
to denote the "global" Carleson norm on $\mathbb{D}$. We furthermore set $\mathbb{D}_{Q}^{\text {short }}:=\mathbb{D}_{Q} \backslash\{Q\}$, and given a family $\mathcal{F} \subset \mathbb{D}$ of pairwise disjoint cubes, we define the "restriction of $\mathfrak{m}$ to the sawtooth $\mathbb{D}_{\mathcal{G}} "$ by

$$
\mathfrak{m}_{\mathcal{G}}\left(\mathbb{D}^{\prime}\right):=\mathfrak{m}\left(\mathbb{D}^{\prime} \cap \mathbb{D}_{\mathcal{G}}\right)=\sum_{Q \in \mathbb{D}^{\prime} \backslash\left(\cup_{g} \mathbb{D}_{Q_{j}}\right)} \alpha_{Q} .
$$

We fix $Q_{0} \in \mathbb{D}$, and construct a "tree-graph" with a vertex for each $Q \in \mathbb{D}_{Q_{0}}$, and with edges connecting a given $Q$ to each of its dyadic "children" (these are the subcubes of $Q$ which lie in the very next dyadic generation $\mathbb{D}_{k(Q)+1}$ ). We consider a random walk along the graph, in which it is permitted to move only to the descendant generation, but not to the ancestral generation (nor to any other cube in the same generation), and we suppose that from a given $Q \in \mathbb{D}_{Q_{0}}$, there is an equal probability of arriving at any of its children. We set $P(Q, Q)=1$, and in general for $Q^{\prime} \subseteq Q \in \mathbb{D}_{Q_{0}}$, we denote by $P\left(Q, Q^{\prime}\right)$ the probability that such a random walk beginning at $Q$ arrives at $Q^{\prime}$ (thus also if $Q$ is strictly contained in $Q^{\prime}$, or if $Q$ and $Q^{\prime}$ are disjoint, we have $P\left(Q, Q^{\prime}\right)=0$ ). See Figure 1 .

Lemma 7.2. - Suppose that $\partial \Omega$ is $A D R$. Fix $Q_{0} \in \mathbb{D}$ and $\mathfrak{m}$ as above. Let $a \geq 0$ and $b>0$, and suppose that $\mathfrak{m}\left(\mathbb{D}_{Q_{0}}\right) \leq(a+b) \sigma\left(Q_{0}\right)$. Then there is a family $\mathcal{F}=\left\{Q_{j}\right\} \subset \mathbb{D}_{Q_{0}}$ of pairwise disjoint cubes, and a constant $C$ depending only on dimension and the ADR constants such that

$$
\begin{align*}
\left\|\mathfrak{m}_{\mathcal{G}}\right\|_{\mathcal{E}\left(Q_{0}\right)} & \leq C b,  \tag{7.3}\\
\sigma(B) & \leq \frac{a+b}{a+2 b} \sigma\left(Q_{0}\right), \tag{7.4}
\end{align*}
$$

where $B$ is the union of those $Q_{j} \in \mathcal{F}$ such that $\mathfrak{m}\left(\mathbb{D}_{Q_{j}}^{\text {short }}\right)>a \sigma\left(Q_{j}\right)$.
Remark 7.5. - In the proof of this result, the only feature of $\sigma$ that we shall use, is that it is a non-negative Borel measure satisfying the "dyadically doubling property on $Q_{0}$ " (by this
we mean that there is a uniform constant $c_{\sigma}$ such that $\sigma(\widetilde{Q}) \leq c_{\sigma} \sigma(Q)$, whenever $\widetilde{Q} \in \mathbb{D}_{Q_{0}}$ is the dyadic parent of $Q$ ). Notice that this property follows at once for our measure $\sigma=\left.H^{n}\right|_{\partial \Omega}$ by the ADR property. Therefore, Lemma 7.2 admits an extension in which $\sigma$ can be any nonnegative dyadically doubling Borel measure on $Q_{0}$.

Proof. - We note that $\mathfrak{m}\left(\mathbb{D}_{Q_{0}}\right)=\mathfrak{m}\left(\mathbb{D}_{Q_{0}}^{\text {short }}\right)+\alpha_{Q_{0}}$. Thus, if $\alpha_{Q_{0}}>b \sigma\left(Q_{0}\right)$, the result is trivial: in this case $\mathfrak{m}\left(\mathbb{D}_{Q_{0}}^{\text {short }}\right) \leq a \sigma\left(Q_{0}\right)$, so we may set $\mathcal{F}=\left\{Q_{0}\right\}$, and $B=\varnothing$.

Suppose now that $\alpha_{Q_{0}} \leq b \sigma\left(Q_{0}\right)$. For $Q^{\prime} \in \mathbb{D}_{Q_{0}}$, we set

$$
\beta\left(Q^{\prime}\right):=\sum_{Q: Q^{\prime} \subseteq Q \subseteq Q_{0}} P\left(Q, Q^{\prime}\right) \alpha_{Q}
$$

In particular, $\beta\left(Q_{0}\right)=\alpha_{Q_{0}} \leq b \sigma\left(Q_{0}\right)$. We now perform a standard stopping time argument to select the collection $\mathcal{F}=\left\{Q_{j}\right\}$, comprised of the subcubes of $Q_{0}$ which are maximal with respect to the property that

$$
\begin{equation*}
\beta\left(Q_{j}\right)>2 b \sigma\left(Q_{j}\right) \tag{7.6}
\end{equation*}
$$

If $\mathcal{F}$ is empty, we simply have that $\mathbb{D}_{\mathcal{G}}=\mathbb{D}, \mathfrak{m}_{\mathcal{F}}=\mathfrak{m}$ and $B=\varnothing$.
We now verify that $\mathcal{F}$ satisfies the desired properties. We start by proving (7.4). To this end let us record some useful facts. We first note that, given a fixed $Q \in \mathbb{D}_{Q_{0}}$,

$$
\begin{equation*}
\sum_{Q_{j} \in \mathscr{F}} P\left(Q, Q_{j}\right)=\sum_{Q_{j} \in \mathscr{F}: Q_{j} \subseteq Q} P\left(Q, Q_{j}\right) \leq 1 \tag{7.7}
\end{equation*}
$$

since the cubes in $\mathcal{F}$, and therefore also the events in the sum, are disjoint. Next, we note that since $P\left(Q_{j}, Q_{j}\right)=1$,

$$
\begin{equation*}
\mathfrak{m}\left(\mathbb{D}_{Q_{j}}^{\text {short }}\right)+\beta\left(Q_{j}\right)=\mathfrak{m}\left(\mathbb{D}_{Q_{j}}\right)+\sum_{Q: Q_{j \subsetneq Q \subseteq Q_{0}}} P\left(Q, Q_{j}\right) \alpha_{Q} \tag{7.8}
\end{equation*}
$$

where the last sum runs over those $Q \in \mathbb{D}_{Q_{0}}$ that strictly contain $Q_{j}$. Consequently,

$$
\begin{align*}
\sum_{Q_{j} \in \mathscr{F}}\left(\mathfrak{m}\left(\mathbb{D}_{Q_{j}}^{\text {short }}\right)\right. & \left.+\beta\left(Q_{j}\right)\right)  \tag{7.9}\\
& =\sum_{Q_{j} \in \mathscr{F}} \mathfrak{m}\left(\mathbb{D}_{Q_{j}}\right)+\sum_{Q \in \mathbb{D}_{\mathcal{G}, Q_{0}}} \alpha_{Q} \sum_{Q_{j} \in \mathscr{F}: Q_{j} \subsetneq Q} P\left(Q, Q_{j}\right) \leq \mathfrak{m}\left(\mathbb{D}_{Q_{0}}\right)
\end{align*}
$$

by (7.7) and the definition of $\mathbb{D}_{\mathcal{F}, Q_{0}}$ (cf. (3.41)).
We now set $\mathcal{F}_{\text {bad }}:=\left\{Q_{j} \in \mathcal{F}: \mathfrak{m}\left(\mathbb{D}_{Q_{j}}^{\text {short }}\right)>a \sigma\left(Q_{j}\right)\right\}$. Then by definition of $B$ and the stopping time construction,

$$
\begin{aligned}
(a+2 b) \sigma(B) & =(a+2 b) \sum_{Q_{j} \in \mathscr{F}_{\text {bad }}} \sigma\left(Q_{j}\right) \\
& \leq \sum_{Q_{j} \in \mathscr{F}_{\text {bad }}}\left(\mathfrak{m}\left(\mathbb{D}_{Q_{j}}^{\text {short }}\right)+\beta\left(Q_{j}\right)\right) \leq \mathfrak{m}\left(\mathbb{D}_{Q_{0}}\right) \leq(a+b) \sigma\left(Q_{0}\right)
\end{aligned}
$$

where in the last line we have used (7.9) and our hypothesis. Estimate (7.4) follows.
We now turn to the proof of (7.3). Let us fix $Q \in \mathbb{D}_{Q_{0}}$. We consider

$$
\mathfrak{m}_{\mathcal{G}}\left(\mathbb{D}_{Q}\right)=\mathfrak{m}\left(\mathbb{D}_{\mathcal{F}, Q}\right)=\sum_{Q^{\prime} \in \mathbb{D}_{\mathcal{F}, Q}} \alpha_{Q^{\prime}}=\lim _{N \rightarrow \infty} \sum_{Q^{\prime} \in \mathbb{D}_{\mathcal{F}_{N}, Q}} \alpha_{Q^{\prime}}=\lim _{N \rightarrow \infty} \mathfrak{m}\left(\mathbb{D}_{\mathcal{F}_{N}, Q}\right)
$$

where $\mathcal{F}_{N}:=\mathscr{F}\left(2^{-N-1}\right)$ is derived from $\mathcal{F}$ as in the discussion at the beginning of Section 4; i.e., $\mathcal{F}_{N}=\left\{Q_{i}^{N}\right\}$ is the collection of maximal cubes of

$$
\mathcal{F} \cup\left\{Q^{\prime} \in \mathbb{D}_{Q}: \ell\left(Q^{\prime}\right) \leq 2^{-N-1}\right\}
$$

Thus,

$$
\mathbb{D}_{\mathcal{G}_{N}, Q}=\left\{Q^{\prime} \in \mathbb{D}_{\mathcal{F}, Q}: \ell\left(Q^{\prime}\right) \geq 2^{-N}\right\}, \quad N \geq k(Q)
$$

It is therefore enough to establish the bound

$$
\begin{equation*}
\mathfrak{m}\left(\mathbb{D}_{\mathcal{G}_{N}, Q}\right) \leq C b \sigma(Q) \tag{7.10}
\end{equation*}
$$

uniformly in $N$. To this end, we observe that equality holds in (7.7), for a given cube $Q$ and pairwise disjoint family $\mathcal{F}$, if $Q$ is covered by a union of cubes in $\mathcal{F}$. Since this is the case for the family $\mathcal{F}_{N}$ and for every $Q^{\prime} \in \mathbb{D}_{\mathcal{G}_{N}, Q}$, we have

$$
\begin{aligned}
\mathfrak{m}\left(\mathbb{D}_{\mathcal{F}_{N}, Q}\right)= & \sum_{Q^{\prime} \in \mathbb{D}_{\mathcal{G}_{N}, Q}} \alpha_{Q^{\prime}} \sum_{Q_{i}^{N} \in \mathcal{F}_{N}: Q_{i}^{N} \subseteq Q^{\prime}} P\left(Q^{\prime}, Q_{i}^{N}\right) \\
& =\sum_{Q_{i}^{N} \in \mathcal{F}_{N} \cap \mathbb{D}_{Q}} \sum_{Q^{\prime}: Q_{i}^{N} \subset Q^{\prime} \in \mathbb{D}_{\mathcal{G}_{N}, Q}} P\left(Q^{\prime}, Q_{i}^{N}\right) \alpha_{Q^{\prime}}=\Sigma_{1}+\Sigma_{2},
\end{aligned}
$$

where in $\Sigma_{1}$ the first sum runs over those $Q_{i}^{N} \in \mathcal{F}_{N} \cap \mathbb{D}_{Q}$ which are equal to some $Q_{j} \in \mathcal{F}$ (i.e., $Q_{i}^{N} \in \mathcal{F}_{N} \cap \mathcal{F} \cap \mathbb{D}_{Q}$ ), while in $\Sigma_{2}$ the first sum runs over the remaining cubes in $\mathcal{F}_{N} \cap \mathbb{D}_{Q}$ (i.e., over $Q_{i}^{N} \in\left(\mathcal{F}_{N} \backslash \mathcal{F}\right) \cap \mathbb{D}_{Q}$, equivalently those $Q_{i}^{N}$ which are not contained in any $Q_{j} \in \mathscr{F}$ ). We then have

$$
\begin{aligned}
\Sigma_{2} & =\sum_{Q_{i}^{N} \in\left(\mathscr{F}_{N} \backslash \mathscr{F}\right) \cap \mathbb{D}_{Q}}\left(\sum_{Q^{\prime}: Q_{i}^{N} \subset Q^{\prime} \in \mathbb{D}_{g_{N}, Q}} P\left(Q^{\prime}, Q_{i}^{N}\right) \alpha_{Q^{\prime}}\right) \\
& \leq \sum_{Q_{i}^{N} \in\left(\mathscr{F}_{N} \backslash \mathscr{F}\right) \cap \mathbb{D}_{Q}} \beta\left(Q_{i}^{N}\right) \leq 2 b \sum_{Q_{i}^{N} \in\left(\mathscr{F}_{N} \backslash \mathscr{F}\right) \cap \mathbb{D}_{Q}} \sigma\left(Q_{i}^{N}\right) \leq 2 b \sigma(Q)
\end{aligned}
$$

by the stopping time construction of $\mathcal{F}$, since $Q_{i}^{N}$ is not contained in any $Q_{j} \in \mathcal{F}$, and $Q_{i}^{N} \in \mathbb{D}_{Q}$.

We now consider $\Sigma_{1}$. We first note that no $Q^{\prime}$ appearing in the sum can be contained in any $Q_{j} \in \mathcal{F}$, since $\mathbb{D}_{\mathcal{F}_{N}, Q} \subset \mathbb{D}_{\mathcal{F}, Q}$. Therefore, if some $Q_{j} \in \mathcal{F}$ is contained in any such $Q^{\prime}$, then so is its dyadic parent $\widetilde{Q}_{j}$. Moreover,

$$
P\left(Q^{\prime}, Q_{j}\right) \leq P\left(Q^{\prime}, \widetilde{Q}_{j}\right), \quad \forall Q_{j} \in \mathcal{F}
$$

We then have that, by definition,

$$
\begin{aligned}
\Sigma_{1} & \leq \sum_{Q_{j} \in \mathcal{Y}} \sum_{Q^{\prime}: \widetilde{Q}_{j} \subseteq Q^{\prime} \subseteq Q} P\left(Q^{\prime}, \widetilde{Q}_{j}\right) \alpha_{Q^{\prime}} \leq \sum_{Q_{j} \in \mathcal{F}: Q_{j} \subset Q} \beta\left(\widetilde{Q}_{j}\right) \\
& \leq 2 b \sum_{Q_{j} \in \mathcal{F}: Q_{j} \subset Q} \sigma\left(\widetilde{Q}_{j}\right) \leq C b \sigma(Q),
\end{aligned}
$$

where the next-to-last inequality holds because the cubes $Q_{j}$ are maximal with respect to the property (7.6), and the last one holds by the dyadic doubling property of $\sigma$ (see Remark 7.5), and the pairwise disjointness of the cubes in $\mathcal{F}$.

## 8. Proofs of Theorems $\mathbf{1 . 2 6}$ and 1.27

### 8.1. Relating geometric and discrete Carleson measures

We recall that the UR property may be characterized in terms of the Carleson measure estimate (1.10), which we shall invoke with $E=\partial \Omega$. We also remind the reader that we may assume that for every Whitney cube $I \in \mathcal{W}$, we have $\operatorname{dist}(4 I, \partial \Omega) \approx \ell(I)$ (cf. (3.42)). In this case, the "fattened" Whitney cubes $4 I$ have bounded overlaps. From this fact, properties of Whitney cubes, and the ADR property, it follows that the fattened Whitney regions $U_{Q}^{\mathrm{fat}}=\bigcup_{W_{Q}^{*}} 4 I$ (cf. (3.43)-(3.49) and (4.1)) also have the bounded overlap property:

$$
\begin{equation*}
\sum_{Q \in \mathbb{D}} 1_{U_{Q}^{\mathrm{fat}}}(X) \leq C \tag{8.1}
\end{equation*}
$$

We now set

$$
\begin{equation*}
\alpha_{Q}:=\iint_{U_{Q}^{\mathrm{fat}}}\left|\nabla^{2} \phi 1(X)\right|^{2} \delta(X) d X \tag{8.2}
\end{equation*}
$$

and for any sub-collection $\mathbb{D}^{\prime} \subset \mathbb{D}$, we define

$$
\begin{equation*}
\mathfrak{m}\left(\mathbb{D}^{\prime}\right):=\sum_{Q \in \mathbb{D}^{\prime}} \alpha_{Q} \tag{8.3}
\end{equation*}
$$

as in the Section 7. By (8.1), for every pairwise disjoint family $\mathcal{F} \subset \mathbb{D}$, and every $Q \in \mathbb{D}_{\mathcal{F}}$, we have

$$
\begin{equation*}
\mathfrak{m}_{\mathcal{G}}\left(\mathbb{D}_{Q}\right)=\mathfrak{m}\left(\mathbb{D}_{\mathcal{F}, Q}\right) \approx \iint_{\Omega_{\mathcal{F}, Q}^{\mathrm{fat}}}\left|\nabla^{2} \phi 1(X)\right|^{2} \delta(X) d X \tag{8.4}
\end{equation*}
$$

where $\Omega_{\mathcal{F}, Q}^{\mathrm{fat}}$ is defined in (4.3) (we have used in (8.4) the rather trivial fact that $\partial \Omega_{\mathcal{F}, Q}^{\mathrm{fat}}$ has $(n+1)$-dimensional Lebesgue measure 0 ). In particular, taking $\mathcal{F}=\varnothing$, in which case $\mathbb{D}_{\mathcal{F}, Q}=\mathbb{D}_{Q}$, and $\Omega_{\mathcal{G}, Q}^{\mathrm{fat}}=T_{Q}^{\text {fat }}$ (cf. (4.2)), we obtain from (5.12) that $\mathfrak{m}$ inherits the Carleson measure property (7.1) from (1.10), and that the Carleson norm $\|\mathfrak{m}\|_{\mathscr{C}}$ depends only on dimension and the various ADR, UR, Corkscrew and Harnack chain constants for $\Omega$ (including $K_{0}$ ).

### 8.2. Proof of Theorem 1.26 with "qualitative assumptions"

In this subsection, we present the proof of Theorem 1.26, in the special case that the qualitative exterior Corkscrew condition holds in $\Omega$ (and therefore also in its sawtooths and Carleson boxes). As we observed in Section 3, this qualitative hypothesis (along with our standard quantitative assumptions), were enough to imply the doubling condition for the harmonic measure for $\Omega$ and the sawtooth regions, and also to allow us to obtain the "Dyadic Sawtooth" Lemma 6.15. We shall remove the qualitative assumptions, and also give the proof of Theorem 1.27, in Subsection 8.3.

We shall use the method of "extrapolation of Carleson measures", based on ideas originating in [12] and [36] (cf. [25, 3], [4, 27]). In more precise detail, we follow our related work in the Euclidean setting [27]. In the sequel, we say that a measure $\mu$ is "dyadically doubling on $Q_{0}$ " if there is a uniform constant $c_{\mu}$ such that $\mu(\widetilde{Q}) \leq c_{\mu} \mu(Q)$, whenever $\widetilde{Q} \in \mathbb{D}_{Q_{0}}$ is the dyadic parent of $Q$.

Lemma 8.5. - We fix $Q_{0} \in \mathbb{D}$. Let $\sigma$ and $\omega$ be a pair of non-negative, dyadically doubling Borel measures on $Q_{0}$, and let $\mathfrak{m}$ be a discrete Carleson measure with respect to $\sigma$ (cf. Section 7) with

$$
\|\mathfrak{m}\|_{\mathscr{C}\left(Q_{0}\right)} \leq M_{0}
$$

Suppose that there is $a \gamma>0$ such that for every $Q \in \mathbb{D}_{Q_{0}}$ and every family of pairwise disjoint dyadic subcubes $\mathcal{F}=\left\{Q_{j}\right\} \subset \mathbb{D}_{Q}$ verifying

$$
\begin{equation*}
\left\|\mathfrak{m}_{\mathcal{G}}\right\|_{\mathscr{C}(Q)} \leq \gamma \tag{8.6}
\end{equation*}
$$

we have that $\mathscr{P}_{\mathscr{G}} \omega$ satisfies the following property:

$$
\begin{equation*}
\forall \varepsilon \in(0,1), \exists C_{\varepsilon}>1 \text { such that }\left(F \subset Q, \frac{\sigma(F)}{\sigma(Q)} \geq \varepsilon \quad \Longrightarrow \quad \frac{\mathscr{P}_{\mathscr{G}} \omega(F)}{\mathscr{P}_{\mathscr{G}} \omega(Q)} \geq \frac{1}{C_{\varepsilon}}\right) \tag{8.7}
\end{equation*}
$$

Then, there exist $\eta_{0} \in(0,1)$ and $C_{0}<\infty$ such that, for every $Q \in \mathbb{D}_{Q_{0}}$,

$$
\begin{equation*}
F \subset Q, \quad \frac{\sigma(F)}{\sigma(Q)} \geq 1-\eta_{0} \quad \Longrightarrow \quad \frac{\omega(F)}{\omega(Q)} \geq \frac{1}{C_{0}} \tag{8.8}
\end{equation*}
$$

I.e., $\omega \in A_{\infty}^{\text {dyadic }}\left(Q_{0}\right)$.

REmARK 8.9. - Notice that in the statement of the lemma, $\sigma$ and $\omega$ are allowed to be any pair of non-negative, dyadically doubling Borel measures on $Q_{0}$, and that $\sigma$ plays the role of underlying measure. Therefore, $\sigma$ appears implicitly in the Carleson conditions, $\mathscr{P}_{\mathcal{F}}$ and in the definition of the class $A_{\infty}^{\text {dyadic }}\left(Q_{0}\right)$. In the present paper, we shall apply this result in the special case that $\sigma=\left.H^{n}\right|_{\partial \Omega}$, and $\omega=\omega^{X_{Q_{0}}}$, the harmonic measure with pole at the Corkscrew point $X_{Q_{0}}$.

Remark 8.10. - It is known that (8.8), for every $Q \in \mathbb{D}_{Q_{0}}$, self-improves to (1.21), but this fact may also be gleaned from Remark B. 10 below.

Remark 8.11. - The key hypothesis of the lemma, and the main point that must be verified in applications, is that (8.6) implies (8.7), for sufficiently small $\gamma$.

In the remainder of this subsection, we shall use Lemma 8.5 to prove Theorem 1.26, assuming the extra qualitative exterior Corkscrew condition. The qualitative hypothesis will then be removed in the next subsection. We defer the proof of the lemma to Section 9.

Proof of Theorem 1.26 with qualitative hypothesis. - To begin, we let $\sigma=\left.H^{n}\right|_{\partial \Omega}$, which is dyadically doubling by the ADR property. Let us fix $Q_{0} \in \mathbb{D}$, and set $\omega:=\omega^{X_{Q_{0}}}$, where as usual $X_{Q_{0}}$ is a Corkscrew point relative to $Q_{0}$. Given the qualitative hypothesis, it holds in particular that $\omega$ is a doubling measure (cf. Corollary 3.36), and therefore also dyadically doubling, on $Q_{0}$ (to obtain dyadic doubling when $\ell(Q) \approx \ell\left(Q_{0}\right)$, we may need to invoke the Harnack chain condition); moreover, the doubling constants depend only upon the constants in the quantitative hypotheses of Theorem 1.26 (i.e., dimension, UR, ADR, Harnack chain and Corkscrew, including the constant $K_{0}$ which ultimately depended only upon the other stated parameters). We define $\mathfrak{m}$ as in (8.3), with $\alpha_{Q}$ as in (8.2). As observed in the previous subsection, this $\mathfrak{m}$ inherits the discrete Carleson measure property (7.1) from (1.10). Therefore, once we have verified that (8.6), with $\gamma$ small enough, implies (8.7), we may then conclude from Lemma 8.5 that $\omega=\omega^{X_{Q_{0}}} \in A_{\infty}^{\text {dyadic }}\left(Q_{0}\right)$, for every $Q_{0} \in \mathbb{D}$, and thus by the Harnack chain condition that $\omega^{X_{Q_{0}}} \in A_{\infty}^{\text {dyadic }}\left(Q_{1}\right)$, for every $Q_{1}$ of the same generation
$4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - No 3
as $Q_{0}$ such that $\operatorname{dist}\left(Q_{0}, Q_{1}\right) \leq 100 \operatorname{diam}\left(Q_{0}\right)$. Since this is true for every $Q_{0} \in \mathbb{D}$, and since $\omega^{X_{Q}}$ is concentrically doubling, we may conclude that $\omega^{X_{\Delta}} \in A_{\infty}(\Delta)$, for every surface ball $\Delta=\Delta(x, r), x \in \partial \Omega$ and $r \leq \operatorname{diam}(\partial \Omega)$, with $A_{\infty}$ constants uniformly controlled, and depending only upon dimension and the UR, ADR, Harnack chain and Corkscrew constants. We reached this conclusion by imposing the extra qualitative exterior Corkscrew condition, but as our estimates do not depend quantitatively on that hypothesis, we shall be able, in Subsection 8.3, to remove it by an approximation argument, but at the the loss of the doubling property of $\omega$.

To complete our task in this subsection, it now remains only to verify that (8.6), with $\gamma$ small enough, implies (8.7). To this end, we fix $Q \in \mathbb{D}_{Q_{0}}$ and a pairwise disjoint family $\mathcal{F} \subset \mathbb{D}_{Q}$, and we suppose that (8.6) holds for some small $\gamma$ to be chosen momentarily. By (8.4), we deduce that

$$
\begin{equation*}
\sup _{Q^{\prime} \in \mathbb{D}_{Q}} \frac{1}{\sigma\left(Q^{\prime}\right)} \iint_{\Omega_{G, Q^{\prime}}^{\mathrm{fat}}}\left|\nabla^{2} \delta 1(X)\right|^{2} \delta(X) d X \leq C \gamma . \tag{8.12}
\end{equation*}
$$

Consequently, if $\gamma$ is small enough, depending only upon the allowable quantitative parameters, we may apply Corollary 5.17, with $Q$ in place of $Q_{0}$, to obtain, for every surface ball $\Delta_{\star}=B \cap \partial \Omega_{\mathcal{G}, Q}$, with $B=B(x, r), x \in \partial \Omega_{\mathcal{G}, Q}$, and $r \leq \operatorname{diam}(Q)$, that $\omega_{\star}^{X \Delta_{\star}} \in A_{\infty}\left(\Delta_{\star}\right)$, where $\omega_{\star}^{X}$ denotes harmonic measure for $\Omega_{\mathcal{G}, Q}$. Moreover, the $A_{\infty}$ constants are uniformly controlled by the stated parameters. By the Harnack chain condition, we obtain that $\omega_{\star}^{A_{Q}} \in A_{\infty}\left(\partial \Omega_{\mathcal{G}, Q}\right)$ (meaning that we view $\partial \Omega_{\mathcal{F}, Q}$ itself as a surface ball $\Delta_{\star}^{Q}$ of radius $r\left(\Delta_{\star}^{Q}\right) \approx K_{0} \ell(Q)$, and that $\left.\omega_{\star}^{A_{Q}} \in A_{\infty}\left(\Delta_{\star}^{Q}\right)\right)$ where $A_{Q}$ is the simultaneous Corkscrew point produced in Corollary 6.6, applied with $Q$ in place of $Q_{0}$.

Let $P_{\mathscr{G}} \nu$ be defined as in Lemma 6.15, but again with $Q$ in place of $Q_{0}$. We shall prove in Lemma B. 6 (Appendix B below) that $P_{\mathcal{G}} \nu \in A_{\infty}^{\text {dyadic }}(Q)$. Thus, by Lemma 6.15, with $Q$ in place of $Q_{0}$, we obtain that $P_{\mathscr{G}} \omega^{A_{Q}} \in A_{\infty}^{\text {dyadic }}(Q)$. We may then use Corollary 3.69 (here we only consider the case that $Q$ is not contained in any $Q_{j}$, otherwise $\mathscr{P}_{\mathscr{F}} \omega^{X_{Q_{0}}} \in A_{\infty}^{\text {dyadic }}(Q)$ trivially) along with the Harnack chain condition and a differentiation argument, to replace the pole $A_{Q}$ by $X_{Q_{0}}$, the Corkscrew point for the ambient cube $Q_{0}$, and to conclude that $\mathscr{P}_{g} \omega=\mathscr{P}_{\mathscr{G}} \omega^{X_{Q_{0}}} \in A_{\infty}^{\text {dyadic }}(Q)$. In particular, (8.7) holds, by Lemmas B. 7 and B. 1 in Appendix B.

### 8.3. Removing the qualitative hypothesis, and conclusion of the proofs of Theorems 1.26 and 1.27

In this subsection, we first complete the proof of Theorem 1.26 (modulo the proof of Lemma 8.5 and the technical lemmata that we have deferred to Appendices), by removing the qualitative exterior Corkscrew condition. We then conclude by giving the proof of Theorem 1.27.

We define approximating domains as follows. For each large integer $N$, set $\mathscr{F}_{N}:=\mathbb{D}_{N}$. We then let $\Omega_{N}:=\Omega_{\mathcal{F}_{N}}$ denote the usual (global) sawtooth with respect to the family $\mathscr{F}_{N}$ (cf. (3.49), (3.47) and (3.53).) Thus,

$$
\begin{equation*}
\Omega_{N}=\operatorname{int}\left(\bigcup_{Q \in \mathbb{D}: \ell(Q) \geq 2^{-N+1}} U_{Q}\right), \tag{8.13}
\end{equation*}
$$

so that $\overline{\Omega_{N}}$ is the union of fattened Whitney cubes $I^{*}=(1+\lambda) I$, with $\ell(I) \gtrsim 2^{-N}$, and the boundary of $\Omega_{N}$ consists of portions of faces of $I^{*}$ with $\ell(I) \approx 2^{-N}$. By virtue of Lemma 3.61, each $\Omega_{N}$ satisfies the ADR, Corkscrew and Harnack chain properties. Moreover, $\partial \Omega_{N}$ is UR. We defer the proof of the UR property to Appendix C. We note that, for each of these properties, the constants are uniform in $N$, and depend only on dimension and on the corresponding constants for $\Omega$. In addition, by construction, $\Omega_{N}$ has exterior Corkscrew points at all scales $\lesssim 2^{-N}$. By Lemma 3.62, the same statement applies to the Carleson boxes $T_{Q}$ and $T_{\Delta}$, and to the sawtooth domains $\Omega_{\mathcal{F}}$ and $\Omega_{\mathcal{F}, Q}$ (all of them relative to $\Omega_{N}$ ) and even to Carleson boxes within saw tooths.

Consequently, by the arguments in the previous subsection, we conclude that for every surface ball $\Delta_{\star}=\Delta_{\star}^{N} \subset \partial \Omega_{N}$, the harmonic measure $\omega_{N}^{X_{\Delta_{\star}}} \in A_{\infty}\left(\Delta_{\star}\right)$, uniformly in $N$. We now consider the limiting case. Fix a surface ball $\Delta:=\Delta(x, r) \subset \partial \Omega$, and a Borel subset $A \subset \Delta$. Assuming the hypotheses of Theorem 1.26, we claim that

$$
\begin{equation*}
\sigma(A) \geq \eta \sigma(\Delta) \Longrightarrow \omega^{X_{\Delta}}(A) \geq c_{0} \eta^{\theta}, \quad \forall \eta \in(0,1) \tag{8.14}
\end{equation*}
$$

for some uniform positive constants $c_{0}$ and $\theta$, where as usual $X_{\Delta}$ denotes a Corkscrew point relative to $\Delta$. By the outer regularity property of $\omega^{X}$ (cf. (3.3)), we may assume that $A$ is (relatively) open. It then follows that we may write $A=\bigcup_{k} Q_{k}$, where $\left\{Q_{k}\right\} \subset \mathbb{D}$ is a pairwise disjoint collection. We set $\Delta_{k}=\Delta\left(x_{k}, r_{k}\right):=\Delta_{Q_{k}}$, where $\Delta_{Q_{k}}:=B_{Q_{k}} \cap \partial \Omega \subset Q_{k}$ is the surface ball defined in (1.16)-(1.17), so that $r_{k} \approx \ell\left(Q_{k}\right)$ and $\sigma\left(Q_{k}\right) \approx \sigma\left(\Delta_{k}\right)$. Then

$$
\eta \sigma(\Delta) \leq \sigma(A)=\sum_{k} \sigma\left(Q_{k}\right) \approx \sum_{k} \sigma\left(\Delta_{k}\right)
$$

We now set $A^{\prime}:=\bigcup_{k=1}^{M} \Delta_{k}$, where $M$ is chosen large enough (depending on $A$ ) so that

$$
\begin{equation*}
\sigma\left(A^{\prime}\right)=\sum_{k=1}^{M} \sigma\left(\Delta_{k}\right) \geq \frac{1}{C} \eta \sigma(\Delta) \tag{8.15}
\end{equation*}
$$

By the ADR property and a covering lemma argument, we may further suppose that the Euclidean balls $B_{k}:=B_{Q_{k}}, 1 \leq k \leq M$, are pairwise disjoint. We now fix $N$ so large that $2^{-N} \ll \min _{1 \leq k \leq M} r_{k}$, and $2^{N} \gg \operatorname{diam}(\Delta)$. Fix also a point $\hat{x} \in \partial \Omega_{N}$, with $|x-\hat{x}| \approx 2^{-N}$ (such a point exists, with implicit constants possibly depending on $K_{0}$, since $\Omega$ satisfies the Corkscrew condition). We shall approximate $\Omega$ by a domain $\widehat{\Omega}_{N}$, which is defined as follows, and whose harmonic measure we denote by $\hat{\omega}^{X}$. If $\Omega$ is bounded, we set $\widehat{\Omega}_{N}=\Omega_{N}$. Otherwise, we define $\widehat{\Omega}_{N}:=T_{\Delta_{\star}^{N}}$, where $T_{\Delta_{\star}^{N}} \subset \Omega_{N}$ denotes the Carleson box corresponding to $\Delta_{\star}^{N}=B\left(\hat{x}, 2^{N}\right) \cap \partial \Omega_{N}$ for the domain $\Omega_{N}$. Then by the arguments in Subsection 8.2, for every surface ball $\Delta_{\star}=\Delta_{\star, N} \subset \partial \widehat{\Omega}_{N}$, the harmonic measure $\hat{\omega}_{N}^{X_{\Delta_{\star}}} \in A_{\infty}\left(\Delta_{\star}\right)$, uniformly in $N$, since $\widehat{\Omega}_{N}$ is either equal to $\Omega_{N}$, or else is a subdomain of $\Omega_{N}$ which inherits all of the requisite properties as observed above.

We set $B_{k}^{\prime}:=c B_{k}$, where $c \in(0,1)$ is the constant in Lemma 3.6. As noted above, the collection $\left\{B_{k}\right\}_{1 \leq k \leq M}$, hence also $\left\{B_{k}^{\prime}\right\}_{1 \leq k \leq M}$, may be taken to be pairwise disjoint. Let us also note that, since $2^{-N} \ll \min _{1 \leq k \leq M} r_{k}$, by the ADR properties of $\partial \widehat{\Omega}_{N}$ and $\partial \Omega$, we have

$$
\begin{equation*}
H^{n}\left(\bigcup_{k=1}^{M} B_{k}^{\prime} \cap \partial \widehat{\Omega}_{N}\right) \gtrsim \sum_{k=1}^{M} r_{k}^{n} \gtrsim \eta \sigma(\Delta) \gtrsim \eta H^{n}\left(\Delta_{\star}(\hat{x}, C r)\right) \tag{8.16}
\end{equation*}
$$

where in the last pair of inequalities we have used (8.15) and ADR (for both $\partial \Omega$ and $\partial \Omega_{N}$ ). Moreover, since $2^{-N} \ll \operatorname{diam}(\Delta) \ll 2^{N}$, we have that $X_{\Delta} \in \widehat{\Omega}_{N}$ is also a Corkscrew point for $\widehat{\Omega}_{N}$ with respect to the surface ball $\Delta_{\star}(\hat{x}, C r)$, where $\hat{x}$ is as above, and where $C$ is chosen large enough that $\bigcup_{k=1}^{M} B_{k}^{\prime} \cap \partial \widehat{\Omega}_{N} \subset \Delta_{\star}(\hat{x}, C r)$.

We observe that $u(X):=\omega^{X}\left(A^{\prime}\right)$ is harmonic in $\Omega$, and thus also in the bounded subdomain $\widehat{\Omega}_{N}$. Since in bounded domains we have uniqueness by the maximum principle, we obtain

$$
\begin{align*}
\omega^{X_{\Delta}}\left(A^{\prime}\right) & =\int_{\partial \hat{\Omega}_{N}} \omega^{Y}\left(A^{\prime}\right) d \hat{\omega}_{N}^{X_{\Delta}}(Y) \geq \sum_{k=1}^{M} \int_{B_{k}^{\prime} \cap \partial \hat{\Omega}_{N}} \omega^{Y}\left(\Delta_{k}\right) d \hat{\omega}_{N}^{X_{\Delta}}(Y)  \tag{8.17}\\
& \gtrsim \sum_{k=1}^{M} \int_{B_{k}^{\prime} \cap \partial \widehat{\Omega}_{N}} d \hat{\omega}_{N}^{X_{\Delta}}(Y)=\hat{\omega}_{N}^{X}\left(\bigcup_{k=1}^{M} B_{k}^{\prime} \cap \partial \widehat{\Omega}_{N}\right) \gtrsim \eta^{\theta},
\end{align*}
$$

where in the last line we have used Lemma 3.6 and then (8.16) and the $A_{\infty}$ property of $\hat{\omega}^{X_{\Delta}}$ (recall that $X_{\Delta}$ serves as a Corkscrew point for $\Delta_{\star}(\hat{x}, C r)$, as we have noted above). Since $A^{\prime} \subset A$, we then obtain (8.14).

We now note that (8.14) trivially implies the following weak version of itself: for bounded $\Omega$ (the unbounded case is treated below) satisfying the hypotheses of Theorem 1.26, there exist uniform constants $\eta \in(0,1)$ and $c_{0}>0$ such that

$$
\begin{equation*}
\sigma(A) \geq \eta \sigma(\Delta) \Longrightarrow \omega^{X_{\Delta}}(A) \geq c_{0} \tag{8.18}
\end{equation*}
$$

We remark here that to establish (8.18) has really been our main goal. Indeed, given (8.18), the remainder of the proof of Theorem 1.26 will follow the arguments in [7]. We further remark that in [7], (8.18) is essentially taken as a starting point: by the maximum principle, and the result of [15], an appropriate version of (8.18) (cf. (8.26) below) follows immediately from the main hypothesis in [7], that $\Omega$ has "interior big pieces" (in the sense of Definition 1.14) of Lipschitz sub-domains of $\Omega$, with uniform constants. Eventually, we shall see that (8.18), suitably interpreted, continues to hold under the hypotheses of Theorem 1.27.

We now proceed to describe the remaining steps needed to deduce the weak $-A_{\infty}$ property of harmonic measure. By [7, Lemma 3.1], it suffices to show that for each $\varepsilon \in(0,1 / 1000)$, there are uniform constants $\eta_{\varepsilon} \in(0,1)$ and $C_{\varepsilon} \in(1, \infty)$, such that given balls $B, B^{\prime}$, centered on $\partial \Omega$, with $2 B^{\prime} \subset B$, and corresponding surface balls $\Delta:=B \cap \partial \Omega$ and $\Delta^{\prime}:=B^{\prime} \cap \partial \Omega$, and a Borel subset $A \subset 2 \Delta^{\prime}$ with $\sigma(A) \geq \eta_{\varepsilon} \sigma\left(2 \Delta^{\prime}\right)$, we have

$$
\begin{equation*}
\omega^{X_{\Delta}}\left(\Delta^{\prime}\right) \leq \varepsilon \omega^{X_{\Delta}}\left(2 \Delta^{\prime}\right)+C_{\varepsilon} \omega^{X_{\Delta}}(A) \tag{8.19}
\end{equation*}
$$

In fact, [7, Lemma 3.1] is a purely real variable result which says that any positive Borel measure $\mu$ on $\partial \Omega$ satisfying (8.19) belongs to weak- $A_{\infty}(\Delta)$ (equivalently, satisfies the weak reverse Hölder estimate (1.25) for some $q>1$ ), assuming only that $\partial \Omega$ is ADR. Under the hypotheses of Theorem 1.26, we shall establish (8.19) with $\eta_{\varepsilon}:=\eta$, the constant in (8.18) (independently of $\varepsilon$ ).

Let us now give the proof of (8.19). We prove the desired bound first in the case that $\Omega$ is bounded. This restriction will be removed at the end of the proof. We follow the argument in [7, Lemma 2.2] almost verbatim, with some small simplifications permitted by our hypothesis that the Harnack chain condition holds in Theorem 1.26. Let $B^{\prime}=B(z, s)$,
$\Delta^{\prime}:=B^{\prime} \cap \partial \Omega$, and suppose $2 B^{\prime} \subseteq B:=B(x, r)$. We cover $\frac{3}{2} \Delta^{\prime} \backslash \frac{5}{4} \Delta^{\prime}$ by annuli of thickness $\approx \varepsilon s$. More precisely, we set

$$
\begin{align*}
U_{k} & :=\{y \in \partial \Omega:(5 / 4+\varepsilon k) s \leq|y-z|<(5 / 4+\varepsilon(k+1)) s\}  \tag{8.20}\\
S_{k} & :=\{X \in \Omega:|X-z|=(5 / 4+\varepsilon(k+1 / 2)) s\}
\end{align*}
$$

where $0 \leq k \lesssim 1 /(4 \varepsilon)$. Suppose now $A \subset 2 \Delta^{\prime}$, with $\sigma(A) \geq \eta \sigma\left(2 \Delta^{\prime}\right)$, for $\eta$ as in (8.18). Let $c \in(0,1)$ be the constant in Lemma 3.6. By the Harnack chain condition and (8.18), applied to $2 \Delta^{\prime}$ in place of $\Delta$, we have

$$
\begin{equation*}
\omega^{X}(A) \geq c_{\varepsilon} c_{0}, \quad \forall X \in S_{k} \cap\{X: \delta(X) \geq c \varepsilon s / 100\} \tag{8.21}
\end{equation*}
$$

uniformly in $k$. On the other hand, if $X \in S_{k} \cap\{X: \delta(X)<c \varepsilon s / 100\}$, then for a suitable uniform constant $C$, we have

$$
\begin{equation*}
C \omega^{X}\left(U_{k}\right) \geq C \omega^{X}(\Delta(\hat{x}, \varepsilon s / 10)) \geq 1 \tag{8.22}
\end{equation*}
$$

by Lemma 3.6, where $\hat{x} \in \partial \Omega$ is chosen so that $|X-\hat{x}|=\delta(X)$. Thus,

$$
\begin{equation*}
\omega^{X}\left(\Delta^{\prime}\right) \leq 1 \leq C \omega^{X}\left(U_{k}\right)+C_{\varepsilon} \omega^{X}(A), \quad \forall X \in S_{k} \tag{8.23}
\end{equation*}
$$

where $C_{\varepsilon}=1 /\left(c_{\varepsilon} c_{0}\right)$. By the maximum principle, this implies in particular that (8.23) continues to hold for $X \in \Omega \backslash \frac{7}{4} B^{\prime}$, since $S_{k} \subset \frac{7}{4} B^{\prime}$, if $\varepsilon$ is small, for every relevant $k$ (i.e., those for which $U_{k}$ meets $\frac{3}{2} B^{\prime} \backslash \frac{5}{4} B^{\prime}$ ). Since this set of $k$ 's has cardinality $\approx 1 / \varepsilon$, summing in $k$ we obtain

$$
\begin{equation*}
\frac{1}{\varepsilon} \omega^{X}\left(\Delta^{\prime}\right) \leq C \omega^{X}\left(2 \Delta^{\prime}\right)+C_{\varepsilon} \omega^{X}(A), \quad \forall X \in \Omega \backslash \frac{7}{4} B^{\prime} \tag{8.24}
\end{equation*}
$$

since the $U_{k}$ 's are pairwise disjoint and contained in $2 \Delta^{\prime}$. The desired bound (8.19) now follows, at least in the case that $\Omega$ is bounded.

Now suppose that $\Omega$ is unbounded. Given a surface ball $\Delta=\Delta(x, r)$, we choose $R \gg r$, set $\Delta_{R}=\Delta(x, R)$, and consider the domain $\Omega_{R}:=T_{\Delta_{R}}$, the Carleson Box associated to $\Delta_{R}$. For each such $R$, the argument above may be applied, to obtain (8.19) for each of the corresponding harmonic measures $\omega_{R}^{X_{\Delta}}$. For any fixed Borel subset $F \subset \Delta$, we have that the solutions $u_{R}(X):=\omega_{R}^{X}(F)$ are monotone increasing on any fixed $\Omega_{R_{0}}$, as $R_{0} \leq R \rightarrow \infty$, by the maximum principle. We then obtain that $u_{R}(X) \rightarrow u(X):=\omega^{X}(F)$, uniformly on compacta, by Harnack's convergence theorem (as in the discussion at the beginning of Section 3), whence (8.19) follows. The proof of Theorem 1.26 is now complete, modulo the deferred arguments.

Proof of Theorem 1.27. - Finally, we discuss the modifications needed to prove Theorem 1.27. By [7, Lemma 3.1] (and a limiting process to treat the case of an unbounded domain), it again suffices to establish, for bounded $\Omega$ now satisfying the hypotheses of Theorem 1.27, an appropriate version of (8.19). That is, we seek to show that for each $\varepsilon \in(0,1 / 1000)$, there are uniform constants $\eta_{\varepsilon} \in(0,1)$ and $C_{\varepsilon} \in(1, \infty)$, such that given balls $B=B(x, r)$ and $B^{\prime}=B(z, s)$, with $2 B^{\prime} \subset B$, and corresponding surface balls $\Delta:=B \cap \partial \Omega$ and $\Delta^{\prime}:=B^{\prime} \cap \partial \Omega$, if $A \subset 2 \Delta^{\prime}$ with $\sigma(A) \geq \eta_{\varepsilon} \sigma\left(2 \Delta^{\prime}\right)$, then

$$
\begin{equation*}
\omega^{X}\left(\Delta^{\prime}\right) \leq \varepsilon \omega^{X}\left(2 \Delta^{\prime}\right)+C_{\varepsilon} \omega^{X}(A), \quad \forall X \in \Omega \backslash B \tag{8.25}
\end{equation*}
$$

To this end, we first establish a suitable variant of (8.18). Given $X \in \Omega$, under the hypotheses of Theorem 1.27, there is a point $x \in \partial \Omega$, with $|X-x|=\delta(X)$, and a subdomain $\Omega^{\prime} \subset \Omega$ satisfying the hypotheses of Theorem 1.26 , with the property that for some constants $K>1$ and $\alpha>0$, we have

$$
\sigma\left(\partial \Omega^{\prime} \cap \Delta_{X}\right) \geq \alpha \sigma\left(\Delta_{X}\right)
$$

where $\Delta_{X}:=\Delta(x, K \delta(X))$. We may further suppose that $X$ serves as a Corkscrew point for $\Omega^{\prime}$ relative to a surface ball $\Delta_{\star}:=B(y, 2 K \delta(X)) \cap \partial \Omega^{\prime}$, with $y \in \Delta_{X} \cap \partial \Omega^{\prime}$. That $\Omega^{\prime}$ exists, with uniform control of the various constants involved, is simply a re-statement of the "big pieces" hypothesis of Theorem 1.27 (cf. Definition 1.14). We claim that there exist uniform constants $\eta \in(0,1)$ and $c_{0}>0$, such that for any Borel subset $A \subset \Delta_{X}$,

$$
\begin{equation*}
\sigma(A) \geq \eta \sigma\left(\Delta_{X}\right) \Longrightarrow \omega^{X}(A) \geq c_{0} \tag{8.26}
\end{equation*}
$$

Let us now prove this claim. Suppose that $A \subset \Delta_{X}$, with $\sigma(A) \geq(1-\alpha / 2) \sigma\left(\Delta_{X}\right)$. Then

$$
\sigma^{\prime}\left(\partial \Omega^{\prime} \cap A\right) \geq \frac{\alpha}{2} \sigma\left(\Delta_{X}\right) \approx \alpha \sigma^{\prime}\left(\Delta_{\star}\right)
$$

where $\sigma^{\prime}:=\left.H^{n}\right|_{\partial \Omega^{\prime}}$ denotes surface measure on $\partial \Omega^{\prime}$ (so $\sigma=\sigma^{\prime}$ on $\partial \Omega \cap \partial \Omega^{\prime}$ ), and where we have used that $\partial \Omega$ and $\partial \Omega^{\prime}$ are both ADR. Since the hypotheses of Theorem 1.26 apply in $\Omega^{\prime}$, we deduce from (8.14) and a formal application of the maximum principle that

$$
\begin{equation*}
\alpha^{\theta} \lesssim \omega_{\Omega^{\prime}}^{X}\left(\partial \Omega^{\prime} \cap A\right) \leq \omega^{X}(A), \tag{8.27}
\end{equation*}
$$

where $\omega_{\Omega^{\prime}}$ is harmonic measure for $\Omega^{\prime}$. Thus, we obtain (8.26), with $\eta=(1-\alpha / 2)$. We caution the reader that our use of the maximum principle to obtain the second inequality in (8.27) is not routine, since we are working in a regime where the Wiener test may fail, and our solutions $X \rightarrow \omega^{X}(A)$ and $X \rightarrow \omega_{\Omega^{\prime}}^{X}\left(\partial \Omega^{\prime} \cap A\right)$ are not Perron solutions for the same domain, nor are they continuous on the closures of the respective domains under consideration. We shall give a rigorous justification of the essential inequality in (8.27) (namely, that $\alpha^{\theta} \lesssim \omega^{X}(A)$ ), at the end of this section.

It remains to establish (8.25). To this end, we again follow the argument in [7, Lemma 2.2]. Fix $B$ and $B^{\prime}$ as above, and define $U_{k}$ and $S_{k}$ as in (8.20). In fact, we proceed as we did under the hypotheses of Theorem 1.26, except that the proof of (8.21) will now be somewhat more delicate, as we may no longer simply invoke the Harnack chain condition. Instead, we return to the original approach of [7]. It is enough to verify (8.23), as the remainder of the proof is unchanged. In particular, we obtain (8.24), which in turn yields (8.25), since $2 B^{\prime} \subset B$.

As before, (8.23) is a direct consequence of (8.21) and (8.22). The latter always holds, by Lemma 3.6, so we consider (8.21). Again we follow [7] essentially verbatim. We suppose first that there exists $Y \in S_{k}$ with $\delta(Y)=c \varepsilon s /(100 K)$, where $c$ is the constant in Lemma 3.6. For each such $Y$, we fix $y \in \partial \Omega$, with $|Y-y|=\delta(Y)$, and set $\Delta_{Y}:=\Delta(y, K \delta(Y))$. If $\eta_{\varepsilon} \in(0,1)$ is chosen close enough to 1 , depending on $\varepsilon$ and the ADR constants of $\partial \Omega$, and if $A \subset 2 \Delta^{\prime}$ with $\sigma(A) \geq \eta_{\varepsilon} \sigma\left(2 \Delta^{\prime}\right)$, then

$$
\sigma\left(A \cap \Delta_{Y}\right) \geq \eta \sigma\left(\Delta_{Y}\right)
$$

for $\eta$ as in (8.26), so that $\omega^{Y}(A) \geq \omega^{Y}\left(A \cap \Delta_{Y}\right) \geq c_{0}$. Thus, (8.21) holds in this case (with $c \varepsilon s / 100$ now multiplied by $1 / K$ ), by Harnack's inequality, because even in the absence of the

Harnack chain condition, there is a Harnack path from any

$$
X \in S_{k} \cap\{X: \delta(X) \geq c \varepsilon s /(100 K)\}
$$

to a point $Y$ in $S_{k}$ with $\delta(Y)=c \varepsilon s /(100 K)$, if the latter exists (just follow a geodesic path on $S_{k}$ from $X$ to the nearest such Y).

On the other hand, suppose that there is no such $Y$. Then either

$$
S_{k} \subset\{X \in \Omega: \delta(X)>c \varepsilon s /(100 K)\}
$$

or $S_{k} \subset\{X \in \Omega: \delta(X)<c \varepsilon s /(100 K)\}$. In the latter case, (8.22) holds now for all $X \in S_{k}$, so (8.23) follows trivially. Otherwise, by continuity of $\delta$, there is a number $\rho>0$ such that

$$
\begin{align*}
\{X \in \Omega: \rho \leq|X-z| \leq(5 / 4+\varepsilon(k+1 / 2)) s\} &  \tag{8.28}\\
& \subset\{X \in \Omega: \delta(X) \geq c \varepsilon s /(100 K)\}
\end{align*}
$$

and $\delta(Y)=c \varepsilon s /(100 K)$ for some $Y \in S(\rho):=\{X \in \Omega:|X-z|=\rho\}$. In this case, we may repeat the analysis above, in which there was such a $Y$ on $S_{k}$. In the present scenario, we have that (8.21) holds for all $X \in S(\rho) \cap\{X \in \Omega: \delta(X) \geq c \varepsilon s /(100 K)\}$, which in fact is all of $S(\rho)$ by (8.28). But then by Harnack's inequality we obtain (8.21) (for all $X \in S_{k}$ ), because the containment in (8.28) allows us to form a radial Harnack path between any $X \in S_{k}$, and its projection onto $S(\rho)$. We conclude that (8.23) holds under all circumstances.

To finish the proof of Theorem 1.27, it remains only to provide a rigorous justification of (8.27). We shall make up for the lack of continuity of the solutions by proceeding as in the removal of the qualitative hypothesis in the proof of Theorem 1.26, with a few minor modifications. We fix $\varepsilon_{1}>0$ to be chosen momentarily, and set $F:=A \cap \partial \Omega^{\prime}$. We recall that $H^{n}(F) \geq(\alpha / 2) H^{n}\left(\Delta_{X}\right)$. By outer regularity of Hausdorff measure and $\omega$, there is a set $\theta$, relatively open in $\partial \Omega$, such that $F \subset \theta \subset \Delta_{X} \subset \partial \Omega$, and

$$
H^{n}(\theta \backslash F)+\omega^{X}(\theta \backslash F)<\varepsilon_{1} .
$$

We let $\mathcal{F} \subset \mathbb{D}(\partial \Omega)$ be a family of non-overlapping dyadic cubes whose union equals $\theta$, so that $H^{n}(\vartheta)=\sum_{\mathcal{G}} H^{n}\left(Q_{k}\right)$, and we set

$$
\mathcal{F}^{\prime}:=\left\{Q_{k} \in \mathcal{F}: H^{n}\left(Q_{k} \cap F\right) \geq \frac{1}{4} H^{n}\left(Q_{k}\right)\right\} .
$$

We claim that

$$
\begin{equation*}
\alpha H^{n}\left(\Delta_{X}\right) \lesssim \sum_{\mathcal{G}^{\prime}} H^{n}\left(Q_{k}\right) \tag{8.29}
\end{equation*}
$$

Indeed, we have that

$$
H^{n}(F)=\sum_{\mathscr{G} \backslash \mathcal{G}^{\prime}} H^{n}\left(Q_{k} \cap F\right)+\sum_{\mathcal{G}^{\prime}} H^{n}\left(Q_{k} \cap F\right) \leq \frac{1}{4} H^{n}(\theta)+\sum_{\mathcal{G}^{\prime}} H^{n}\left(Q_{k}\right),
$$

whence (8.29) follows, if we choose $\varepsilon_{1} \ll H^{n}(F)$.
Since each $Q_{k} \in \mathcal{F}^{\prime}$ has an ample intersection with $F$, by Lemma ?? (vi), we may choose a point $x_{k} \in Q_{k} \cap F \subset \partial \Omega^{\prime} \cap \partial \Omega$, and a radius $r_{k} \approx \ell\left(Q_{k}\right)$, such that $Q_{k} \supset \partial \Omega \cap B_{k}$, where $B_{k}:=B\left(x_{k}, r_{k}\right)$. We emphasize that, in particular, each $B_{k}$ is centered on $\partial \Omega^{\prime} \cap \partial \Omega$. Set $B_{k}^{\prime}:=c B_{k}$, where $c \in(0,1)$ is the constant in Lemma 3.6. Set $\mathcal{F}_{N}^{\prime}:=\mathbb{D}_{N}\left(\partial \Omega^{\prime}\right)$, and let $\Omega_{N}^{\prime}:=\Omega_{\mathcal{F}_{N}}^{\prime}$ be the corresponding approximating domain relative to $\Omega^{\prime}$. By the ADR
property and a covering lemma argument, and by choice of $N$ sufficiently large, we can select a finite, pairwise disjoint sub-collection $\left\{B_{k}\right\}_{1 \leq k \leq M}$, such that

$$
\alpha H^{n}\left(\Delta_{\star}^{N}\right) \approx \alpha H^{n}\left(\Delta_{X}\right) \lesssim \sum_{k=1}^{M} H^{n}\left(B_{k}^{\prime} \cap \partial \Omega_{N}^{\prime}\right)
$$

where $\Delta_{\star}^{N}$ is a surface ball on $\partial \Omega_{N}^{\prime}$ of radius $\approx \delta(X)$, such that $X$ is a Corkscrew point for $\Delta_{\star}^{N}$ in $\Omega_{N}^{\prime}$, and $\bigcup_{k=1}^{M}\left(B_{k}^{\prime} \cap \partial \Omega_{N}^{\prime}\right) \subset \Delta_{\star}^{N}$. We may apply Theorem 1.26 in $\Omega_{N}^{\prime}$ (see the discussion immediately following (8.13) above), to obtain that $\omega_{N}^{X}$, the harmonic measure for the approximating domain $\Omega_{N}^{\prime}$, belongs to $A_{\infty}\left(\Delta_{\star}^{N}\right)$ with bounds that are independent of $N$.

We now set $A^{\prime}:=\left(\bigcup_{k=1}^{M} B_{k}\right) \cap \partial \Omega$, and observe that $A^{\prime} \subset \theta$. Since $X \rightarrow \omega^{X}\left(A^{\prime}\right)$ is continuous on $\overline{\Omega_{N}^{\prime}}$, we may repeat the argument in (8.17), mutatis mutandis, to obtain that

$$
\omega^{X}(A)+\varepsilon_{1} \geq \omega^{X}(F)+\varepsilon_{1} \geq \omega^{X}(\theta) \geq \omega^{X}\left(A^{\prime}\right) \gtrsim \alpha^{\theta} .
$$

We choose $\varepsilon_{1} \ll \alpha^{\theta}$, and it follows that $\alpha^{\theta} \lesssim \omega^{X}(A)$, as desired.

## 9. Proof of the extrapolation lemma

To finish the proofs of Theorems 1.26 and 1.27, it remains to prove Lemma 8.5.
Proof of Lemma 8.5. - The proof follows the strategy introduced in [36], and developed further in [25, 3] and [4]. In more precise detail, the argument is based on the systematic treatment given in [27] in the Euclidean setting.

The proof uses an induction argument with continuous parameter. The induction hypothesis is the following: given $a \geq 0$,

There exist $\eta_{a} \in(0,1)$ and $C_{a}<\infty$ such that for every $Q \in \mathbb{D}_{Q_{0}}$ satisfying $\mathfrak{m}\left(\mathbb{D}_{Q}\right) \leq a \sigma(Q)$, it follows that

$$
F \subset Q, \quad \frac{\sigma(F)}{\sigma(Q)} \geq 1-\eta_{a} \quad \Longrightarrow \quad \frac{\omega(F)}{\omega(Q)} \geq \frac{1}{C_{a}} .
$$

The induction argument is split in two steps.
Step 1. - Show that $H(0)$ holds.
Step 2. - Show that there exists $b$ depending on $\gamma$, dimension, and the ADR property such that for all $0 \leq a \leq M_{0}, H(a)$ implies $H(a+b)$.

Once these steps have been carried out, the proof follows easily: pick $k \geq 1$ such that $(k-1) b<M_{0} \leq k b$ (note that $k$ only depends on $b$ and $M_{0}$ ). By Step 1 and Step 2, it follows that $H(k b)$ holds. Observe that $\|\mathfrak{m}\|_{\mathscr{G}\left(Q_{0}\right)} \leq M_{0} \leq k b$ implies $\mathfrak{m}\left(\mathbb{D}_{Q}\right) \leq k b \sigma(Q)$ for all $Q \in \mathbb{D}_{Q_{0}}$, and by $H(k b)$ we conclude (8.8).

## Step 1. $H(0)$ holds

If $\mathfrak{m}\left(\mathbb{D}_{Q}\right)=0$ then we take $\mathcal{F}$ to be empty, so that $\mathbb{D}_{Q} \cap \mathbb{D}_{\mathcal{F}}=\mathbb{D}_{Q}$, and $\mathscr{P}_{\mathscr{G}} \omega=\omega$. Then (8.6) holds (since $0 \leq \gamma$ ) and therefore we can use (8.7) with $\omega$ in place of $\mathscr{P}_{\mathcal{G}} \omega$, which is the desired property.

## Step 2. $H(a)$ implies $H(a+b)$

Fix $0 \leq a \leq M_{0}$ and $Q \in \mathbb{D}_{Q_{0}}$ such that $\mathfrak{m}\left(\mathbb{D}_{Q}\right) \leq(a+b) \sigma(Q)$, where we choose $b$ so that $C b:=\gamma$ and $C$ is the constant in the right hand side of (7.3). We also fix $F \subset Q$ with $\sigma(F) \geq(1-\eta) \sigma(Q)$, where $0<\eta \leq \eta_{a, b}$ and $\eta_{a, b}$ is to be chosen. We may now apply Lemma 7.2 and Remark 7.5 to the cube $Q$, to construct the non-overlapping family of cubes $\mathcal{F}=\left\{Q_{j}\right\} \subset \mathbb{D}_{Q}$ with the stated properties. Set

$$
E_{0}=Q \backslash \bigcup_{Q_{j} \in \mathcal{G}} Q_{j}, \quad G=\bigcup_{Q_{j} \in \mathcal{F}_{\text {good }}} Q_{j}, \quad B=\bigcup_{Q_{j} \in \mathscr{T} \backslash \mathcal{G}_{\text {good }}} Q_{j},
$$

where $\mathcal{F}_{\text {good }}=\left\{Q_{j} \in \mathscr{F}: \mathfrak{m}\left(\mathbb{D}_{Q_{j}}^{\text {short }}\right) \leq a \sigma\left(Q_{j}\right)\right\}$. We recall that by (7.4), we have $\sigma(B) / \sigma(Q) \leq(a+b) /(a+2 b)$.

We shall also require the following "pigeon hole" lemma, which says that "most" of the cubes $Q_{j}$ have an ample overlap with $F$.

Lemma 9.1. - Given $0<\tilde{\eta}<1$, we set

$$
\mathcal{F}_{1}=\left\{Q_{j} \in \mathcal{F}_{\text {good }}: \sigma\left(F \cap Q_{j}\right) \geq(1-\tilde{\eta}) \sigma\left(Q_{j}\right)\right\}, \quad \quad G_{1}=\bigcup_{Q_{j} \in \mathscr{F}_{1}} Q_{j}
$$

If $0<\eta \leq \eta_{1}:=\tilde{\eta} \frac{1}{2}\left(1-\frac{M_{0}+b}{M_{0}+2 b}\right)$, then $\sigma\left(E_{0} \cup G_{1}\right) \geq \eta_{1} \sigma(Q)$.
Proof. - Take $\theta$ such that $\sigma(B)=\theta \sigma(Q)$, and $\theta_{0}=\left(M_{0}+b\right) /\left(M_{0}+2 b\right)$. By (7.4) and since $a \leq M_{0}$ we obtain that $\theta \leq \theta_{0}$ :

$$
\theta \sigma(Q)=\sigma(B) \leq \frac{a+b}{a+2 b} \sigma(Q) \leq \theta_{0} \sigma(Q)
$$

We set $B_{1}=\bigcup_{Q_{j} \in \mathcal{F}_{\text {good }} \backslash \mathcal{I}_{1}} Q_{j}$ and observe that $B_{1} \subset G \subset Q \backslash B$. Hence,

$$
\begin{aligned}
\sigma\left(F \cap B_{1}\right) & =\sum_{Q_{j} \in \mathcal{F}_{\text {good }} \backslash \mathscr{I}_{1}} \sigma\left(F \cap Q_{j}\right)<(1-\tilde{\eta}) \sum_{Q_{j} \in \mathcal{G}_{\text {good }} \backslash \mathscr{F}_{1}} \sigma\left(Q_{j}\right) \\
& =(1-\tilde{\eta}) \sigma\left(B_{1}\right) \leq(1-\tilde{\eta}) \sigma(Q \backslash B)=(1-\tilde{\eta})(1-\theta) \sigma(Q) .
\end{aligned}
$$

Thus, using that $\theta \leq \theta_{0}$, we have

$$
\begin{aligned}
(1-\eta) \sigma(Q) & \leq \sigma(F)=\sigma\left(F \cap E_{0}\right)+\sigma(F \cap B)+\sigma\left(F \cap G_{1}\right)+\sigma\left(F \cap B_{1}\right) \\
& \leq \sigma\left(E_{0}\right)+\sigma(B)+\sigma\left(G_{1}\right)+(1-\tilde{\eta})(1-\theta) \sigma(Q) \\
& =\sigma\left(E_{0}\right)+\sigma\left(G_{1}\right)+[\theta+(1-\tilde{\eta})(1-\theta)] \sigma(Q) \\
& \leq \sigma\left(E_{0}\right)+\sigma\left(G_{1}\right)+\left[1-\tilde{\eta}\left(1-\theta_{0}\right)\right] \sigma(Q)
\end{aligned}
$$

and therefore

$$
\sigma\left(E_{0} \cup G_{1}\right)=\sigma\left(E_{0}\right)+\sigma\left(G_{1}\right) \geq\left[\tilde{\eta}\left(1-\theta_{0}\right)-\eta\right] \sigma(Q) \geq \frac{1}{2} \tilde{\eta}\left(1-\theta_{0}\right) \sigma(Q)=\eta_{1} \sigma(Q),
$$

where we have used that $\eta \leq \tilde{\eta}\left(1-\theta_{0}\right) / 2=\eta_{1}$.
We now return to the proof of Step 2. We apply Lemma 9.1, with $\tilde{\eta} \in(0,1)$ to be chosen. Given $Q_{j} \in \mathscr{F}_{1} \subset \mathcal{F}_{\text {good }}$ we have that $\mathfrak{m}\left(\mathbb{D}_{Q_{j}}^{\text {short }}\right) \leq a \sigma\left(Q_{j}\right)$. Moreover,

$$
\mathbb{D}_{Q_{j}}^{\text {short }}=\mathbb{D}_{Q_{j}} \backslash\left\{Q_{j}\right\}=\bigcup_{i} \mathbb{D}_{Q_{i}^{j}},
$$

where $\left\{Q_{i}^{j}\right\}_{i}$ is the family of dyadic "children" of $Q_{j}$ (these are the subcubes of $Q_{j}$ which lie in the very next dyadic generation $\left.\mathbb{D}_{k\left(Q_{j}\right)+1}\right)$. Then by pigeon-holing, there exists at least one $i_{0}$ such that $Q_{i_{0}}^{j}=: Q_{j}^{\prime}$ satisfies

$$
\begin{equation*}
\mathfrak{m}\left(\mathbb{D}_{Q_{j}^{\prime}}\right) \leq a \sigma\left(Q_{j}^{\prime}\right) \tag{9.2}
\end{equation*}
$$

(there could be more than one $i_{0}$ with this property, but we just pick one). We define $\widetilde{\mathcal{F}}_{1}$ to be the collection of those selected "children" $Q_{j}^{\prime}$, with $Q_{j} \in \mathscr{F}_{1}$. Let $C_{0}$ be the dyadically doubling constant of $\sigma$, i.e., $\sigma(Q) \leq C_{0} \sigma\left(Q^{\prime}\right)$ for every $Q \in \mathbb{D}_{Q_{0}}$, and for every "child" $Q^{\prime}$ of $Q$. Then, for each such $Q_{j}^{\prime}$, using the definition of $\mathscr{F}_{1}$, and taking $0<\tilde{\eta}=\eta_{a} / C_{0}$ (where $0<\eta_{a}<1$ is provided by $H(a)$ ), we have

$$
\sigma\left(Q_{j}^{\prime} \backslash F\right) \leq \sigma\left(Q_{j} \backslash F\right) \leq \tilde{\eta} \sigma\left(Q_{j}\right) \leq \tilde{\eta} C_{0} \sigma\left(Q_{j}^{\prime}\right)=\eta_{a} \sigma\left(Q_{j}^{\prime}\right),
$$

which yields $\sigma\left(Q_{j}^{\prime} \cap F\right) \geq\left(1-\eta_{a}\right) \sigma\left(Q_{j}^{\prime}\right)$. With this estimate and (9.2) in hand, we can use the induction hypothesis $H(a)$ to deduce:

$$
\begin{equation*}
\omega\left(Q_{j}^{\prime} \cap F\right) \geq \frac{1}{C_{a}} \omega\left(Q_{j}^{\prime}\right), \quad \forall Q_{j}^{\prime} \in \widetilde{\mathscr{F}}_{1} \tag{9.3}
\end{equation*}
$$

On the other hand, if we set $\widetilde{G}_{1}=\bigcup_{Q_{j}^{\prime} \in \widetilde{\mathscr{G}}_{1}} Q_{j}^{\prime}$, then

$$
\sigma\left(\widetilde{G}_{1}\right)=\sum_{Q_{j}^{\prime} \in \widetilde{\mathscr{F}}_{1}} \sigma\left(Q_{j}^{\prime}\right) \geq C_{0}^{-1} \sum_{Q_{j} \in \mathcal{F}_{1}} \sigma\left(Q_{j}\right)=C_{0}^{-1} \sigma\left(G_{1}\right) .
$$

Thus, by Lemma 9.1, having now fixed $\tilde{\eta}$ above, we have that

$$
\sigma\left(E_{0} \cup \widetilde{G}_{1}\right)=\sigma\left(E_{0}\right)+\sigma\left(\widetilde{G}_{1}\right) \geq C_{0}^{-1} \sigma\left(E_{0} \cup G_{1}\right) \geq C_{0}^{-1} \eta_{1} \sigma(Q)=: \eta_{2} \sigma(Q),
$$

if $\eta \leq \eta_{1}$, from which it follows that

$$
\sigma\left(F \cap\left(E_{0} \cup \widetilde{G}_{1}\right)\right) \geq \frac{1}{2} \eta_{2} \sigma(Q)=: \eta_{a, b} \sigma(Q),
$$

if $\eta \leq \eta_{2} / 2$, since $\sigma(Q \backslash F) \leq \eta \sigma(Q)$.
We recall that the family $\mathcal{F}$ was constructed using Lemma 7.2 with $C b:=\gamma$. Consequently, by (7.3), we may deduce that (8.6) holds, so in turn, by hypothesis, we can apply (8.7) to the set $F \cap\left(E_{0} \cup \widetilde{G}_{1}\right)$, obtaining

$$
\frac{\mathscr{P}_{g} \omega\left(F \cap\left(E_{0} \cup \widetilde{G}_{1}\right)\right)}{\mathscr{P}_{\mathscr{G}} \omega(Q)} \geq \frac{1}{C_{\eta_{a}, b}} .
$$

As observed before, $\mathscr{P}_{\mathscr{G}} \omega(Q)=\omega(Q)$. Thus, in order to establish the conclusion of $H(a+b)$, and consequently to complete the proof of Lemma 8.5 , it remains only to show that

$$
\mathscr{P}_{\mathscr{F}} \omega\left(F \cap\left(E_{0} \cup \widetilde{G}_{1}\right)\right) \leq C \omega(F) .
$$

To this end, we use first the definition of $\mathscr{P}_{\mathscr{F}}$, then that $\omega$ is dyadically doubling and finally (9.3) to obtain

$$
\begin{aligned}
\mathscr{P}_{\mathscr{F}} \omega\left(F \cap\left(E_{0} \cup \widetilde{G}_{1}\right)\right) & =\mathscr{P}_{\mathscr{G}} \omega\left(F \cap E_{0}\right)+\mathscr{P}_{\mathscr{G}} \omega\left(F \cap \widetilde{G}_{1}\right) \\
& =\omega\left(F \cap E_{0}\right)+\sum_{Q_{j} \in \mathcal{F}_{1}} \frac{\sigma\left(Q_{j}^{\prime} \cap F\right)}{\sigma\left(Q_{j}\right)} \omega\left(Q_{j}\right) \\
& \leq \omega\left(F \cap E_{0}\right)+C_{\omega} \sum_{Q_{j}^{\prime} \in \widetilde{\mathscr{F}}_{1}} \omega\left(Q_{j}^{\prime}\right) \\
& \leq \omega\left(F \cap E_{0}\right)+C_{\omega} C_{a} \sum_{Q_{j}^{\prime} \in \widetilde{\mathscr{G}}_{1}} \omega\left(Q_{j}^{\prime} \cap F\right) \leq C \omega(F)
\end{aligned}
$$

This concludes the proof of Lemma 8.5.

## Appendix A

## Inheritance of properties by Carleson and sawtooth regions

This section is devoted to the proof of Lemma 3.61, which states that Carleson and Sawtooth regions inherit the Corkscrew, Harnack chain and ADR properties from the original domain $\Omega$. Moreover, in the presence of the Corkscrew, Harnack chain and ADR properties, the UR property is transmitted to the Carleson boxes $T_{Q}$ and $T_{\Delta}$. We discuss these properties one at a time. We shall find it convenient for our purposes in this section to continue to let $\Delta$ denote a surface ball on $\partial \Omega$, while $\Delta_{\star}$ will denote a surface ball on the boundary of the sub-domain under consideration. Similarly $\delta(X)$ will continue to denote the distance from $X$ to $\partial \Omega$, while $\delta_{\star}(X)$ will denote the distance from $X$ to the boundary of the sub-domain under consideration.

In order to avoid possible confusion, let us emphasize that the construction of our sawtooth and Carleson sub-domains is always based on the Whitney decomposition of the domain under consideration at that moment, even if that domain happens to be, say, an approximating domain $\Omega_{N}$ which had been constructed in the first place from Whitney cubes of the original domain $\Omega$.

## A.1. Corkscrew

For the sake of specificity, we treat only the case of a local sawtooth region $\Omega_{\mathcal{F}, Q}$. The proof for the global sawtooth $\Omega_{\mathscr{F}}$ is almost identical. Moreover, specializing to the case that $\mathcal{F}=\varnothing$, we see that the result for a sawtooth $\Omega_{\mathcal{F}, Q}$ applies immediately to the Carleson box $T_{Q}$, and therefore also almost immediately to any box $T_{\Delta}$, since the latter is a union of a bounded number of $T_{Q}$ 's.

We fix $Q_{0} \in \mathbb{D}$, and a pairwise disjoint family $\left\{Q_{j}\right\}=\mathscr{F} \subset \mathbb{D}_{Q_{0}}$, and let $\Omega_{\mathcal{F}, Q_{0}}$ denote the associated local sawtooth region (cf. (3.39)-(3.54)). Set

$$
\Delta_{\star}:=\Delta_{\star}(x, r):=B(x, r) \cap \partial \Omega_{\mathcal{F}, Q_{0}}
$$

$4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - No 3
with $r \lesssim \ell\left(Q_{0}\right)$ and $x \in \partial \Omega_{\mathcal{F}, Q_{0}}$. We suppose first that $x \in \partial \Omega_{\mathscr{F}, Q_{0}} \cap \partial \Omega$. Then by construction of $\Omega_{\mathcal{G}, Q_{0}}$, there is a $Q \in \mathbb{D}_{\mathcal{G}, Q_{0}}$, with $x \in \bar{Q}$, and $r \approx 100 K_{0} \ell(Q)$ (see Proposition 6.1). Consequently, by (3.47)-(3.49), we have

$$
X_{Q} \in U_{Q} \subset B(x, r) \cap \Omega_{\mathcal{G}, Q_{0}}
$$

where $X_{Q}$ is a Corkscrew point for $\Omega$, relative to $Q$, and which we have assumed (without loss of generality) to be the center of some $I \in \mathcal{W}_{Q}^{*}$. This same $X_{Q}$ then serves as a Corkscrew point for $\Omega_{\mathcal{G}, Q_{0}}$, relative to $\Delta_{\star}(x, r)$, with Corkscrew constant $c \approx 1 /\left(100 K_{0}\right)$.

Next, we suppose that $x \in \partial \Omega_{\mathcal{F}, Q_{0}} \backslash \partial \Omega$, where as above $\Delta_{\star}:=\Delta_{\star}(x, r)$. Then by definition of the sawtooth region, $x$ lies on a face of a fattened Whitney cube $I^{*}=(1+\lambda) I$, with $I \in W_{Q}^{*}$, for some $Q \in \mathbb{D}_{\mathcal{G}, Q_{0}}$. If $r \lesssim \ell(I)$, then trivially there is a point $X^{\star} \in I^{*}$ such that $B\left(X^{\star}, c r\right) \subset B(x, r) \cap \operatorname{int}\left(I^{*}\right) \subset B(x, r) \cap \Omega_{\mathcal{G}, Q_{0}}$. This $X^{\star}$ is then a Corkscrew point for $\Delta_{\star}$. On the other hand, if $\ell(I)<r /\left(M K_{0}\right)$, with $M$ sufficiently large to be chosen momentarily, then there is a $Q^{\prime} \in \mathbb{D}_{\mathcal{G}, Q_{0}}$, with $\ell\left(Q^{\prime}\right) \approx r /\left(M K_{0}\right)$, and $Q \subseteq Q^{\prime}$. Now fix $I^{\prime} \in \mathcal{W}_{Q^{\prime}} \subset \mathcal{W}_{Q^{\prime}}^{*}$, and observe that

$$
\left|x-X\left(I^{\prime}\right)\right| \lesssim \operatorname{dist}(I, Q)+\operatorname{dist}\left(Q^{\prime}, I^{\prime}\right) \lesssim K_{0} \ell(I)+K_{0} \ell\left(I^{\prime}\right) \lesssim r / M .
$$

Note that $B\left(X\left(I^{\prime}\right), c r\right) \subset \operatorname{int}\left(I^{\prime}\right)$, for $c \approx\left(M K_{0}\right)^{-1}$. Moreover, for $M$ large enough we have that $B\left(X\left(I^{\prime}\right), c r\right) \subset B(x, r) \cap \operatorname{int}\left(I^{\prime}\right) \subset B(x, r) \cap \Omega_{\mathcal{G}, Q_{0}}$, so that $X\left(I^{\prime}\right)$ is a Corkscrew point for $\Delta_{\star}$.

## A.2. Harnack chain

We establish the Harnack chain condition for a local sawtooth $\Omega_{\mathcal{G}, Q}$, of which, as noted above, the Carleson box $T_{Q}$ is a special case (with $\mathcal{F}=\varnothing$ ). The proof for a global sawtooth is almost the same, and we omit it. We shall discuss the Carleson boxes $T_{\Delta}$ at the end of this subsection.

Fix $Q \in \mathbb{D}$, and a pairwise disjoint family $\mathcal{F} \subset \mathbb{D}_{Q}$, and let $\Omega_{\mathscr{G}, Q}$ be the corresponding local sawtooth region. Let $X_{1}, X_{2} \in \Omega_{\mathcal{G}, Q}$. By definition of the sawtooth regions, there exist $Q_{1}, Q_{2} \in \mathbb{D}_{\mathcal{G}, Q}$, with $X_{i} \in(1+\lambda) I_{i}=I_{i}^{*}$ where $I_{i} \in W_{Q_{i}}^{*}, i=1,2$. Without loss of generality we may suppose that $\ell\left(Q_{1}\right) \leq \ell\left(Q_{2}\right)$. We first observe that the desired result is clear if $I_{1}=I_{2}$, or more generally, if $I_{1}^{*}$ and $I_{2}^{*}$ overlap. Therefore, we may suppose that

$$
\begin{equation*}
\operatorname{dist}\left(I_{1}^{*}, I_{2}^{*}\right) \gtrsim \ell\left(I_{2}\right) \gtrsim \ell\left(I_{1}\right) \tag{A.1}
\end{equation*}
$$

(cf. (3.50).) In order to construct a Harnack chain under these circumstances, relative to $\Omega_{\mathcal{F}, Q}$, from $X_{1}$ to $X_{2}$, it is convenient to make a few simple reductions and observations, as follows.
(1) It is enough to treat the case that $X_{i}$ is the center of $I_{i}$. If $X_{i}$ is near the boundary of the sawtooth (and hence also near the boundary of $I_{i}^{*}$ ), then $\operatorname{dist}\left(X_{i}, \partial I_{i}^{*}\right) \approx$ $\operatorname{dist}\left(X_{i}, \partial \Omega_{\mathcal{G}, Q}\right)$, so that the Harnack chain within $I_{i}^{*}$, that connects $X_{i}$ to the center $X\left(I_{i}\right)$, is also a Harnack chain for the sawtooth. On the other hand, if $X_{i}$ is not near the boundary of the sawtooth, then we can easily join $X_{i}$ with $X\left(I_{i}\right)$ by a bounded number of balls of radius $\approx \ell\left(I_{i}\right)$ with distance to the boundary of $\partial \Omega_{\mathcal{G}, Q}$ comparable to $\delta_{\star}\left(X_{i}\right)$.
(2) By construction (cf. (3.47)-(3.49)), we may then further suppose that $X_{i}=X_{Q_{i}}$, the designated Corkscrew point (for the ambient domain $\Omega$ ), relative to $Q_{i}$.
(3) Recall that by construction, if $Q^{\prime} \subset Q^{\prime \prime}$ belong to consecutive generations in $\mathbb{D}$ (i.e., $k\left(Q^{\prime \prime}\right)=k\left(Q^{\prime}\right)-1$ ), then $U_{Q^{\prime}} \cap U_{Q^{\prime \prime}}$ contains the Corkscrew point $X_{Q^{\prime}}$ (cf. (3.48)) and is therefore non-empty. Thus, by (3.47)-(3.49) there is a Harnack chain joining the respective Corkscrew points $X_{Q^{\prime}}$ and $X_{Q^{\prime \prime}}$.
(4) We note that by definition, if $Q_{i} \in \mathbb{D}_{\mathcal{G}, Q}$, then also $Q^{\prime} \in \mathbb{D}_{\mathcal{G}, Q}$ for every $Q^{\prime}$ such that $Q_{i} \subseteq Q^{\prime} \subseteq Q$.
(5) If $X(I)$ denotes the center of a Whitney cube $I$, then $\delta(X(I)) \approx \delta_{\star}(X(I)) \approx \ell(I)$.

With these observations in mind, we consider three cases. Set $R:=\left|X_{1}-X_{2}\right|$.
CASE 1. - $Q_{1} \subseteq Q_{2}$. In this case, $R \lesssim \ell\left(Q_{2}\right)$ (with $R \approx \ell\left(Q_{2}\right)$ if $\ell\left(Q_{2}\right) \gg \ell\left(Q_{1}\right)$ ), and $\min \left(\delta_{\star}\left(X_{1}\right), \delta_{\star}\left(X_{2}\right)\right) \gtrsim \ell\left(Q_{1}\right)$. Consequently, we may form a Harnack chain of cardinality $\approx k\left(Q_{1}\right)-k\left(Q_{2}\right)+1$ that connects the Corkscrew points of every $Q^{\prime}$, with $Q_{1} \subseteq Q^{\prime} \subseteq Q_{2}$.

Before proceeding to the remaining cases, we observe that if Case 1 does not hold, then $Q_{1}$ and $Q_{2}$ are disjoint, whence it follows from (A.1) and observations (1) and (5) above that

$$
\begin{equation*}
R \gtrsim \delta_{\star}\left(X_{2}\right) \approx \ell\left(Q_{2}\right) \geq \ell\left(Q_{1}\right) \approx \delta_{\star}\left(X_{1}\right) \tag{A.2}
\end{equation*}
$$

Of course, we also have $R \lesssim \ell(Q)$.
CASE 2. $-Q_{1} \cap Q_{2}=\varnothing$, but have a common ancestor $Q^{*} \subseteq Q$, with $\ell\left(Q^{*}\right) \approx R$. We may then proceed as in Case 1, to construct respective Harnack chains, connecting each of $X_{1}$ and $X_{2}$, to $X_{Q^{*}}$. The union of these two chains connects $X_{1}$ to $X_{2}$.

Case 3. - $Q_{1}$ and $Q_{2}$ have no common ancestor of length $\approx R$. In this case, we may suppose that $R<\ell(Q) /\left(M K_{0}\right)$, where $M$ is a sufficiently large number to be chosen momentarily. Indeed, if not, then $\ell(Q) /\left(M K_{0}\right) \leq R \lesssim \ell(Q)$, in which case $Q$ would be a common ancestor with $\ell(Q) \approx R$.

Thus, since $R<\ell(Q) /\left(M K_{0}\right)$, there exist $Q_{1}^{*}, Q_{2}^{*} \in \mathbb{D}_{Q}$ such that, for $i=1,2, Q_{i}^{*}$ is an ancestor of $Q_{i}$, with $\ell\left(Q_{1}^{*}\right)=\ell\left(Q_{2}^{*}\right) \approx M K_{0} R$. Since $\operatorname{dist}\left(X_{i}, Q_{i}\right) \lesssim K_{0} \ell\left(Q_{i}\right)$ by construction (cf. (3.49)), we then have that $\operatorname{dist}\left(Q_{1}, Q_{2}\right) \lesssim K_{0} R$ (by (A.2) and the triangle inequality), and therefore also that

$$
\operatorname{dist}\left(Q_{1}^{*}, Q_{2}^{*}\right) \leq C \ell\left(Q_{1}^{*}\right) / M \leq \ell\left(Q_{1}^{*}\right)=\ell\left(Q_{2}^{*}\right),
$$

by choice of $M$ large enough. Consequently, by (3.44), $W_{Q_{1}^{*}}^{*} \cap W_{Q_{2}^{*}}^{*}$ is non-empty, whence there is a Harnack chain connecting the respective Corkscrew points $X_{Q_{1}^{*}}$ and $X_{Q_{2}^{*}}$. We may then proceed as above to construct a Harnack chain from $X_{i}$ to $X_{Q_{i}^{*}}, i=1,2$, and the proof of the Harnack chain condition for the sawtooth $\Omega_{\mathcal{G}, Q}$ is now complete.

We finish this subsection by verifying the Harnack chain property for a Carleson box $T_{\Delta}$. Let $X_{1}, X_{2} \in T_{\Delta}$, with $a:=\delta_{\star}\left(X_{1}\right) \leq \delta_{\star}\left(X_{2}\right)=: b$. As above, we may suppose that $I_{1}^{*}$ and $I_{2}^{*}$ are separated, and thus as in observation (1), that each $X_{i}, i=1,2$ is the respective center of the Whitney cube $I_{i}$ whose dilate contains it.

By definition of $T_{\Delta}$ (cf. (3.58)-(3.59)), and since $X_{i}$ is the center of $I_{i}$, we have $X_{i} \in T_{Q^{i}}$, where $Q^{i} \in \mathbb{D}^{\Delta}, i=1,2$. By (3.44), $W_{Q^{1}} \cap W_{Q^{2}}$ is non-empty. Consequently, there is a Harnack chain connecting the respective Corkscrew points $X_{Q^{1}}$ and $X_{Q^{2}}$, so in the case that $\left|X_{1}-X_{2}\right|=: \Lambda a \approx r_{\Delta}$, we may connect $X_{1}$ to $X_{Q^{1}}$ to $X_{Q^{2}}$ to $X_{2}$.

Therefore, we may now suppose that $\left|X_{1}-X_{2}\right|=\Lambda a \leq r_{\Delta} /\left(M K_{0}\right)$, for some sufficiently large $M$ to be chosen momentarily. We note that there is a uniform constant $c>0$ such that $\Lambda \geq c$, since $X_{1}$ and $X_{2}$ are the respective centers of non-overlapping Whitney cubes
(cf. observation (5) above). We now claim that we also have $b \lesssim \Lambda a$. Indeed, if $b \gg \Lambda a$, then by the triangle inequality, $a \geq b-\Lambda a \gg \Lambda a$, which contradicts the uniform lower bound for $\Lambda$. Therefore, by observation (5) above, and by construction of each $T_{Q^{i}}$, there exist $\widetilde{Q}^{1} \subset Q^{1}, \widetilde{Q}^{2} \subset Q^{2}$ such that $X_{i} \in T_{\widetilde{Q}^{i}}, \ell\left(\widetilde{Q}^{1}\right)=\ell\left(\widetilde{Q}^{2}\right) \approx M K_{0} \Lambda a$, and

$$
\operatorname{dist}\left(X_{1}, \widetilde{Q}^{1}\right) \lesssim K_{0} a, \quad \operatorname{dist}\left(X_{2}, \widetilde{Q}^{2}\right) \lesssim K_{0} b \lesssim K_{0} \Lambda a
$$

By the triangle inequality, we then have

$$
\operatorname{dist}\left(\widetilde{Q}^{1}, \widetilde{Q}^{2}\right) \leq C \ell\left(\widetilde{Q}^{1}\right) / M \leq \ell\left(\widetilde{Q}^{1}\right)=\ell\left(\widetilde{Q}^{2}\right)
$$

by choice of $M$ large enough. By (3.44), $W_{\widetilde{Q}^{1}} \cap W_{\widetilde{Q}^{2}}$ is non-empty, so that we may construct a Harnack chain from $X_{1}$ to $X_{2}$, by a now familiar argument, via the Corkscrew points $X_{\widetilde{Q}^{1}}$ and $X_{\widetilde{Q}^{2}}$.

## A.3. ADR

Suppose that $\partial \Omega$ is ADR , and we show first that for each $Q \in \mathbb{D}(\partial \Omega)$, the boundary of the "Carleson box" $T_{Q}$ is also ADR. We begin with the upper bound. Let $x \in \partial T_{Q}$, and let $\Delta_{\star}:=\Delta_{\star}(x, r):=B(x, r) \cap \partial T_{Q}$, with $r \lesssim \operatorname{diam} Q$. If $B(x, r)$ meets $\partial \Omega$, then there is a point $x^{\prime} \in \partial \Omega$ such that $B(x, r) \subset B\left(x^{\prime}, 2 r\right)$. Consequently,

$$
H^{n}\left(B(x, r) \cap \partial \Omega \cap \partial T_{Q}\right) \leq H^{n}\left(\Delta\left(x^{\prime}, 2 r\right)\right) \lesssim r^{n}
$$

since $\partial \Omega$ is ADR.
Now consider $\Delta_{\star} \backslash \partial \Omega$. This portion of $\Delta_{\star}$ is contained in a union of faces (or partial faces) of fattened Whitney cubes $I^{*}=(1+\lambda) I$. Let $J_{Q}$ denote the collection of Whitney cubes $I$ for which $\partial I^{*}$ meets $\partial T_{Q}$, and $\operatorname{int}\left(I^{*}\right) \subset T_{Q}$. Suppose that $I \in \mathcal{I}_{Q}$ is a Whitney cube such that $\partial I^{*}$ meets $\Delta_{\star}$. Then

$$
\begin{equation*}
H^{n}\left(\Delta_{\star} \cap \partial I^{*}\right) \leq H^{n}\left(B(x, r) \cap \partial I^{*}\right) \lesssim \min \left(\ell(I)^{n}, r^{n}\right) \tag{A.3}
\end{equation*}
$$

Therefore,

$$
\sum_{I \in \mathcal{I}_{Q}: \ell(I) \geq r /\left(M K_{0}\right)} H^{n}\left(\Delta_{\star} \cap \partial I^{*}\right) \lesssim r^{n}
$$

because only a bounded number of terms can appear in this sum. Here, $M$ is a sufficiently large number to be chosen, and $K_{0}$ is the same constant appearing in (3.49). It remains to consider

$$
\sum_{I \in \mathcal{I}_{Q}: \ell(I)<r /\left(M K_{0}\right)} H^{n}\left(\Delta_{\star} \cap \partial I^{*}\right)=\sum_{k: 2^{-k}<r /\left(M K_{0}\right)} \sum_{I \in \mathcal{J}_{Q}^{k}} H^{n}\left(\Delta_{\star} \cap \partial I^{*}\right)
$$

where $\mathcal{J}_{Q}^{k}:=\left\{I \in J_{Q}: \ell(I)=2^{-k}\right\}$. It is then enough to show that there is an $\varepsilon>0$ such that for each $k$ with $2^{-k}<r /\left(M K_{0}\right)$, we have

$$
\begin{equation*}
\sum_{I \in \mathcal{J}_{Q}^{k}} H^{n}\left(\Delta_{\star} \cap \partial I^{*}\right) \lesssim 2^{-k \varepsilon} r^{n-\varepsilon} \tag{A.4}
\end{equation*}
$$

It follows from (A.3) that the latter bound will hold if the cardinality of the set of $I$ 's which make a non-trivial contribution to the sum is no larger than

$$
\begin{equation*}
C\left(2^{k} r\right)^{n-\varepsilon} \tag{A.5}
\end{equation*}
$$

We recall that by the definition of $T_{Q}$ (cf. (3.43)-(3.52)), for each $I \in \mathcal{J}_{Q}^{k}$, there is a $Q_{I} \in \mathbb{D}_{Q}$ such that $\ell\left(Q_{I}\right) \approx \ell(I)=2^{-k}$, and $\operatorname{dist}\left(I, Q_{I}\right) \lesssim K_{0} \ell(I)$. Since $2^{-k}<r$, there is a uniform constant $C$ such that $B(x, C r)$ contains each such $Q_{I}$, for every $I$ such that $\partial I^{*}$ meets $\Delta_{\star}$. We may then cover $B(x, C r) \cap Q$ by a bounded number of subcubes $Q^{\prime} \in \mathbb{D}_{Q}$, with $\ell\left(Q^{\prime}\right) \approx r$, so that each relevant $Q_{I}$ is contained in some $Q^{\prime}$. It is enough to consider those $Q_{I}$ contained in one such $Q^{\prime}$. We therefore now fix $Q^{\prime}$ and $k$, and distinguish two types of $Q_{I} \subset Q^{\prime}$ :

$$
\begin{array}{ll}
\text { Type 1 : } & \operatorname{dist}\left(Q_{I},\left(Q^{\prime}\right)^{c}\right)>2^{-\gamma k} r^{1-\gamma} \\
\text { Type 2 : } & \operatorname{dist}\left(Q_{I},\left(Q^{\prime}\right)^{c}\right) \leq 2^{-\gamma k} r^{1-\gamma}
\end{array}
$$

where we have fixed $\gamma \in(0,1)$. We note that there are at most a bounded number of $I$ 's corresponding to each $Q_{I}$. Thus, since $2^{-k}<r /\left(M K_{0}\right) \ll r$, by Lemma ?? (vi) we have that the cardinality of the set of $I$ 's for which $Q_{I}$ is of Type 2 is no larger than $C\left(2^{k} r\right)^{-\gamma \eta} r^{n} /\left(2^{-k n}\right) \approx$ $\left(2^{k} r\right)^{n-\gamma \eta}$, which is (A.5), with $\varepsilon=\gamma \eta$.

We now claim that for $M$ chosen large enough, depending on $\gamma$ and $K_{0}$, the collection of $I$ such that $\partial I^{*}$ meets $\Delta_{\star}$, and for which $Q_{I}$ is of Type 1 , is empty. Indeed, if $Q_{I}$ is of Type 1 , and if $M$ is sufficiently large, we then have

$$
\operatorname{dist}\left(Q_{I},\left(Q^{\prime}\right)^{c}\right)>2^{-\gamma k} r^{1-\gamma}>\left(M K_{0}\right)^{(1-\gamma)} 2^{-k} \gg K_{0} \ell(I) \gtrsim \operatorname{dist}\left(I, Q_{I}\right)
$$

Consequently, if $y \in \partial \Omega$ satisfies $\operatorname{dist}(I, y) \lesssim K_{0} \ell(I)$, then

$$
\operatorname{dist}\left(y, Q_{I}\right) \lesssim K_{0} \ell(I) \ll \operatorname{dist}\left(Q_{I},\left(Q^{\prime}\right)^{c}\right)
$$

so that $y \in Q^{\prime}$, and $\operatorname{dist}\left(y,\left(Q^{\prime}\right)^{c}\right) \gg K_{0} \ell(I)$. Now consider any Whitney cube $J \in \mathcal{W}$ that touches $I$. Then $\operatorname{dist}(J, \partial \Omega) \approx \ell(J) \approx \ell(I) \approx \operatorname{dist}(I, \partial \Omega)$, so that for some $y_{J} \in \partial \Omega$, we have $\operatorname{dist}\left(y_{J}, J\right) \approx \operatorname{dist}\left(y_{J}, I\right) \ll C_{0} \ell(I) \leq K_{0} \ell(I)$ (cf. (3.43) and (3.49).) Thus, $y_{J} \in Q^{\prime}$, and $\operatorname{dist}\left(y_{J},\left(Q^{\prime}\right)^{c}\right) \gg K_{0} \ell(I) \approx K_{0} \ell(J)$. It follows that there is a $Q_{J} \in \mathbb{D}(\partial \Omega)$, with $Q_{J} \subset Q^{\prime} \subset Q, y_{J} \in \bar{Q}_{J}, \ell\left(Q_{J}\right)=\ell(J)$, and $\operatorname{dist}\left(Q_{J}, J\right) \leq \operatorname{dist}\left(y_{J}, J\right) \leq C_{0} \ell(J)$. Therefore, $J \in \mathcal{W}_{Q_{J}}$. Since this is true for all Whitney cubes $J$ that touch $I$, we have in particular that every point on $\partial I^{*}$ is an interior point of $T_{Q}$, hence $\Delta_{\star} \cap \partial I^{*}=\varnothing$. We have now established the upper bound $H^{n}\left(\Delta_{\star}(x, r)\right) \lesssim r^{n}$.

The lower bound is easy. Consider $B:=B(x, r), r \lesssim \operatorname{diam} Q$, with $x \in \partial T_{Q}$. If $B \cap Q$ contains a surface ball $\Delta \subset \partial \Omega$, with radius $r_{\Delta} \gtrsim r$, then we are done, by the ADR property of $\partial \Omega$. Otherwise, if $B \cap Q$ contains no such surface ball, then $\operatorname{dist}(x, Q) \gtrsim r$, whence it follows that $x \in \partial I^{*}$, where $I$ is a Whitney cube with $\ell(I) \gtrsim r / K_{0}$ (cf. (3.49)), and where $x$ lies in a subset $F$ of a (closed) face of $I^{*}$, with $F \subset B \cap \partial T_{Q}$, and $H^{n}(F \cap B)=H^{n}(F) \gtrsim$ $\left(r / K_{0}\right)^{n}$, as desired.

Next, we discuss the ADR property of a Carleson region $T_{\Delta}$. By definition (cf. (3.59)), $T_{\Delta}$ is a union of a bounded number of regions $T_{Q}$. The upper bound in the ADR condition is then an immediate consequence of the corresponding bound for $T_{Q}$. The lower bound is proved in the same way as it was for $T_{Q}$ depending on whether or not the ball $B$ has an ample intersection with some $Q \in \mathbb{D}^{\Delta}$. We omit the routine details.

Finally, we establish the ADR property for the global (3.53) and local (3.54) sawtooth regions. The proofs are similar, so for the sake of specificity, we treat the global sawtooth $\Omega_{\mathcal{G}}$. We first prove the upper bound in the ADR condition. Fix $B:=B(x, r)$. The desired bound for $H^{n}\left(B \cap \partial \Omega \cap \partial \Omega_{\mathscr{G}}\right)$ is an immediate consequence of the fact that $\partial \Omega$ is ADR.
$4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - $\mathrm{N}^{\mathrm{o}} 3$

Now consider $\Sigma:=\partial \Omega_{\mathcal{G}} \backslash \partial \Omega$. We observe that this portion of the boundary consists of (portions of) faces of certain fattened Whitney cubes $J^{*}=(1+\lambda) J$, with $\operatorname{int}\left(J^{*}\right) \subset \Omega_{\mathcal{G}}$, which meet some $I \in W$ for which $I \notin \mathcal{W}_{Q}^{*}$, for any $Q \in \mathbb{D}_{\mathcal{G}}$ (so that $\tau I \subset \Omega \backslash \Omega_{\mathcal{G}}$ for some $\tau \in(1 / 2,1)$; cf. (3.51).) Necessarily, $I \in \mathcal{W}_{Q^{\prime}}^{*}$, where $Q^{\prime} \in \mathbb{D}_{Q_{j}}$ for some $Q_{j} \in \mathscr{F}$. For each $Q_{j} \in \mathcal{F}$, we set

$$
\mathcal{R}_{Q_{j}}:=\bigcup_{Q^{\prime} \in \mathbb{D}_{Q_{j}}} w_{Q^{\prime}}^{*},
$$

and denote by $\mathscr{F}_{B}$ the sub-collection of those $Q_{j} \in \mathscr{F}$ such that there is an $I \in \mathcal{R}_{Q_{j}}$ for which $B \cap \Sigma$ meets $I$. We then split the latter collection into $\mathscr{F}_{B}=\mathscr{F}_{1} \cup \mathscr{F}_{2}$, where $Q_{j} \in \mathscr{F}_{1}$ if $\ell\left(Q_{j}\right)<r$, and $Q_{j} \in \mathcal{F}_{2}$ if $\ell\left(Q_{j}\right) \geq r$. We consider the contribution of the latter first. Suppose that $Q_{j}$ and $Q_{k}$ are both in $\mathcal{F}_{2}$, and without loss of generality that $r \leq \ell\left(Q_{j}\right) \leq \ell\left(Q_{k}\right)$. Since $B$ meets some $I \in \mathcal{R}_{Q_{j}}$, we obtain in particular that $\operatorname{dist}(y, \partial \Omega) \lesssim \ell\left(Q_{j}\right), \forall y \in B$. Thus, $B \cap \Sigma$ lies within $C \ell\left(Q_{j}\right)$ of $\partial \Omega$, and therefore meets no Whitney cubes of side length greater than $C \ell\left(Q_{j}\right)$. Consequently, any such Whitney cube $I^{\prime} \in \mathcal{R}_{Q_{k}}$, which meets $B$, must lie within $C K_{0} \ell\left(Q_{j}\right)$ of $Q_{k}$. Therefore, for any pair $Q_{j}, Q_{k} \in \mathcal{F}_{2}$, we have that $\operatorname{dist}\left(Q_{j}, Q_{k}\right) \lesssim \min \left(\ell\left(Q_{j}\right), \ell\left(Q_{k}\right)\right)$ (with implicit constants depending on $K_{0}$.) Since the cubes in $\mathcal{F}$ are pairwise disjoint, it follows that the cardinality of $\mathscr{F}_{2}$ is uniformly bounded, hence

$$
H^{n}\left(B \cap \Sigma \cap\left(\bigcup_{Q_{j} \in \mathcal{F}_{2}} \cup_{I \in \mathscr{R}_{Q_{j}}} I\right)\right) \lesssim \sup _{Q_{j} \in \mathcal{F}} H^{n}\left(B \cap \Sigma_{j}\right),
$$

where $\Sigma_{j}:=\Sigma \cap\left(\bigcup_{I \in \mathscr{R}_{Q_{j}}} I\right)$. The desired bound for the contribution of $\mathscr{F}_{2}$ is an immediate consequence of following estimate, which holds for every $Q_{j} \in \mathcal{F}$ :

$$
\begin{equation*}
H^{n}\left(B \cap \Sigma_{j}\right) \lesssim\left(\min \left(r, \ell\left(Q_{j}\right)\right)\right)^{n} \tag{A.6}
\end{equation*}
$$

Let us take the latter bound for granted momentarily, and consider the contribution of $\mathscr{F}_{1}$. If $Q_{j} \in \mathscr{F}_{1}$, then $Q_{j} \subset B^{*}:=C K_{0} B$ for some uniform constant $C$. By the case $r>\ell\left(Q_{j}\right)$ of (A.6), we have that $H^{n}\left(\Sigma_{j} \cap B\right) \lesssim H^{n}\left(Q_{j}\right)$. Therefore,

$$
H^{n}\left(B \cap\left(\bigcup_{Q_{j} \in \mathcal{F}_{1}} \Sigma_{j}\right)\right) \lesssim \sum_{\mathscr{G}_{1}} H^{n}\left(Q_{j}\right) \leq H^{n}\left(B^{*} \cap \partial \Omega\right) \approx\left(K_{0} r\right)^{n},
$$

since the $Q_{j}$ 's are pairwise disjoint.
Thus, to finish proving the upper ADR bound for the sawtooth regions, it remains only to establish (A.6). Suppose first that $\ell\left(Q_{j}\right) \lesssim r$. We write

$$
\Sigma_{j}:=\bigcup_{k: 2^{-k} \leq \ell\left(Q_{j}\right)} \Sigma_{j}^{k}
$$

where $\Sigma_{j}^{k}=\Sigma \cap\left(\bigcup_{\left\{I \in \mathscr{R}_{Q_{j}}: \ell(I)=2^{-k}\right\}} I\right)=\Sigma_{j} \cap\left(\bigcup_{\left\{I \in \mathscr{R}_{Q_{j}}: \ell(I)=2^{-k}\right\}} I\right)$. We observe that for any $I \in \mathcal{W}$,

$$
\begin{equation*}
H^{n}(\Sigma \cap I) \lesssim \ell(I)^{n} . \tag{A.7}
\end{equation*}
$$

Moreover, there are at most a bounded number of $I \in \mathcal{R}_{Q_{j}}$ for which $\ell(I) \approx \ell\left(Q_{j}\right)$, so that

$$
H^{n}\left(\sum_{k: 2^{-k} \approx \ell\left(Q_{j}\right)} \Sigma_{j}^{k}\right) \lesssim \ell\left(Q_{j}\right)^{n}
$$

as desired. On the other hand, suppose $I \in \mathcal{R}_{Q_{j}}$, with $\ell(I)=2^{-k} \ll \ell\left(Q_{j}\right)$. Then there is a $Q_{I} \in \mathbb{D}_{Q_{j}}$, with $I \in \mathcal{W}_{Q_{I}}^{*}$. In addition, if $I$ meets $\Sigma_{j}$, then $I$ meets $J^{*}$, for some $J \in \mathcal{W}_{Q^{\prime}}^{*}$, with $Q^{\prime} \in \mathbb{D}_{\mathcal{F}}$, and $\ell\left(Q^{\prime}\right) \approx \ell(J) \approx \ell(I) \ll \ell\left(Q_{j}\right)$. We note that $Q^{\prime} \cap Q_{j}=\varnothing$, by definition of $\mathbb{D}_{\mathcal{F}}$, and the fact that $\ell\left(Q^{\prime}\right)<\ell\left(Q_{j}\right)$. Consequently,

$$
\operatorname{dist}\left(Q_{I},\left(Q_{j}\right)^{c}\right) \leq \operatorname{dist}\left(Q_{I}, Q^{\prime}\right) \lesssim K_{0} 2^{-k}
$$

Notice that for each such $Q_{I}$, there are at most a bounded number of $I^{\prime} \in \mathcal{W}_{Q_{I}}^{*}$ (indeed, by definition of $W_{Q_{I}}^{*}$, all such $I^{\prime}$ satisfy $\ell\left(I^{\prime}\right) \approx \ell\left(Q_{I}\right) \approx \operatorname{dist}\left(I^{\prime}, Q_{I}\right)$ ). By Lemma ?? (vi) we therefore have that

$$
\#\left\{I \in \mathscr{R}_{Q_{j}}: \ell(I)=2^{-k}, I \cap \Sigma_{j}^{k} \neq \varnothing\right\} \lesssim\left(2^{k} \ell\left(Q_{j}\right)\right)^{n-\eta}
$$

(where the implicit constant depends upon $K_{0}$ ), whence it follows that

$$
H^{n}\left(\sum_{k: 2^{-k} \ll \ell\left(Q_{j}\right)} \Sigma_{j}^{k}\right) \lesssim \ell\left(Q_{j}\right)^{n}
$$

Thus (A.6) holds in the case $r \gtrsim \ell\left(Q_{j}\right)$.
Now suppose that $r \ll \ell\left(Q_{j}\right)$. If $x \notin \partial \Omega$, then $B=B(x, r)$ is centered on a face of some $J_{x}^{*}$, with $\operatorname{int}\left(J_{x}^{*}\right) \subset \Omega_{g}$. If $\ell\left(J_{x}\right) \gg r$, we are done, by the nature of Whitney cubes. On the other hand, if $\ell\left(J_{x}\right) \lesssim r$, or if $x \in \partial \Omega$, then for each $I \in \mathcal{R}_{Q_{j}}$ which meets $B$, we have that $\ell(I) \lesssim r$, and also that $B(x, C r)$ meets $Q_{I}$, for some uniform constant $C$, where $Q_{I} \in \mathbb{D}_{Q_{j}}$ is defined as in the previous paragraph. We may then cover $B(x, C r) \cap Q_{j}$ by a bounded number of subcubes $Q^{i} \subset Q_{j}$, with $\ell\left(Q^{i}\right) \approx M r$, so that each relevant $Q_{I}$ is contained in some $Q^{i}$. Here, $M$ is a sufficiently large number, to be fixed momentarily. Now suppose that $I$ meets $\Sigma_{j}$. We may then proceed as in the previous paragraph, except that in this case we consider $\operatorname{dist}\left(Q_{I},\left(Q^{i}\right)^{c}\right)$, and $\ell\left(Q_{j}\right)$ is replaced by $\ell\left(Q^{i}\right) \approx M r$. As above, we find that $I$ meets some $J^{*}$, with $J \in \mathcal{W}_{Q^{\prime}}^{*}$ and $Q^{\prime} \in \mathbb{D}_{\mathcal{F}}$, so that $\ell\left(Q^{\prime}\right) \approx \ell(J) \approx \ell(I)$. In the present scenario, we have $\ell(I) \lesssim r$, therefore $\ell\left(Q^{\prime}\right)<\ell\left(Q^{i}\right)$, for $M$ chosen large enough and consequently $Q^{\prime} \cap Q^{i}=\varnothing$. The rest of the argument follows as before. We omit the details.

Finally, to complete our discussion of the ADR property, it remains only to prove the lower ADR bound for the sawtooth regions. For the sake of specificity, we treat only the case of a local sawtooth, as the proof in the global case is similar.

Fix now $Q_{0} \in \mathbb{D}, r \lesssim \operatorname{diam} Q_{0}$ and $x \in \partial \Omega_{\mathcal{G}, Q_{0}}$, where $\mathcal{F} \subset \mathbb{D}$ is a disjoint family, and set $B:=B(x, r)$ and $\Delta_{\star}=\Delta_{\star}(x, r):=B \cap \partial \Omega_{\mathscr{F}, Q_{0}}$. We consider two main cases. As usual, $M$ denotes a sufficiently large number to be chosen.

CASE 1: $\delta(x) \geq r /\left(M K_{0}\right)$. In this case, for some $J$ with $\operatorname{int}\left(J^{*}\right) \subset \Omega_{\mathcal{G}, Q_{0}}$, we have that $x$ lies on a subset $F$ of a (closed) face of $J^{*}$, satisfying $H^{n}(F) \gtrsim\left(r /\left(M K_{0}\right)\right)^{n}$, and $F \subset \partial \Omega_{\mathcal{F}, Q_{0}}$. Thus, $H^{n}\left(B \cap \partial \Omega_{\mathscr{G}, Q_{0}}\right) \geq H^{n}(B \cap F) \gtrsim\left(r /\left(M K_{0}\right)\right)^{n}$, as desired.

CASE 2: $\delta(x)<r /\left(M K_{0}\right)$. In this case, we have that $\operatorname{dist}\left(x, Q_{0}\right) \lesssim r / M$. Indeed, if $x \in \partial \Omega \cap \partial \Omega_{\mathcal{G}, Q_{0}}$, then by Proposition $6.1, x \in \overline{Q_{0}}$, so that $\operatorname{dist}\left(x, Q_{0}\right)=0$. Otherwise,
there is some cube $Q \in \mathbb{D}_{\mathcal{G}, Q_{0}}$ such that $x$ lies on the face of a fattened Whitney cube $I^{*}$, with $I \in \mathcal{W}_{Q}^{*}$, and $\ell(Q) \approx \ell(I) \approx \delta(x)<r /\left(M K_{0}\right)$. Thus,

$$
\operatorname{dist}\left(x, Q_{0}\right) \lesssim \operatorname{dist}(I, Q) \lesssim K_{0} \ell(Q) \lesssim r / M .
$$

Consequently, we may choose $\hat{x} \in Q_{0}$ such that $|x-\hat{x}| \lesssim r / M$. Fix now $\widehat{Q} \in \mathbb{D}_{Q_{0}}$ with $\hat{x} \in \widehat{Q}$ and $\ell(\widehat{Q}) \approx r / M$. Then for $M$ chosen large enough we have that $\widehat{Q} \subset B(\hat{x}, r / \sqrt{M}) \subset$ $B(x, r)$. We now consider two sub-cases.

Sub-case 2A: $B(\hat{x}, r / \sqrt{M})$ meets a $Q_{j} \in \mathscr{F}$ with $\ell\left(Q_{j}\right) \geq r / M$. Then in particular, there is a $Q \subseteq Q_{j}$, with $\ell(Q) \approx r / M$, and $Q \subset B(\hat{x}, 2 r / \sqrt{M})$. By Lemma 5.9, there is a ball $B^{\prime} \subset \mathbb{R}^{n+1} \backslash \overline{\Omega_{\mathcal{G}, Q_{0}}}$, with radius $r^{\prime} \approx \ell(Q) / K_{0} \approx r /\left(K_{0} M\right)$, such that $B^{\prime} \cap \partial \Omega \subset Q$, and thus also $B^{\prime} \subset B$ (for $M$ large enough). On the other hand, we have already established above that $\Omega_{\mathscr{G}, Q_{0}}$ satisfies the (interior) Corkscrew condition, so there is another ball $B^{\prime \prime} \subset B \cap \Omega_{\mathcal{G}, Q_{0}}$, with radius $r^{\prime \prime} \approx r$. Therefore, by the isoperimetric inequality and the structure theorem for sets of locally finite perimeter (cf. [20], pp. 190 and 205, resp.) we have $H^{n}\left(\Delta_{\star}\right) \gtrsim c_{K_{0}} r^{n}$.

Sub-case 2b: There is no $Q_{j}$ as in sub-case 2a. Thus, if $Q_{j} \in \mathcal{F}$ meets $B(\hat{x}, r / \sqrt{M})$, then $\ell\left(Q_{j}\right) \leq r / M$. Since $\hat{x} \in Q_{0}$, there is a surface ball

$$
\Delta_{1}:=\Delta\left(x_{1}, c r / \sqrt{M}\right) \subset Q_{0} \cap B(\hat{x}, r / \sqrt{M}) \subset Q_{0} \cap B .
$$

Let $\mathcal{F}_{1}$ denote the collection of those $Q_{j} \in \mathcal{F}$ which meet $\Delta_{1}$. We then have the covering

$$
\Delta_{1} \subset\left(\bigcup_{\mathscr{G}_{1}} Q_{j}\right) \cup\left(\Delta_{1} \backslash\left(\bigcup_{\mathscr{G}_{1}} Q_{j}\right)\right)
$$

If

$$
\begin{equation*}
\sigma\left(\frac{1}{2} \Delta_{1} \backslash\left(\bigcup_{\mathscr{F}_{1}} Q_{j}\right)\right) \geq \frac{1}{2} \sigma\left(\frac{1}{2} \Delta_{1}\right) \approx r^{n} \tag{A.8}
\end{equation*}
$$

then we are done, since $\Delta_{1} \backslash\left(\bigcup_{\mathcal{F}_{1}} Q_{j}\right) \subset\left(Q_{0} \backslash\left(\bigcup_{\mathcal{G}} Q_{j}\right)\right) \cap B \subset \Delta_{\star}$, by Proposition 6.1.
Otherwise, if (A.8) fails, then

$$
\begin{equation*}
\sum_{Q_{j} \in \mathscr{F}_{1}^{\prime}} \sigma\left(Q_{j}\right) \gtrsim r^{n} \tag{A.9}
\end{equation*}
$$

where $\mathscr{F}_{1}^{\prime}$ denotes those $Q_{j} \in \mathcal{F}_{1}$ which meet $\frac{1}{2} \Delta_{1}$. Let us remind the reader that $x_{j}^{\star}$ is the center of the $n$-dimensional cube $P_{j}$ constructed in Proposition 6.7, and we recall (6.10) and the related discussion. We claim that there is a uniform constant $C$ such that for each such $Q_{j}$, the ball $B_{Q_{j}}^{*}:=B\left(x_{j}^{\star}, C K_{0} \ell\left(Q_{j}\right)\right)$ contains both an interior and an exterior Corkscrew point for $\Omega_{\mathcal{G}, Q_{0}}$, with respect to the surface ball $B_{Q_{j}}^{*} \cap \partial \Omega_{\mathscr{F}, Q_{0}}$ (with Corkscrew constants that may depend upon $K_{0}$ ).

Indeed, the exterior point exists by virtue of Lemma 5.9 , while the interior point may be taken to be the center of some $I \in \mathcal{W}_{\widetilde{Q}_{j}}^{*}$ with $\ell(I) \approx \ell\left(\widetilde{Q}_{j}\right)$, where $\widetilde{Q}_{j}$ is the dyadic parent of $Q_{j}$, so that $\widetilde{Q}_{j} \in \mathbb{D}_{\mathcal{G}, Q_{0}}$ and therefore $I \subset \operatorname{int}\left(I^{*}\right) \subset \Omega_{\mathscr{G}, Q_{0}}$. Consequently, by the isoperimetric inequality and the structure theorem for sets of locally finite perimeter, we have

$$
\begin{equation*}
H^{n}\left(B_{Q_{j}}^{*} \cap \partial \Omega_{\mathcal{F}, Q_{0}}\right) \gtrsim \ell\left(Q_{j}\right)^{n} \approx \sigma\left(Q_{j}\right) . \tag{A.10}
\end{equation*}
$$

Now, by the ADR property and a covering lemma argument, and (A.9), there is a subcollection $\mathscr{F}_{1}^{\prime \prime} \subset \mathscr{F}_{1}^{\prime}$ such that the balls in $\left\{B_{Q_{j}}^{*}\right\}_{Q_{j} \in \mathcal{F}_{1}^{\prime \prime}}$ are pairwise disjoint and

$$
\begin{equation*}
\sum_{Q_{j} \in \mathcal{T}_{1}^{\prime \prime}} \sigma\left(Q_{j}\right) \gtrsim r^{n} \tag{A.11}
\end{equation*}
$$

Combining (A.10) and (A.11), we obtain that $H^{n}\left(\Delta_{\star}\right) \gtrsim r^{n}$, since for $M$ large enough, each $B_{Q_{j}}^{*} \subset B$, by construction.

## A.4. UR

In this subsection, we show that the Carleson box $T_{Q}$ inherits the UR property from $\Omega$. This fact extends routinely to any $T_{\Delta}$, and we omit the details.

Let us note that, since $\partial \Omega$ is UR, we have the global $L^{2}$ bound

$$
\begin{equation*}
\iint_{\mathbb{R}^{n+1}}\left|\nabla^{2} \delta f(X)\right|^{2} \delta(X) d X \leq C\|f\|_{L^{2}(\partial \Omega)}^{2} \tag{A.12}
\end{equation*}
$$

which is equivalent to the Carleson measure condition (1.10) by " $T 1$ reasoning".
Fix now $Q \in \mathbb{D}(\partial \Omega)$, and as usual let $\delta_{\star}(X):=\operatorname{dist}\left(X, \partial T_{Q}\right)$ (in the present context, $X$ need not belong to $T_{Q}$, but of course $\delta_{\star}(X)$ is still well-defined). By "local $T b$ " theory (see [24] in the present context), it is enough to verify that for every $\Delta_{\star}=\Delta_{\star}(x, r):=$ $B(x, r) \cap \partial T_{Q}$, with $x \in \partial T_{Q}$ and $r \lesssim \operatorname{diam}(Q)$, there is a function $b_{\Delta_{\star}}$, supported in $\Delta_{\star}$, and satisfying

$$
\begin{align*}
\left|f_{\Delta_{\star}} b_{\Delta_{\star}} d H^{n}\right| & \geq \frac{1}{C}  \tag{A.13}\\
f_{\Delta_{\star}}\left|b_{\Delta_{\star}}\right|^{2} d H^{n} & \leq C  \tag{A.14}\\
\iint_{B(x, 2 r)}\left|\nabla^{2} \& b_{\Delta_{\star}}(X)\right|^{2} \delta_{\star}(X) d X & \leq C r^{n} \tag{A.15}
\end{align*}
$$

where $C$ is a uniform constant, independent of $Q$. We fix a large constant $M$ to be chosen. There are two cases:

CASE 1: $\operatorname{dist}(x, \partial \Omega) \leq r /\left(M K_{0}\right)$.
In this case, either $x \in \bar{Q}$, or $x$ lies on a face of some $I^{*}=(1+\lambda) I$, with $I \in \mathcal{W}_{Q_{I}}^{*}$, for some $Q_{I} \in \mathbb{D}_{Q}$, where $\ell\left(Q_{I}\right) \approx \ell(I) \lesssim r /\left(M K_{0}\right)$, and $\operatorname{dist}\left(Q_{I}, I\right) \lesssim K_{0} \ell(I) \lesssim r / M$.

We claim that there is a surface ball $\Delta^{\prime}=B^{\prime} \cap \partial \Omega \subset \Delta_{\star} \cap \partial \Omega$, with $r_{\Delta^{\prime}} \approx r / M$, and with $B^{\prime} \cap \Omega \subset T_{Q}$. Indeed, if $x \in \bar{Q}$, then there is a $Q^{\prime} \in \mathbb{D}_{Q}$ such that $x \in \overline{Q^{\prime}}$ and $\ell\left(Q^{\prime}\right) \approx r / M$, and we may then set $\Delta^{\prime}:=B^{\prime} \cap \partial \Omega$, where $B^{\prime}:=B_{Q^{\prime}}^{\prime}$ is the ball promised by Lemma 3.55, applied with $Q^{\prime}$ in place of $Q$, so that $B^{\prime} \cap \Omega \subset T_{Q^{\prime}} \subset T_{Q}$. On the other hand, if $x \in \partial I^{*}$, with $I \in \mathcal{W}_{Q_{I}}^{*}$ as above, then there is a $Q^{\prime} \in \mathbb{D}_{Q}$ with $Q_{I} \subseteq Q^{\prime}$, and $\ell\left(Q^{\prime}\right) \approx r / M$. Moreover, by the triangle inequality, $|x-y| \lesssim r / M$, for every $y \in Q^{\prime}$, so that for $M$ large enough we have $Q^{\prime} \subset B(x, r) \cap \partial T_{Q}=\Delta_{\star}$ by Proposition 6.1. Thus, we may again set $B^{\prime}:=B_{Q^{\prime}}^{\prime}$, as in Lemma 3.55, and the claim is established.

We fix $\Delta^{\prime}=\Delta\left(x_{Q^{\prime}}, r_{\Delta^{\prime}}\right)$ as in the previous paragraph, and then set $b_{\Delta_{\star}}:=1_{\Delta^{\prime \prime}}$, where $\Delta^{\prime \prime}:=\Delta\left(x_{Q^{\prime}}, r_{\Delta^{\prime}} / 4\right)$. Then (A.14) is trivial, and (A.13) holds by the ADR properties
of $\Omega$ and $T_{Q}$. It remains to establish (A.15). To this end, we claim that $\delta_{\star}(X)=\delta(X)$, for $X \in 2 B^{\prime \prime}=\frac{1}{2} B^{\prime}$. Momentarily taking this claim for granted, we obtain that

$$
\iint_{2 B^{\prime \prime}}\left|\nabla^{2} \& b_{\Delta_{\star}}(X)\right|^{2} \delta_{\star}(X) d X \leq C r^{n}
$$

by (A.12), since $b_{\Delta_{\star}}=1_{\Delta^{\prime \prime}}$ is supported in $\partial \Omega$. Otherwise, for $X \in B(x, 2 r) \backslash 2 B^{\prime \prime}$, we have

$$
\begin{equation*}
\left|\nabla^{2} \delta b_{\Delta_{\star}}(X)\right| \lesssim\left|\int_{|X-y| \gtrsim r_{\Delta^{\prime}}}\right| X-\left.y\right|^{-n-1} 1_{\Delta^{\prime \prime}}(y) d \sigma(y) \mid \lesssim 1 / r_{\Delta^{\prime}} \approx M / r \tag{A.16}
\end{equation*}
$$

by the ADR property, and (A.15) follows.
Let us now verify the claim. Fix $X \in \frac{1}{2} B^{\prime}$. We note that

$$
\operatorname{dist}\left(X, \partial \Omega \cap \partial T_{Q}\right) \leq \frac{1}{2} r_{\Delta^{\prime}}
$$

since $B^{\prime}$ is centered on $\partial \Omega \cap \partial T_{Q}$. On the other hand, $\partial T_{Q} \backslash \partial \Omega \subset \Omega$, and therefore lies outside of $B^{\prime}$, since, by construction, $B^{\prime} \cap \Omega \subset T_{Q}$. Thus,

$$
\operatorname{dist}\left(X, \partial T_{Q} \backslash \partial \Omega\right) \geq \frac{1}{2} r_{\Delta^{\prime}} .
$$

Consequently, $\delta_{\star}(X)=\operatorname{dist}\left(X, \partial \Omega \cap \partial T_{Q}\right)$. Similarly, we shall have that $\delta(X)=$ $\operatorname{dist}\left(X, \partial \Omega \cap \partial T_{Q}\right)$, and thus $\delta(X)=\delta_{\star}(X)$ as claimed, once we show that $\partial \Omega \backslash \partial T_{Q}$ lies outside $B^{\prime}$, or equivalently, that $\partial \Omega \cap B^{\prime} \subset \partial T_{Q}$. So, fix $y \in \partial \Omega \cap B^{\prime}$. Since $B^{\prime}$ is open, we have that $B\left(y, \varepsilon_{0}\right) \subset B^{\prime}$ for $\varepsilon_{0}$ small enough. Note that $B\left(y, \varepsilon_{k}\right)$ meets $\Omega$ for a sequence $\varepsilon_{k} \rightarrow 0$, with $\varepsilon_{k}<\varepsilon_{0}$. Thus, there exists a sequence $\left\{y_{k}\right\}_{k} \subset B^{\prime} \cap \Omega \subset T_{Q}$, with $y_{k} \rightarrow y$, whence $y \in \partial T_{Q}$.

CASE 2: $\operatorname{dist}(x, \partial \Omega)>r /\left(M K_{0}\right)$.
In this case, we can find a Whitney cube $I$ with $x \in \partial I^{*}$ and $\operatorname{int}\left(I^{*}\right) \subset T_{Q}$, and a ball $B^{\prime}=B\left(x^{\prime}, r^{\prime}\right)$, with $r^{\prime} \approx r /\left(M K_{0}\right)$, such that some face $F$ of $I^{*}$ contains the surface ball $\Delta_{\star}^{\prime}:=B^{\prime} \cap \partial T_{Q}$. We define $b_{\Delta_{\star}}:=1_{\Delta_{\star}^{\prime \prime}}$, where $\Delta_{\star}^{\prime \prime}=B^{\prime \prime} \cap \partial T_{Q}$ and $B^{\prime \prime}:=\frac{1}{4} B^{\prime}$. We may now proceed as in Case 1 , using that of course (A.12) holds when $\partial \Omega$ is replaced by the hyper-plane $\mathscr{H}$ that contains $F$, and $\delta(X)=\operatorname{dist}(X, \mathscr{H})$. We omit the routine details.

## Appendix B

## Dyadically doubling and Muckenhoupt weights

Recall that, for a fixed cube $Q_{0} \in \mathbb{D}$, we say that $\omega$ is dyadically doubling on $Q_{0}$ if there exists $C_{\omega}$ such that $\omega(Q) \leq C_{\omega} \omega\left(Q^{\prime}\right)<\infty$ for every $Q \in \mathbb{D}_{Q_{0}}$, and for every dyadic "child" $Q^{\prime}$ of $Q$. We write $C_{\sigma}$ for the dyadic doubling constant of $\sigma$ (which depends on the ADR property). Throughout Appendix B, $Q_{0}$ will denote a fixed cube in $\mathbb{D}$. Let us also recall that the projection operators $\mathscr{P}_{\mathscr{G}}$ have been introduced in Section 6 .

Lemma B.1. - Fix $Q_{0}$. Let $\omega$ be a dyadically doubling measure on $Q_{0}$ with constant $C_{\omega}$. Then for every family $\mathcal{F} \subset \mathbb{D}_{Q_{0}}$ of pairwise disjoint dyadic cubes, $\mathscr{P}_{\mathscr{G}} \omega$ is dyadically doubling on $Q_{0}$, indeed $\mathscr{P}_{\mathscr{g}} \omega(Q) \leq \max \left(C_{\omega}, C_{\sigma}\right) \mathscr{P}_{g} \omega\left(Q^{\prime}\right)$ for every $Q \in \mathbb{D}_{Q_{0}}$, and for every dyadic "child" $Q^{\prime}$ of $Q$.

Proof. - We follow the proof of [27, Lemma B.1]. Let us fix $Q \in \mathbb{D}_{Q_{0}}$ and one of its dyadic "children" $Q^{\prime}$. We consider several cases.

Case 1: There exists $Q_{k} \in \mathcal{F}$ with $Q \subset Q_{k}$. The estimate is trivial in this case:

$$
\mathscr{P}_{\mathscr{F}} \omega(Q)=\frac{\sigma(Q)}{\sigma\left(Q_{k}\right)} \omega\left(Q_{k}\right) \leq C_{\sigma} \frac{\sigma\left(Q^{\prime}\right)}{\sigma\left(Q_{k}\right)} \omega\left(Q_{k}\right)=C_{\sigma} \mathscr{P}_{\mathscr{F}} \omega\left(Q^{\prime}\right)<\infty
$$

Case 2: $Q^{\prime} \in \mathcal{F}$. Notice that $\mathscr{P}_{\mathscr{F}} \omega\left(Q^{\prime}\right)=\omega\left(Q^{\prime}\right)$. Let $\mathcal{F}_{1}$ be the family of cubes $Q_{k} \in \mathcal{F}$ with $Q_{k} \cap Q \neq \varnothing$ and observe that if $Q_{k} \in \mathcal{F}_{1}$ then $Q_{k} \subsetneq Q$. Thus,

$$
\begin{aligned}
\mathscr{P}_{\mathscr{G}} \omega(Q) & =\omega\left(Q \backslash\left(\bigcup_{Q_{k} \in \mathscr{F}} Q_{k}\right)\right)+\sum_{Q_{k} \in \mathscr{F}_{1}} \frac{\sigma\left(Q_{k} \cap Q\right)}{\sigma\left(Q_{k}\right)} \omega\left(Q_{k}\right) \\
& =\omega\left(Q \backslash\left(\bigcup_{Q_{k} \in \mathscr{F}} Q_{k}\right)\right)+\sum_{Q_{k} \in \mathscr{F}_{1}} \omega\left(Q_{k}\right) \\
& =\omega(Q) \leq C_{\omega} \omega\left(Q^{\prime}\right)=C_{\omega} \mathscr{P}_{\mathscr{G}} \omega\left(Q^{\prime}\right)<\infty
\end{aligned}
$$

CASE 3: None of the conditions in the previous cases occur. We take the same set $\mathcal{F}_{1}$ and observe that if $Q_{k} \in \mathcal{F}_{1}$ then $Q_{k} \subsetneq Q$ (otherwise we are driven to Case 1). Let $\mathcal{F}_{2}$ be the family of cubes $Q_{k} \in \mathcal{F}$ with $Q_{k} \cap Q^{\prime} \neq \varnothing$. Notice that if $Q_{k} \in \mathcal{F}_{2}$ then $Q_{k} \subsetneq Q^{\prime}$ : otherwise, either $Q_{k}=Q^{\prime}$ which leads us to Case 2 , or $Q^{\prime} \subsetneq Q_{k}$ which implies $Q \subset Q_{k}$ and this is Case 1. Then proceeding as in the previous case one obtains that $\mathscr{P}_{\mathscr{G}} \omega(Q)=\omega(Q)$ and $\mathscr{P}_{\mathscr{G}} \omega\left(Q^{\prime}\right)=\omega\left(Q^{\prime}\right)$ which in turn imply

$$
\mathscr{P}_{\mathscr{F}} \omega(Q)=\omega(Q) \leq C_{\omega} \omega\left(Q^{\prime}\right)=C_{\omega} \mathscr{D}_{\mathscr{F}} \omega\left(Q^{\prime}\right)<\infty
$$

Lemma B.2. - Under the hypotheses of Lemma 6.15, $\nu$ and $\mathscr{P}_{\mathscr{G}} \nu$ are dyadically doubling on $Q_{0}$.

Proof. - We proceed as in [27, Lemma B.2]. Let us first consider $\nu$. Fix $Q \in \mathbb{D}_{Q_{0}}$, and one of its dyadic "children" $Q^{\prime}$. We recall Proposition 6.7, which for each $Q_{k} \in \mathcal{F}$ promises the existence of an $n$-dimensional cube $P_{k} \subset \partial \Omega_{\mathcal{F}, Q_{0}}$, with $\ell\left(P_{k}\right) \approx \ell\left(Q_{k}\right) \approx \operatorname{dist}\left(P_{k}, Q_{k}\right) \approx$ $\operatorname{dist}\left(P_{k}, \partial \Omega\right)$.

Case 1: There exists $Q_{k} \in \mathcal{F}$ with $Q \subset Q_{k}$. The estimate is trivial in this case since $\omega$ is dyadically doubling:

$$
\nu(Q)=\frac{\omega(Q)}{\omega\left(Q_{k}\right)} \omega_{\star}\left(P_{k}\right) \leq C_{\omega} \frac{\omega\left(Q^{\prime}\right)}{\omega\left(Q_{k}\right)} \omega_{\star}\left(P_{k}\right)=C_{\omega} \nu\left(Q^{\prime}\right)<\infty
$$

Case 2: $Q^{\prime} \in \mathcal{F}$. Write $Q^{\prime}=Q_{1} \in \mathcal{F}$. Notice that $\nu\left(Q^{\prime}\right)=\omega_{\star}\left(P_{1}\right)$. Let $\mathcal{F}_{1}$ be the family of cubes $Q_{k} \in \mathcal{F}$ with $Q_{k} \cap Q \neq \varnothing$ and observe that if $Q_{k} \in \mathcal{F}_{1}$ then $Q_{k} \subsetneq Q$. Thus,

## Remark 6.9 implies

$$
\begin{aligned}
\nu(Q) & =\omega_{\star}\left(Q \backslash\left(\bigcup_{Q_{k} \in \mathcal{F}} Q_{k}\right)\right)+\sum_{Q_{k} \in \mathcal{F}_{1}} \frac{\omega\left(Q_{k} \cap Q\right)}{\omega\left(Q_{k}\right)} \omega_{\star}\left(P_{k}\right) \\
& =\omega_{\star}\left(Q \backslash\left(\bigcup_{Q_{k} \in \mathcal{F}} Q_{k}\right)\right)+\sum_{Q_{k} \in \mathcal{F}_{1}} \omega_{\star}\left(P_{k}\right) \\
& \lesssim \omega_{\star}\left(\left(Q \backslash\left(\bigcup_{Q_{k} \in \mathcal{F}} Q_{k}\right)\right) \cup\left(\bigcup_{Q_{k} \in \mathcal{F}_{1}} P_{k}\right)\right) .
\end{aligned}
$$

We note that there are uniform positive constants $c$ and $C$ such that

$$
\begin{equation*}
\left(Q \backslash\left(\bigcup_{Q_{k} \in \mathcal{F}} Q_{k}\right)\right) \cup\left(\bigcup_{Q_{k} \in \mathscr{F}_{1}} P_{k}\right) \subset \Delta_{\star}\left(x_{1}^{\star}, C \ell\left(P_{1}\right)\right) \tag{B.3}
\end{equation*}
$$

and

$$
\begin{equation*}
\Delta_{\star}\left(x_{1}^{\star}, c \ell\left(P_{1}\right)\right) \subset P_{1}, \tag{B.4}
\end{equation*}
$$

where as usual $x_{1}^{\star}$ denotes the center of the $n$-dimensional cube $P_{1}$. Indeed, (B.4) is trivial, since by construction (cf. Proposition 6.7), $P_{1} \subset \partial \Omega_{\text {ๆ. } Q_{0}}$. To verify (B.3), it is enough to observe that, by Proposition 6.7, and the fact that $Q$ is the dyadic parent of $Q_{1}$, for $Q_{k} \in \mathcal{F}_{1}$ we have

$$
\ell\left(P_{k}\right) \approx \operatorname{dist}\left(P_{k}, Q_{k}\right) \approx \operatorname{dist}\left(P_{k}, Q\right) \lesssim \ell\left(P_{1}\right) \approx \ell(Q) \approx \operatorname{dist}\left(Q, P_{1}\right)
$$

Consequently, since $\omega_{\star}$ is doubling, we have

$$
\nu(Q) \lesssim \omega_{\star}\left(\Delta_{\star}\left(x_{1}^{\star}, C \ell\left(P_{1}\right)\right)\right) \lesssim \omega_{\star}\left(\Delta_{\star}\left(x_{1}^{\star}, c \ell\left(P_{1}\right)\right)\right) \leq \omega_{\star}\left(P_{1}\right)=\nu\left(Q^{\prime}\right) .
$$

Case 3: None of the conditions in the previous cases occur. We take the same set $\mathscr{F}_{1}$ and observe that if $Q_{k} \in \mathcal{F}_{1}$ then $Q_{k} \subsetneq Q$ (otherwise we are driven to Case 1). Let $\mathcal{F}_{2}$ be the family of cubes $Q_{k} \in \mathcal{F}$ with $Q_{k} \cap Q^{\prime} \neq \varnothing$. Notice that if $Q_{k} \in \mathcal{F}_{2}$ then $Q_{k} \subsetneq Q^{\prime}$ : otherwise, either $Q_{k}=Q^{\prime}$ which leads us to Case 2 , or $Q^{\prime} \subsetneq Q_{k}$ which implies $Q \subset Q_{k}$ and this is Case 1. We claim that for some uniform constant $C$, we have

$$
\begin{equation*}
\left(Q \backslash\left(\bigcup_{Q_{k} \in \mathcal{G}} Q_{k}\right)\right) \cup\left(\bigcup_{Q_{k} \in \mathcal{F}_{1}} P_{k}\right) \subset \Delta_{\star}\left(x_{Q^{\prime}}^{\star}, C t_{Q^{\prime}}\right) \tag{B.5}
\end{equation*}
$$

(see Proposition 6.12 for the notation). Indeed, since $Q$ is the dyadic parent of $Q^{\prime}$, by the construction in Proposition 6.12 (applied to $Q^{\prime}$ ), we have that

$$
\operatorname{dist}\left(x_{Q^{\prime}}^{\star}, Q\right) \leq \operatorname{dist}\left(x_{Q^{\prime}}^{\star}, Q^{\prime}\right) \lesssim \ell\left(Q^{\prime}\right) \approx \ell(Q) \approx t_{Q^{\prime}},
$$

whence (B.5) follows immediately.
Then we proceed as in the previous case and obtain that

$$
\begin{aligned}
\nu(Q) & \lesssim \omega_{\star}\left(\left(Q \backslash\left(\bigcup_{Q_{k} \in \mathcal{G}} Q_{k}\right)\right) \cup\left(\bigcup_{Q_{k} \in \mathcal{F}_{1}} P_{k}\right)\right) \\
& \leq \omega_{\star}\left(\Delta_{\star}\left(x_{Q^{\prime}}^{\star}, C t_{Q^{\prime}}\right)\right) \lesssim \omega_{\star}\left(\Delta_{\star}^{Q^{\prime}}\right) \lesssim \nu\left(Q^{\prime}\right)
\end{aligned}
$$

where we have used that $\omega_{\star}$ is doubling and where the last inequality follows as in (6.19):

$$
\nu\left(Q^{\prime}\right)=\omega_{\star}\left(Q^{\prime} \cap E_{0}\right)+\sum_{Q_{k} \in \mathcal{F}_{2}} \omega_{\star}\left(P_{k}\right) \gtrsim \omega_{\star}\left(\Delta_{\star}^{Q^{\prime}}\right) .
$$

One might show that $\mathscr{P}_{\mathscr{G}} \nu$ is dyadically doubling by invoking Lemma B.1, but then the doubling constant would depend on $\omega$ and $\omega_{\star}$. This is not the right approach as we have already observed that $\mathscr{P}_{\mathscr{g}} \nu$ does not depend on $\omega$. On the other hand, following the previous argument for $\nu$ we can see that the doubling constant does not depend on $\omega$. In Cases 2 and 3, we have that $\mathscr{P}_{\mathscr{g}} \nu(Q)=\nu(Q)$ and $\mathscr{P}_{\mathscr{G}} \nu\left(Q^{\prime}\right)=\nu\left(Q^{\prime}\right)$ so the doubling condition follows at once from the previous computations, and depends quantitatively only upon the doubling constant for $\omega_{\star}$, but not on $\omega$. In Case 1 we obtain

$$
\mathscr{P}_{\mathscr{G}} \nu(Q)=\frac{\sigma(Q)}{\sigma\left(Q_{k}\right)} \omega_{\star}\left(P_{k}\right) \leq C_{\sigma} \frac{\sigma\left(Q^{\prime}\right)}{\sigma\left(Q_{k}\right)} \omega_{\star}\left(P_{k}\right)=C_{\sigma} \mathscr{P}_{\mathscr{G}} \nu\left(Q^{\prime}\right) .
$$

Let us remind the reader that, as explained above, we may view $\partial \Omega_{\mathscr{G}, Q_{0}}$ itself as a surface ball $\Delta_{\star}^{Q_{0}}$ of radius $r\left(\Delta_{\star}^{Q_{0}}\right) \approx K_{0} \ell\left(Q_{0}\right)$, and then $A_{\infty}\left(\partial \Omega_{\mathcal{F}, Q_{0}}\right)$ is identified with $A_{\infty}\left(\Delta_{\star}^{Q_{0}}\right)$.

Lemma B.6. - Under the hypotheses of Lemma 6.15, if $\omega_{\star} \in A_{\infty}\left(\partial \Omega_{\mathcal{G}, Q_{0}}\right)$, then $\mathscr{P}_{\mathscr{G}} \nu \in A_{\infty}^{\text {dyadic }}\left(Q_{0}\right)$.

Proof. - Fix $0<\eta<1 / 2$ and $F \subset Q \in \mathbb{D}_{Q_{0}}$ with $\sigma(F) \geq(1-\eta) \sigma(Q)$.
Case 1: There exists $Q_{k} \in \mathcal{F}$ with $Q \subset Q_{k}$. The estimate is trivial in this case:

$$
\frac{\mathcal{P}_{\mathscr{G}} \nu(F)}{\mathcal{P}_{\mathscr{G}} \nu(Q)}=\frac{\frac{\sigma(F)}{\sigma\left(Q_{k}\right)} \omega_{\star}\left(P_{k}\right)}{\frac{\sigma(Q)}{\sigma\left(Q_{k}\right)} \omega_{\star}\left(P_{k}\right)}=\frac{\sigma(F)}{\sigma(Q)} \geq 1-\eta .
$$

Case 2: $Q$ is not contained in any $Q_{k} \in \mathscr{F}$ (i.e., $Q \in \mathbb{D}_{\mathcal{G}, Q_{0}}$ ). Let $\mathcal{F}_{1}$ be the family of cubes $Q_{k} \in \mathscr{F}$ with $Q_{k} \cap Q \neq \varnothing$ and observe that if $Q_{k} \in \mathcal{F}_{1}$ then $Q_{k} \subsetneq Q$. We set

$$
\tilde{\mathcal{F}}=\left\{Q_{k} \in \mathcal{F}_{1}: \sigma\left(F \cap Q_{k}\right) \geq(1-2 \eta) \sigma\left(Q_{k}\right)\right\},
$$

and

$$
E_{0}=Q \backslash \bigcup_{Q_{k} \in \mathcal{F}} Q_{k}, \quad G=\bigcup_{Q_{k} \in \tilde{\mathcal{F}}} Q_{k}, \quad B=\bigcup_{Q_{k} \in \mathcal{F}_{1} \backslash \tilde{\mathscr{G}}} Q_{k}
$$

Note that

$$
\sigma(F \cap B)=\sum_{Q_{k} \in \mathcal{F}_{1} \backslash \tilde{\mathscr{G}}} \sigma\left(F \cap Q_{k}\right) \leq(1-2 \eta) \sum_{Q_{k} \in \mathscr{F}_{1} \backslash \tilde{\mathscr{G}}} \sigma\left(Q_{k}\right) \leq(1-2 \eta) \sigma(Q) .
$$

Thus,

$$
\begin{aligned}
(1-\eta) \sigma(Q) & \leq \sigma(F) \leq \sigma\left(F \cap E_{0}\right)+\sigma(F \cap B)+\sigma(F \cap G) \\
& \leq \sigma\left(\left(F \cap E_{0}\right) \cup G\right)+(1-2 \eta) \sigma(Q),
\end{aligned}
$$

and therefore $\sigma\left(\left(F \cap E_{0}\right) \cup G\right) \geq \eta \sigma(Q)$.
Note that the ADR property of $\partial \Omega$ and $\partial \Omega_{\mathcal{G}, Q_{0}}$ implies

$$
\sigma(Q) \approx \ell(Q)^{n} \approx\left(\widehat{r}_{Q}\right)^{n} \approx \sigma_{\star}\left(\Delta_{\star}\left(y_{Q}, \widehat{r}_{Q}\right)\right)
$$

$4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - No 3
with $\Delta_{\star}\left(y_{Q}, \widehat{r}_{Q}\right)$ given in (6.11) (see also Proposition 6.4), where as usual we write $\sigma_{\star}$ to denote the "surface measure" on $\partial \Omega_{\mathcal{F}, Q_{0}}$, i.e., $\sigma_{\star}=\left.H^{n}\right|_{\partial \Omega_{\mathcal{G}, Q_{0}}}$. If we set $G_{\star}=\bigcup_{Q_{k} \in \tilde{\mathcal{G}}} P_{k}$ we have that $\sigma(G) \approx \sigma_{\star}\left(G_{\star}\right)$. Indeed, since $\Omega_{\mathscr{F}, Q_{0}}$ is ADR we have that

$$
\sigma_{\star}\left(P_{k}\right) \approx \ell\left(P_{k}\right)^{n} \approx \ell\left(Q_{k}\right)^{n} \approx \sigma\left(Q_{k}\right)
$$

by Proposition 6.7; thus Remark 6.9 yields

$$
\sigma(G)=\sum_{Q_{k} \in \tilde{\mathscr{F}}} \sigma\left(Q_{k}\right) \approx \sum_{Q_{k} \in \tilde{\mathscr{G}}} \sigma_{\star}\left(P_{k}\right) \approx \sigma_{\star}\left(G_{\star}\right) .
$$

On the other hand, Proposition 6.1 gives $\sigma\left(F \cap E_{0}\right)=\sigma_{\star}\left(F \cap E_{0}\right)$ and therefore

$$
\sigma_{\star}\left(\left(F \cap E_{0}\right) \cup G_{\star}\right) \approx \sigma\left(\left(F \cap E_{0}\right) \cup G\right) \geq \eta \sigma(Q) \approx \eta \sigma_{\star}\left(\Delta_{\star}\left(y_{Q}, \widehat{r}_{Q}\right)\right)
$$

Next we use that $\omega_{\star} \in A_{\infty}\left(\partial \Omega_{\mathcal{G}, Q_{0}}\right)$ to obtain

$$
\frac{\omega_{\star}\left(\left(F \cap E_{0}\right) \cup G_{\star}\right)}{\omega_{\star}\left(\Delta_{\star}\left(y_{Q}, \widehat{r}_{Q}\right)\right)} \gtrsim\left(\frac{\sigma_{\star}\left(\left(F \cap E_{0}\right) \cup G_{\star}\right)}{\sigma_{\star}\left(\Delta_{\star}\left(y_{Q}, \widehat{r}_{Q}\right)\right)}\right)^{\theta} \gtrsim \eta^{\theta}
$$

where we have used that $\left(F \cap E_{0}\right) \cup G_{\star} \subset \Delta_{\star}\left(y_{Q}, \widehat{r}_{Q}\right)$ by (6.11). Then,

$$
\begin{aligned}
\mathscr{P}_{\mathscr{G}} \nu(F) & \geq \omega_{\star}\left(F \cap E_{0}\right)+\sum_{Q_{k} \in \tilde{\mathcal{G}}} \frac{\sigma\left(F \cap Q_{k}\right)}{\sigma\left(Q_{k}\right)} \omega_{\star}\left(P_{k}\right) \\
& \geq \omega_{\star}\left(F \cap E_{0}\right)+(1-2 \eta) \sum_{Q_{k} \in \widetilde{\mathscr{G}}} \omega_{\star}\left(P_{k}\right) \\
& \geq(1-2 \eta) \omega_{\star}\left(\left(F \cap E_{0}\right) \cup G_{\star}\right) \\
& \gtrsim(1-2 \eta) \eta^{\theta} \omega_{\star}\left(\Delta_{\star}\left(y_{Q}, \widehat{r}_{Q}\right)\right) \\
& \gtrsim(1-2 \eta) \eta^{\theta} \omega_{\star}\left(\left(Q \cup\left(\bigcup_{Q_{k} \in \mathcal{F}: Q_{k} \subset Q} B\left(x_{k}^{\star}, r_{k}\right)\right)\right) \cap \partial \Omega_{\mathcal{G}, Q_{0}}\right)
\end{aligned}
$$

where the last inequality follows from (6.11). Next we observe that, by Proposition 6.1,

$$
\begin{aligned}
\left(Q \cap E_{0}\right) \cup\left(\bigcup_{Q_{k} \in \mathscr{F}_{1}} P_{k}\right) & \subset\left(Q \cap E_{0}\right) \cup\left(\bigcup_{Q_{k} \in \mathscr{F}_{1}} \Delta_{\star}\left(x_{k}^{\star}, r_{k}\right)\right) \\
& \subset\left(Q \cup\left(\bigcup_{Q_{k} \in \mathscr{F}: Q_{k} \subset Q} B\left(x_{k}^{\star}, r_{j}\right)\right)\right) \cap \partial \Omega_{\mathscr{F}, Q_{0}}
\end{aligned}
$$

Consequently,

$$
\begin{aligned}
\mathscr{P}_{\mathscr{G}} \nu(F) & \gtrsim(1-2 \eta) \eta^{\theta} \omega_{\star}\left(\left(Q \cap E_{0}\right) \cup\left(\bigcup_{Q_{k} \in \mathscr{F}_{1}} P_{k}\right)\right) \\
& \gtrsim(1-2 \eta) \eta^{\theta}\left(\omega_{\star}\left(Q \cap E_{0}\right)+\sum_{Q_{k} \in \mathcal{F}_{1}} \omega_{\star}\left(P_{k}\right)\right)=(1-2 \eta) \eta^{\theta} \mathscr{P}_{\mathscr{G}} \nu(Q) .
\end{aligned}
$$

Thus, in both cases we have shown as desired that $\mathscr{P}_{\mathscr{G}} \nu(F) / \mathscr{P}_{\mathscr{G}} \nu(Q) \geq C_{\eta}$.

Next we give a version of the classical result in [14] valid in our situation. The proof of this result follows the standard arguments in [21] although one has to adapt the ideas to the dyadic and local setting considered here. We give the proof for completeness.

Lemma B.7. - Let $Q_{0}$ be a fixed cube and let $\omega_{1}, \omega_{2}$ be two dyadically doubling measures on $Q_{0}$. Assume that there exist positive constants $C_{0}, \theta_{0}$ such that for all $Q \in \mathbb{D}_{Q_{0}}$ and $F \subset Q$,

$$
\begin{equation*}
\frac{\omega_{2}(F)}{\omega_{2}(Q)} \leq C_{0}\left(\frac{\omega_{1}(F)}{\omega_{1}(Q)}\right)^{\theta_{0}} \tag{B.8}
\end{equation*}
$$

Then, there exist positive constants $C_{1}, \theta_{1}$ such that for all $Q \in \mathbb{D}_{Q_{0}}$ and $F \subset Q$,

$$
\begin{equation*}
\frac{\omega_{1}(F)}{\omega_{1}(Q)} \leq C_{1}\left(\frac{\omega_{2}(F)}{\omega_{2}(Q)}\right)^{\theta_{1}} \tag{B.9}
\end{equation*}
$$

REmARK B.10. - The proof shows that the desired estimate can be obtained from the following (apparently) weaker condition: there exist $0<\alpha, \beta<1$ such that for every cube $Q \in \mathbb{D}_{Q_{0}}$,

$$
\begin{equation*}
F \subset Q, \quad \frac{\omega_{2}(F)}{\omega_{2}(Q)}<\alpha \quad \Longrightarrow \quad \frac{\omega_{1}(F)}{\omega_{1}(Q)}<\beta \tag{B.11}
\end{equation*}
$$

To prove this result we need a local Calderón-Zygmund decomposition for dyadically doubling weights. The proof is standard and we leave it to the interested reader.

Lemma B.12. - Given $Q_{0}$ and $\omega$ a dyadically doubling measure on $Q_{0}$ with constant $C_{\omega}$, we consider the local dyadic Hardy-Littlewood maximal function with respect to $\omega$ :

$$
\mathcal{M}_{\omega} f(x)=\sup _{x \in Q \in \mathbb{D}_{Q_{0}}} \frac{1}{\omega(Q)} \int_{Q}|f(y)| d \omega(y)
$$

For any $0 \leq f \in L^{1}\left(Q_{0}, \omega\right)$ and $\lambda \geq \frac{1}{\omega\left(Q_{0}\right)} \int_{Q_{0}}|f(y)| d \omega(y)$, there exists a collection of maximal and therefore disjoint dyadic cubes $\left\{Q_{j}\right\} \subset \mathbb{D}_{Q_{0}}$ such that

$$
\begin{equation*}
E_{\lambda}=\left\{x \in Q_{0}: \mathcal{M}_{\omega} f(x)>\lambda\right\}=\bigcup_{j} Q_{j} \tag{B.13}
\end{equation*}
$$

$$
\begin{equation*}
f(x) \leq \lambda, \quad \text { for } \omega \text {-a.e. } x \notin E_{\lambda} \tag{B.14}
\end{equation*}
$$

$$
\begin{equation*}
\lambda<\frac{1}{\omega\left(Q_{j}\right)} \int_{Q_{j}} f(y) d \omega(y) \leq C_{\omega} \lambda \tag{B.15}
\end{equation*}
$$

Proof of Lemma B.7. - We proceed as in [27, Lemma B.4]. Pick $0<\alpha<1$ and $\beta=1-\left(\frac{1-\alpha}{C_{0}}\right)^{1 / \theta_{0}}$, and notice that $0<\beta<1$ since $C_{0} \geq 1$. Then for any $F \subset Q$, $Q \in \mathbb{D}_{Q_{0}}$ we apply (B.8) to $Q \backslash F$ and we conclude (B.11). Next we see that this (apparently) weaker condition implies the desired conclusion. Assume momentarily that $\omega_{1} \ll \omega_{2}$. Then the Radon-Nikodym derivative $h=d \omega_{1} / d \omega_{2}$ satisfies that $h \in L^{1}\left(Q_{0}, \omega_{2}\right)$ and $0 \leq h(x)<\infty$ for $\omega_{2}$-a.e. $x \in Q_{0}$.

Fixed $Q \in \mathbb{D}_{Q_{0}}$ we write $\tau=C_{\omega_{2}} / \alpha$,

$$
\lambda_{0}=\frac{1}{\omega_{2}(Q)} \int_{Q} h(x) d \omega_{2}(x)=\frac{\omega_{1}(Q)}{\omega_{2}(Q)}
$$

and $\lambda_{k}=\tau^{k} \lambda_{0}$. Notice that $\lambda_{0}<\lambda_{1}<\lambda_{2}<\cdots$ since $\tau>C_{\omega_{2}} \geq 1$. For every $k \geq 0$ we apply Lemma B. 12 in $Q$ to $h$ with dyadically doubling measure $\omega_{2}$ : let $\left\{Q_{j}^{k}\right\}_{j} \subset \mathbb{D}_{Q} \subset \mathbb{D}_{Q_{0}}$
$4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - No 3
be the corresponding collection of cubes such that $E_{k}=E_{\lambda_{k}}=\bigcup_{j} Q_{j}^{k}$. Fix $Q_{j_{0}}^{k}$ and observe that if $Q_{j_{0}}^{k} \cap Q_{j}^{k+1} \neq \varnothing$ then $Q_{j}^{k+1} \subset Q_{j_{0}}^{k}$ : otherwise we would have $Q_{j_{0}}^{k} \subsetneq Q_{j}^{k+1}$, by (B.15) we observe that $\frac{1}{\omega_{2}\left(Q_{j}^{k+1}\right)} \int_{Q_{j}^{k+1}} h d \omega_{2}>\lambda_{k+1}>\lambda_{k}$ and then $Q_{j_{0}}^{k}$ would not be maximal. Then using (B.13) and (B.15) we obtain

$$
\begin{aligned}
\omega_{2}\left(Q_{j_{0}}^{k} \cap E_{k+1}\right) & =\sum_{j: Q_{j}^{k+1} \subset Q_{j_{0}}^{k}} \omega_{2}\left(Q_{j}^{k+1}\right)<\frac{1}{\lambda_{k+1}} \sum_{j: Q_{j}^{k+1} \subset Q_{j_{0}}^{k}} \int_{Q_{j}^{k+1}} h d \omega_{2} \\
& \leq \frac{1}{\lambda_{k+1}} \int_{Q_{j_{0}}^{k}} h d \omega_{2} \leq \frac{C_{\omega_{2}} \lambda_{k}}{\lambda_{k+1}} \omega_{2}\left(Q_{j_{0}}^{k}\right)=\alpha \omega_{2}\left(Q_{j_{0}}^{k}\right) .
\end{aligned}
$$

This estimate allows us to use (B.11) which in turn gives that $\omega_{1}\left(Q_{j_{0}}^{k} \cap E_{k+1}\right)<\beta \omega_{1}\left(Q_{j_{0}}^{k}\right)$. Next we sum on $j_{0}$ and conclude that $\omega_{1}\left(E_{k+1}\right)<\beta \omega_{1}\left(E_{k}\right)$ since $E_{k+1} \subset E_{k}$. By iterating this expression we obtain $\omega_{1}\left(E_{k}\right)<\beta^{k} \omega_{1}\left(E_{0}\right)$. Similarly, $\omega_{2}\left(E_{k}\right)<\alpha^{k} \omega_{1}\left(E_{0}\right)$, which implies

$$
\omega_{2}\left(\cap_{k} E_{k}\right)=\lim _{k \rightarrow \infty} \omega_{2}\left(E_{k}\right)=0
$$

Let $0<\varepsilon<-\log \beta / \log \tau$. Then $0<\tau^{\varepsilon} \beta<1$ and by (B.14)

$$
\begin{align*}
\frac{1}{\omega_{2}(Q)} & \int_{Q} h(x)^{1+\varepsilon} d \omega_{2}(x)  \tag{B.16}\\
& =\frac{1}{\omega_{2}(Q)} \int_{Q \backslash E_{0}} h(x)^{1+\varepsilon} d \omega_{2}(x)+\frac{1}{\omega_{2}(Q)} \sum_{k=0}^{\infty} \int_{E_{k} \backslash E_{k+1}} h(x)^{1+\varepsilon} d \omega_{2}(x) \\
& \leq \lambda_{0}^{\varepsilon} \frac{1}{\omega_{2}(Q)} \int_{Q} h(x) d \omega_{2}(x)+\frac{1}{\omega_{2}(Q)} \sum_{k=0}^{\infty} \lambda_{k+1}^{\varepsilon} \int_{E_{k}} h(x) d \omega_{2}(x) \\
& =\lambda_{0}^{\varepsilon} \frac{\omega_{1}(Q)}{\omega_{2}(Q)}+\frac{1}{\omega_{2}(Q)} \sum_{k=0}^{\infty} \lambda_{k+1}^{\varepsilon} \omega_{1}\left(E_{k}\right) \\
& \leq \lambda_{0}^{\varepsilon} \frac{\omega_{1}(Q)}{\omega_{2}(Q)}+\lambda_{0}^{\varepsilon} \frac{\omega_{1}\left(E_{0}\right)}{\omega_{2}(Q)} \sum_{k=0}^{\infty} \tau^{(k+1) \varepsilon} \beta^{k} \\
& \leq \lambda_{0}^{\varepsilon} \frac{\omega_{1}(Q)}{\omega_{2}(Q)}\left(1+\tau^{\varepsilon}\left(1-\tau^{\varepsilon} \beta\right)^{-1}\right)=\left(\frac{\omega_{1}(Q)}{\omega_{2}(Q)}\right)^{1+\varepsilon} C_{1}^{1+\varepsilon} .
\end{align*}
$$

This estimate implies that for all $F \subset Q$,

$$
\begin{aligned}
\frac{\omega_{1}(F)}{\omega_{2}(Q)} & =\frac{1}{\omega_{2}(Q)} \int_{Q} \chi_{F} h d \omega_{2} \leq\left(\frac{1}{\omega_{2}(Q)} \int_{Q} h^{1+\varepsilon} d \omega_{2}\right)^{\frac{1}{1+\varepsilon}}\left(\frac{\omega_{2}(F)}{\omega_{2}(Q)}\right)^{\frac{1}{(1+\varepsilon)^{\prime}}} \\
& \leq \frac{\omega_{1}(Q)}{\omega_{2}(Q)} C_{1}\left(\frac{\omega_{2}(F)}{\omega_{2}(Q)}\right)^{\frac{1}{(1+\varepsilon)^{\prime}}}
\end{aligned}
$$

which is (B.9) with $\theta_{1}=1 /(1+\varepsilon)^{\prime}$. Notice that $\varepsilon$ and $C_{1}$ depend only on $\alpha, \beta$ and $C_{\omega_{2}}$.
Next we see how to proceed in the general case starting from (B.11). We define a new measure $\tilde{\omega}_{2}=\omega_{2}+\delta \omega_{1}$ with $\delta>0$. It is clear that $\omega_{1} \ll \tilde{\omega}_{2}$ and also that $\tilde{\omega}_{2}$ is dyadically doubling on $Q_{0}$ with constant $C_{\tilde{\omega}_{2}}=C_{\omega_{1}}+C_{\omega_{2}}$. We claim that setting
$\tilde{\beta}=1-\min \{1-\beta, \alpha / 2\}, \tilde{\alpha}=\alpha / 2$ we have for every $Q \in \mathbb{D}_{Q_{0}}$,

$$
\begin{equation*}
F \subset Q, \quad \frac{\tilde{\omega}_{2}(F)}{\tilde{\omega}_{2}(Q)}<\tilde{\alpha} \quad \Longrightarrow \quad \frac{\omega_{1}(F)}{\omega_{1}(Q)}<\tilde{\beta} . \tag{B.17}
\end{equation*}
$$

Assuming this, (B.11) holds for $\omega_{1}, \tilde{\omega}_{2}$. By the previous case, since $\omega_{1} \ll \tilde{\omega}_{2}$, there exist $\tilde{\varepsilon}, \tilde{C}_{1}$ such that for every $Q \in \mathbb{D}_{Q_{0}}, F \subset Q$ we have

$$
\frac{\omega_{1}(F)}{\omega_{1}(Q)} \leq \tilde{C}_{1}\left(\frac{\tilde{\omega}_{2}(F)}{\tilde{\omega}_{2}(Q)}\right)^{\frac{1}{(1+\varepsilon)}}
$$

As mentioned above $\tilde{\varepsilon}, \tilde{C}_{1}$ depend only on $\tilde{\alpha}, \tilde{\beta}, C_{\tilde{\omega}_{2}}$ and these are ultimately given in terms of $\alpha, \beta, C_{\omega_{1}}, C_{\omega_{2}}$. Next we see that $\omega_{1} \ll \omega_{2}$ : given $F \subset Q_{0}$ with $\omega_{2}(F)=0$, the previous inequality applied to $Q=Q_{0}$ gives as desired

$$
0 \leq \frac{\omega_{1}(F)}{\omega_{1}(Q)} \leq \tilde{C}_{1}\left(\frac{\delta \omega_{1}(F)}{\tilde{\omega}_{2}\left(Q_{0}\right)}\right)^{\frac{1}{(1+\bar{\varepsilon})^{\prime}}} \leq \tilde{C}_{1}\left(\delta \frac{\omega_{1}(F)}{\omega_{2}\left(Q_{0}\right)}\right)^{\frac{1}{(1+\tilde{\varepsilon})^{\prime}}} \longrightarrow 0, \quad \text { as } \delta \rightarrow 0^{+}
$$

Thus, we get back to the first case and obtain (B.16) which eventually leads to (B.9) with $C_{1}$ and $\theta_{1}$ as stated above.

To complete the proof we obtain (B.17). Given $F$ as there, it follows that $\tilde{\omega}_{2}(Q \backslash F) / \tilde{\omega}_{2}(Q)>1-\alpha / 2$. We see that $\omega_{1}(Q \backslash F) / \omega_{1}(Q)>\min \{1-\beta, \alpha / 2\}$, which yields as desired $\omega_{1}(F) / \omega_{1}(Q)<\tilde{\beta}$. If this were not the case then we would have $\omega_{1}(Q \backslash F) / \omega_{1}(Q) \leq \alpha / 2$ and also that $\omega_{1}(F) / \omega_{1}(Q) \geq \beta$. By (B.11), the latter gives $\omega_{2}(F) / \omega_{2}(Q) \geq \alpha$ and therefore $\omega_{2}(Q \backslash F) / \omega_{2}(Q) \leq 1-\alpha$. Gathering these estimates we get a contradiction

$$
\frac{\tilde{\omega}_{2}(Q \backslash F)}{\tilde{\omega}_{2}(Q)}=\frac{\omega_{2}(Q \backslash F)}{\tilde{\omega}_{2}(Q)}+\delta \frac{\omega_{1}(Q \backslash F)}{\tilde{\omega}_{2}(Q)} \leq \frac{\omega_{2}(Q \backslash F)}{\omega_{2}(Q)}+\frac{\omega_{1}(Q \backslash F)}{\omega_{1}(Q)} \leq 1-\alpha / 2 .
$$

Remark B.18. - Let us observe that (B.16) can be equivalently written as

$$
\left(\frac{1}{\omega_{2}(Q)} \int_{Q} h(x)^{1+\varepsilon} d \omega_{2}(x)\right)^{\frac{1}{1+\varepsilon}} \leq C_{1} \frac{1}{\omega_{2}(Q)} \int_{Q} h(x) d \omega_{2}(x),
$$

and this shows that $h \in R H_{1+\varepsilon}^{\text {dyadic }}\left(Q_{0}, \omega_{2}\right)$.

## Appendix C

## The UR property for approximating domains

We establish the UR property (with uniform constants) for the approximating domains $\Omega_{N}$ defined by (8.13). Recall that we have already observed that $\Omega_{N}$ inherits the ADR, Corkscrew and Harnack chain conditions from $\Omega$.

The proof is based on ideas of Guy David, and uses the following singular integral characterization of UR sets, established in [18]. Suppose that $E \subset \mathbb{R}^{n+1}$ is $n$-dimensional ADR. The singular integral operators that we shall consider are those of the form

$$
T_{E, \varepsilon} f(x)=T_{\varepsilon} f(x):=\int_{E} K_{\varepsilon}(x-y) f(y) d H^{n}(y)
$$

where $K_{\varepsilon}(x):=K(x) \Phi(|x| / \varepsilon)$, with $0 \leq \Phi \leq 1, \Phi(\rho) \equiv 1$ if $\rho \geq 2, \Phi(\rho) \equiv 0$ if $\rho \leq 1$, and $\Phi \in C^{\infty}(\mathbb{R})$, and where the singular kernel $K$ is an odd function, smooth on $\mathbb{R}^{n+1} \backslash\{0\}$, and satisfying

$$
\begin{align*}
|K(x)| & \leq C|x|^{-n}  \tag{C.1}\\
\left|\nabla^{m} K(x)\right| & \leq C_{m}|x|^{-n-m}, \quad \forall m \in\{1,2,3, \ldots\} \tag{C.2}
\end{align*}
$$

Then $E$ is UR if and only if for every such kernel $K$, we have that

$$
\begin{equation*}
\sup _{\varepsilon>0} \int_{E}\left|T_{\varepsilon} f\right|^{2} d H^{n} \leq C_{K} \int_{E}|f|^{2} d H^{n} \tag{C.3}
\end{equation*}
$$

We refer the reader to [18] for the proof. We shall also require "non-tangential" estimates for an extension of $T_{\varepsilon}$ defined as follows. For $K$ as above, set

$$
\begin{equation*}
\mathcal{J}_{E} f(X):=\int_{E} K(X-y) f(y) d H^{n}(y), \quad X \in \mathbb{R}^{n+1} \backslash E \tag{C.4}
\end{equation*}
$$

We define non-tangential approach regions $\Gamma_{\tau}(x)$ as follows. Let $\mathcal{W}_{E}$ denote the collection of cubes in the Whitney decomposition of $\mathbb{R}^{n+1} \backslash E$, and set

$$
\mathcal{W}_{\tau}(x):=\left\{I \in \mathcal{W}_{E}: \operatorname{dist}(I, x)<\tau \ell(I)\right\}
$$

Then we define

$$
\Gamma_{\tau}(x):=\bigcup_{I \in W_{\tau}(x)} I^{*}
$$

(thus, roughly speaking, $\tau$ is the "aperture" of $\Gamma_{\tau}(x)$ ). For $F \in \mathscr{C}\left(\mathbb{R}^{n+1} \backslash E\right)$ we may then also define the non-tangential maximal function

$$
N_{*, \tau}(F)(x):=\sup _{Y \in \Gamma_{\tau}(x)}|F(Y)| .
$$

We shall sometimes write simply $N_{*}$ when there is no chance of confusion in leaving implicit the dependence on the aperture $\tau$.

Lemma C.5. - Suppose that $E \subset \mathbb{R}^{n+1}$ is n-dimensional $U R$, and let $\mathcal{J}_{E}$ be defined as in (C.4). Then for each $\tau \in(0, \infty)$, there is a constant $C_{\tau, K}$ depending only on $n, \tau, K$ and the UR constants such that

$$
\begin{equation*}
\int_{E}\left(N_{*, \tau}\left(\mathcal{G}_{E} f\right)\right)^{2} d H^{n} \leq C_{\tau, K} \int_{E}|f|^{2} d H^{n} \tag{C.6}
\end{equation*}
$$

Given (C.3), Lemma C. 5 is a variant of the standard "Cotlar inequality" for maximal singular integrals, and we omit the proof.

We are now ready to prove that $\partial \Omega_{N}$ is UR, uniformly in $N$. It is enough to establish the estimate (C.3), for all $K$ as above, with $E$ replaced by $\partial \Omega_{N}$. On the other hand, we are given that $\partial \Omega$ is UR, whence (C.6) holds with $E=\partial \Omega$. Since $\partial \Omega_{N}$ is ADR, it enjoys the dyadic grid structure promised by Lemma ??. We then make a partition $\partial \Omega_{N}=\cup Q_{j}(N)$, where $Q_{j}(N) \in \mathbb{D}_{N}\left(\partial \Omega_{N}\right)=: \mathbb{D}_{N}(N)$, the dyadic grid on $\partial \Omega_{N}$ at scale $2^{-N}$. We observe that, by the construction of $\Omega_{N}$, for each $Q_{j}(N) \in \mathbb{D}_{N}(N)$, we may choose a $Q_{j} \in \mathbb{D}_{N}(\partial \Omega)=$ : $\mathbb{D}_{N}$ with $\operatorname{dist}\left(Q_{j}(N), Q_{j}\right) \approx 2^{-N}$. By the ADR property of $\partial \Omega$, a given $Q \in \mathbb{D}_{N}$ can serve in this
way for at most a bounded number of $Q_{j}(N) \in \mathbb{D}_{N}(N)$. Therefore, we have the bounded overlap condition

$$
\begin{equation*}
\sum_{Q_{j}(N) \in \mathbb{D}_{N}(N)} 1_{Q_{j}}(x) \leq C, \quad \forall x \in \partial \Omega \tag{C.7}
\end{equation*}
$$

As usual, we set $\sigma:=\left.H^{n}\right|_{\partial \Omega}$, and we now also let $\sigma_{N}:=\left.H^{n}\right|_{\partial \Omega_{N}}$. We then have that for $\tau$ large enough,

$$
\begin{aligned}
\int_{\partial \Omega_{N}}\left|\mathcal{T}_{\partial \Omega} f\right|^{2} d \sigma_{N} & =\sum_{Q_{j}(N) \in \mathbb{D}_{N}(N)} \int_{Q_{j}(N)}\left|\mathcal{T}_{\partial \Omega} f\right|^{2} d \sigma_{N} \\
& =\sum_{Q_{j}(N) \in \mathbb{D}_{N}(N)} \frac{1}{\sigma\left(Q_{j}\right)} \int_{Q_{j}} \int_{Q_{j}(N)}\left|\mathcal{T}_{\partial \Omega} f(x)\right|^{2} d \sigma_{N}(x) d \sigma\left(x^{\prime}\right) \\
& \lesssim \sum_{Q_{j}(N) \in \mathbb{D}_{N}(N)} \int_{Q_{j}}\left(N_{*, \tau}\left(\mathcal{T}_{\partial \Omega} f\right)\right)^{2} d \sigma \leq C_{\tau, K} \int_{\partial \Omega}|f|^{2} d \sigma
\end{aligned}
$$

where in the last line we have used first the ADR properties of $\partial \Omega$ and $\partial \Omega_{N}$, and then (C.7) and (C.6) with $E=\partial \Omega$.

We have thus established that $\mathcal{J}_{\partial \Omega}: L^{2}(\partial \Omega) \rightarrow L^{2}\left(\partial \Omega_{N}\right)$. Since the kernel $K$ is odd, we therefore obtain by duality that

$$
\begin{equation*}
\mathcal{J}_{\partial \Omega_{N}}: L^{2}\left(\partial \Omega_{N}\right) \rightarrow L^{2}(\partial \Omega) \tag{C.8}
\end{equation*}
$$

Now fix $\varepsilon>0$, and $N$ large enough that $2^{-N} \ll \operatorname{diam} \partial \Omega$. Set $\varepsilon_{N}=2^{-N}$. We consider two cases.

Case 1: $\varepsilon<\varepsilon_{N}$. In this case,

$$
\begin{aligned}
& \int_{\partial \Omega_{N}}\left|T_{\partial \Omega_{N}, \varepsilon} f\right|^{2} d \sigma_{N} \\
& \quad \lesssim \int_{\partial \Omega_{N}}\left|T_{\partial \Omega_{N}, \varepsilon} f-T_{\partial \Omega_{N}, \varepsilon_{N}} f\right|^{2} d \sigma_{N}+\int_{\partial \Omega_{N}}\left|T_{\partial \Omega_{N}, \varepsilon_{N}} f\right|^{2} d \sigma_{N}=: I+I I
\end{aligned}
$$

Let $Q_{j} \in \mathbb{D}_{N}$ denote the cube chosen relative to $Q_{j}(N) \in \mathbb{D}_{N}(N)$ as above. Then for $x \in Q_{j}(N)$, and $x^{\prime} \in Q_{j}$, we have by standard Calderón-Zygmund estimates using (C.1) and (C.2), and the ADR property of $\partial \Omega_{N}$, that

$$
\left|T_{\partial \Omega_{N}, \varepsilon_{N}} f(x)-\mathcal{T}_{\partial \Omega_{N}} f\left(x^{\prime}\right)\right| \lesssim M^{N} f(x)
$$

where $M^{N}$ denotes the Hardy-Littlewood maximal function on $\partial \Omega_{N}$. Consequently,

$$
\begin{aligned}
I I & =\sum_{Q_{j}(N) \in \mathbb{D}_{N}(N)} \frac{1}{\sigma\left(Q_{j}\right)} \int_{Q_{j}} \int_{Q_{j}(N)}\left|T_{\partial \Omega_{N}, \varepsilon_{N}} f(x)\right|^{2} d \sigma_{N}(x) d \sigma\left(x^{\prime}\right) \\
& \lesssim \sum_{Q_{j}(N) \in \mathbb{D}_{N}(N)} \int_{Q_{j}(N)}\left(M^{N} f(x)\right)^{2} d \sigma_{N}(x) \\
& +\sum_{Q_{j}(N) \in \mathbb{D}_{N}(N)} \int_{Q_{j}}\left|\mathcal{T}_{\partial \Omega_{N}} f\left(x^{\prime}\right)\right|^{2} d \sigma\left(x^{\prime}\right)=: I I^{\prime}+I I^{\prime \prime}
\end{aligned}
$$

The desired bound for $I I^{\prime}$ follows immediately, and the bound for $I I^{\prime \prime}$ follows directly from (C.7) and (C.8).

We turn now to term $I$. Let us note that since $\varepsilon_{N} \approx \operatorname{diam}\left(Q_{j}(N)\right)$, for $x \in Q_{j}(N)$ we have

$$
\begin{aligned}
& T_{\partial \Omega_{N}, \varepsilon} f(x)-T_{\partial \Omega_{N}, \varepsilon_{N}} f(x) \\
& \quad=\int_{\partial \Omega_{N}} K(x-y)\left(\Phi\left(\frac{|x-y|}{\varepsilon}\right)-\Phi\left(\frac{|x-y|}{\varepsilon_{N}}\right)\right) f(y) 1_{\Delta_{N, j}}(y) d \sigma_{N}(y) \\
& \quad=: T_{\partial \Omega_{N}, \varepsilon, \varepsilon_{N}}\left(f 1_{\Delta_{N, j}}\right)(x)
\end{aligned}
$$

a doubly truncated singular integral on $\partial \Omega_{N}$, where $\Delta_{N, j}:=B_{N, j} \cap \partial \Omega_{N}$, and $B_{N, j}$ is a ball centered at some point in $Q_{j}(N)$, with radius $C \operatorname{diam}\left(Q_{j}(N)\right) \approx 2^{-N}$. By choosing $C$ large enough, we may assume that $Q_{j}(N) \subset \Delta_{N, j}$. We recall that by definition, $\partial \Omega_{N}$ is a union of portions of faces of fattened Whitney cubes $I^{*}$, of side length $\approx 2^{-N}$. Since only a bounded number of these can meet $B_{N, j}$, we have

$$
\Delta_{N, j} \subset \bigcup_{m=1}^{M_{0}} F_{m}^{j}
$$

where $M_{0}$ is a uniform constant and each $F_{m}^{j}$ is either a portion of a face of some $I^{*}$, or else $F_{m}^{j}=\varnothing$ (since $M_{0}$ is not necessarily equal to the number of faces, but is rather an upper bound for the number of faces). Thus,

$$
I \lesssim \sum_{Q_{j}(N) \in \mathbb{D}_{N}(N)} \sum_{1 \leq m, m^{\prime} \leq M_{0}} \int_{F_{m}^{j}}\left|T_{\partial \Omega_{N}, \varepsilon, \varepsilon_{N}}\left(f 1_{F_{m^{\prime}}^{j}}\right)\right|^{2} d \sigma_{N}
$$

The faces $F_{m^{\prime}}^{j}$ have bounded overlaps as we sum in $j$. Therefore, the case $m=m^{\prime}$ reduces to the classical case that $\partial \Omega_{N}$ is a hyperplane. For $m \neq m^{\prime}$, there are two cases as follows. If $\operatorname{dist}\left(F_{m}^{j}, F_{m^{\prime}}^{j}\right) \approx 2^{-N}$, then using (C.1), we may crudely dominate $T_{\partial \Omega_{N}, \varepsilon, \varepsilon_{N}}$ by the HardyLittlewood maximal operator. Otherwise, $\operatorname{dist}\left(F_{m}^{j}, F_{m^{\prime}}^{j}\right) \ll 2^{-N}$, in which case $F_{m}^{j}$ and $F_{m^{\prime}}^{j}$ are contained in respective faces which either lie in the same hyperplane, or else meet at an angle of $\pi / 2$. In the latter scenario, after a possible rotation of co-ordinates, we may view $F_{m}^{j} \cup F_{m^{\prime}}^{j}$ as lying in a Lipschitz graph with Lipschitz constant 1 , so that we may estimate $T_{\partial \Omega_{N}, \varepsilon, \varepsilon_{N}}$ using an extension of the Coifman-McIntosh-Meyer theorem.

CASE 2: $\varepsilon \geq \varepsilon_{N}$. We observe that (C.8) also applies to the modified operator $\mathcal{T}_{\partial \Omega_{N}}^{\varepsilon}$, obtained by replacing the kernel $K$ by the kernel $K_{\varepsilon}$, since the latter is still odd and still satisfies the Calderón-Zygmund estimates (C.1) and (C.2) (uniformly in $\varepsilon$ ). The present case may then be handled just like term $I I$ above, by writing

$$
T_{\partial \Omega_{N}, \varepsilon} f(x)=\left(T_{\partial \Omega_{N}, \varepsilon} f(x)-\mathscr{J}_{\partial \Omega_{N}}^{\varepsilon} f\left(x^{\prime}\right)\right)+\mathscr{G}_{\partial \Omega_{N}}^{\varepsilon} f\left(x^{\prime}\right)
$$

There is no term $I$. We leave the details to the reader.

## BIBLIOGRAPHY

[1] H. Aikawa, Boundary Harnack principle and Martin boundary for a uniform domain, J. Math. Soc. Japan 53 (2001), 119-145.
[2] H. Aikawa, Equivalence between the boundary Harnack principle and the Carleson estimate, Math. Scand. 103 (2008), 61-76.
[3] P. Auscher, S. Hofmann, J. L. Lewis, P. Tchamitchian, Extrapolation of Carleson measures and the analyticity of Kato's square-root operators, Acta Math. 187 (2001), 161-190.
[4] P. Auscher, S. Hofmann, C. Muscalu, T. Tao, C. Thiele, Carleson measures, trees, extrapolation, and $T(b)$ theorems, Publ. Mat. 46 (2002), 257-325.
[5] J. Azzam, R. Schul, Hard Sard: quantitative implicit function and extension theorems for Lipschitz maps, Geom. Funct. Anal. 22 (2012), 1062-1123.
[6] M. Badger, Null sets of harmonic measure on NTA domains: Lipschitz approximation revisited, Math. Z. 270 (2012), 241-262.
[7] B. Bennewitz, J. L. Lewis, On weak reverse Hölder inequalities for nondoubling harmonic measures, Complex Var. Theory Appl. 49 (2004), 571-582.
[8] C. J. Bishop, P. W. Jones, Harmonic measure and arclength, Ann. of Math. 132 (1990), 511-547.
[9] J. Bourgain, On the Hausdorff dimension of harmonic measure in higher dimension, Invent. Math. 87 (1987), 477-483.
[10] L. Caffarelli, E. Fabes, S. Mortola, S. Salsa, Boundary behavior of nonnegative solutions of elliptic operators in divergence form, Indiana Univ. Math. J. 30 (1981), 621-640.
[11] L. Carleson, Interpolations by bounded analytic functions and the Corona problem, Ann. of Math. 76 (1962), 547-559.
[12] L. Carleson, J. Garnett, Interpolating sequences and separation properties, Journal d'Analyse Mathématique 28 (1975), 273-299.
[13] M. Christ, A $T(b)$ theorem with remarks on analytic capacity and the Cauchy integral, Colloq. Math. 60/61 (1990), 601-628.
[14] R. R. Coifman, C. Fefferman, Weighted norm inequalities for maximal functions and singular integrals, Studia Math. 51 (1974), 241-250.
[15] B. E. J. Dahlberg, Estimates of harmonic measure, Arch. Rational Mech. Anal. 65 (1977), 275-288.
[16] B. E. J. Dahlberg, D. S. Jerison, C. E. Kenig, Area integral estimates for elliptic differential operators with nonsmooth coefficients, Ark. Mat. 22 (1984), 97-108.
[17] G. David, D. S. Jerison, Lipschitz approximation to hypersurfaces, harmonic measure, and singular integrals, Indiana Univ. Math. J. 39 (1990), 831-845.
[18] G. David, S. Semmes, Singular integrals and rectifiable sets in $\mathbf{R}^{n}$ : Beyond Lipschitz graphs, Astérisque 193 (1991).
[19] G. David, S. Semmes, Analysis of and on uniformly rectifiable sets, Mathematical Surveys and Monographs 38, Amer. Math. Soc., Providence, RI, 1993.
[20] L. C. Evans, R. F. Gariepy, Measure theory and fine properties of functions, Studies in Advanced Math., CRC Press, Boca Raton, FL, 1992.
[21] J. García-Cuerva, J. L. Rubio de Francia, Weighted norm inequalities and related topics, North-Holland Mathematics Studies 116, North-Holland Publishing Co., Amsterdam, 1985.
$4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - No 3
[22] D. Gilbarg, N. S. Trudinger, Elliptic partial differential equations of second order, second ed., Grundl. Math. Wiss. 224, Springer, Berlin, 1983.
[23] M. Grüter, K.-O. Widman, The Green function for uniformly elliptic equations, Manuscripta Math. 37 (1982), 303-342.
[24] A. Grau de la Herran, M. Mourgoglou, A Tb theorem for square functions in domains with Ahlfors-David regular boundaries, J. Geom. Anal. (2013), 1-22.
[25] S. Hofmann, J. L. Lewis, The Dirichlet problem for parabolic operators with singular drift terms, Mem. Amer. Math. Soc. 151 (2001).
[26] S. Hofmann, J. M. Martell, A note on $A_{\infty}$ estimates via extrapolation of Carleson measures, in The AMSI-ANU Workshop on Spectral Theory and Harmonic Analysis, Proc. Centre Math. Appl. Austral. Nat. Univ. 44, Austral. Nat. Univ., Canberra, 2010, 143-166.
[27] S. Hofmann, J. M. Martell, $A_{\infty}$ estimates via extrapolation of Carleson measures and applications to divergence form elliptic operators, Trans. Amer. Math. Soc. 364 (2012), 65-101.
[28] S. Hofmann, J. M. Martell, I. Uriarte-Tuero, Uniform rectifiability and harmonic measure II: Poisson kernels in $L^{p}$ imply uniform rectifiability, to appear in Duke Math. J.
[29] D. S. Jerison, C. E. Kenig, Boundary behavior of harmonic functions in nontangentially accessible domains, Adv. in Math. 46 (1982), 80-147.
[30] P. W. Jones, Square functions, Cauchy integrals, analytic capacity, and harmonic measure, in Harmonic analysis and partial differential equations (El Escorial, 1987), Lecture Notes in Math. 1384, Springer, Berlin, 1989, 24-68.
[31] C. E. Kenig, Harmonic analysis techniques for second order elliptic boundary value problems, CBMS Regional Conference Series in Mathematics 83, Amer. Math. Soc., 1994.
[32] C. E. Kenig, T. Toro, Harmonic measure on locally flat domains, Duke Math. J. 87 (1997), 509-551.
[33] C. E. Kenig, T. Toro, Free boundary regularity for harmonic measures and Poisson kernels, Ann. of Math. 150 (1999), 369-454.
[34] C. E. Kenig, T. Toro, Poisson kernel characterization of Reifenberg flat chord arc domains, Ann. Sci. École Norm. Sup. 36 (2003), 323-401.
[35] M. Lavrent'ev, Boundary problems in the theory of univalent functions, Amer. Math. Soc. Transl. 32 (1963), 1-35.
[36] J. L. Lewis, M. A. M. Murray, The method of layer potentials for the heat equation in time-varying domains, Mem. Amer. Math. Soc. 114 (1995).
[37] F. Riesz, M. Riesz, Über die Randwerte einer analytischen Funktion, in Comptes rendus du quatrième congrès des mathématiciens scandinaves, Stockholm 1916, Almqvist \& Wilksels, 1920.
[38] E. T. SAWYER, Two weight norm inequalities for certain maximal and integral operators, in Harmonic analysis ( Minneapolis, Minn., 1981), Lecture Notes in Math. 908, Springer, Berlin-New York, 1982, 102-127.
[39] S. Semmes, Analysis vs. geometry on a class of rectifiable hypersurfaces in $\mathbf{R}^{n}$, Indiana Univ. Math. J. 39 (1990), 1005-1035.
[40] E. M. Stein, Singular integrals and differentiability properties of functions, Princeton Mathematical Series, No. 30, Princeton Univ. Press, Princeton, N.J., 1970.
(Manuscrit reçu le 5 octobre 2012; accepté, après révision, le 12 avril 2013.)

Steve Hofmann<br>Department of Mathematics<br>University of Missouri<br>Columbia, MO 65211, USA<br>E-mail: hofmanns@missouri.edu<br>José María Martell<br>Instituto de Ciencias Matemáticas<br>CSIC-UAM-UC3M-UCM<br>Consejo Superior de Investigaciones Científicas<br>C/ Nicolás Cabrera, 13-15<br>E-28049 Madrid, Spain<br>E-mail: chema.martell@icmat.es


[^0]:    The first author was supported by NSF grant DMS-0801079. The second author was supported by MINECO Grant MTM2010-16518 and ICMAT Severo Ochoa project SEV-2011-0087.
    This work has been possible thanks to the support and hospitality of the University of Missouri-Columbia (USA), the Consejo Superior de Investigaciones Cientificas (Spain), the Universidad Autónoma de Madrid (Spain), and the Australian National University (Canberra, Australia). Both authors express their gratitude to these institutions. Both authors are also grateful to P. Auscher and M. Perrin for their help in the French translations.

[^1]:    ${ }^{(1)}$ On the other hand, Azzam and Schul [5] have recently shown that every UR set contains "big pieces of big pieces of Lipschitz graphs" (see [19, pp. 15-16] or [5] for a precise formulation). This is a beautiful result, but seems inapplicable to the estimates for harmonic measure considered here: to enable essential use of the maximum principle, one would need "interior big pieces (cf. Definition 1.14 below) of interior big pieces of Lipschitz subdomains" (say, in the presence of the 1-sided NTA condition), and it is not clear that the methods of [5] would yield such a result. We do expect that the methods of the present paper could be pushed to do so, and we plan to present these arguments, with applications to more general elliptic-harmonic measures, in a forthcoming paper.

[^2]:    ${ }^{(2)}$ We recall that such domains are sometimes denoted "uniform" domains in the literature, but we prefer the terminology " 1 -sided NTA", both because it is more descriptive of the actual properties enjoyed by such domains, and to avoid confusion with the completely different notion of "uniform rectifiability".

[^3]:    ${ }^{(3)}$ Since we have made no assumption as regards Wiener's regularity criterion, our harmonic function is a generalized solution, which may not be continuous up to the boundary.
    $4^{\mathrm{e}}$ SÉRIE - TOME 47 - 2014 - No 3

