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Weiyuan QIU \& Pascale ROESCH \& Xiaoguang WANG \& Yongcheng YIN

Hyperbolic components of McMullen maps

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#### Abstract

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# HYPERBOLIC COMPONENTS OF McMULLEN MAPS 

by Weiyuan QIU, Pascale ROESCH, Xiaoguang WANG and Yongcheng YIN


#### Abstract

In this article, we completely settle a question raised by B. Devaney. We prove that all the hyperbolic components are Jordan domains in the family of rational maps of McMullen type. Moreover, we give a precise description of all the rational maps on the outer boundary. It follows that the cusps are dense on the outer boundary.

Résumé. - Dans cet article nous résolvons complètement une question posée par B. Devaney. Nous montrons que toutes les composantes hyperboliques sont des domaines de Jordan dans la famille de fractions rationnelles de type McMullen. De plus nous donnons une description précise de toutes les fractions du bord de la composante non bornée. Il en découle que les cusps sont denses dans le bord de la composante non bornée.


## 1. Introduction

In his article [22], Curt McMullen presented a family of rational maps with the particularity that, viewed as a dynamical system, it exhibits very rich dynamical behavior. Nevertheless, this family has a very simple form. It consists in a singular perturbation of the monomial $z \mapsto z^{n}$ acting on the Riemann sphere $\widehat{\mathbb{C}}=\mathbb{C} \cup\{\infty\}$ :

$$
\begin{aligned}
& \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}} \\
& z \mapsto z^{n}+\lambda z^{-m}
\end{aligned}
$$

where $\lambda$ varies in $\mathbb{C}^{*}=\mathbb{C} \backslash\{0\}$ and $n, m \geq 1$.
The Julia set, which is the minimal totally invariant compact set of cardinality at least 3 , appears in this family under several various classic fractals. Namely, Curt McMullen pointed out (in [22]) that when $(n, m)=(2,3)$ and $\lambda \in \mathbb{C}^{*}$ is small, the Julia set is a Cantor set of circles. Moreover, in the hyperbolic case, the Julia set can also be homeomorphic to either

[^0]a Cantor set, or a Sierpiński carpet as proven in [10]. A rational map is hyperbolic if every critical point converges under iterations to an attracting cycle.

The parameters $\lambda \in \mathbb{C}^{*}$ can then be divided into two classes, the hyperbolic ones and the others (see Figure 1, where hyperbolic parameters are in blue and yellow). Conjecturally the set of hyperbolic parameters (which is an open set) is dense in $\mathbb{C}^{*}$. In this article we study the boundaries of the hyperbolic components of the McMullen maps:

$$
f_{\lambda}: z \mapsto z^{n}+\lambda z^{-n}, \quad \lambda \in \mathbb{C}^{*}, n \geq 3
$$



Figure 1. Parameter plane of McMullen maps, $n=3,4$.

Bob Devaney has proven in [6] that the boundary of the hyperbolic component containing the punctured neighborhood of the origin is a Jordan curve. He raised the question in 2004 at the Snowbird Conference (on the 25th Anniversary of the Mandelbrot set) whether all the other hyperbolic components of escape type (the free critical orbits escape to $\infty$ ) are Jordan domains.

The following Theorem 1.1 solves completely the question, it is the main result of the paper.

Theorem 1.1. - Fix any $n \geq 3$. The boundary of every hyperbolic component of the family $f_{\lambda}(z)=z^{n}+\lambda / z^{n}, \lambda \in \mathbb{C}^{*}$ is a Jordan curve.

Moreover, we give a complete description of the dynamics of the McMullen maps lying at the boundary of the unbounded hyperbolic component $\mathscr{H}_{0}$. In this component, the Julia sets of the maps are Cantor sets, but at the boundary the Julia set is connected (see Section 2). As a corollary we obtain our second main result :

Theorem 1.2. - Cusps are dense in $\partial \mathscr{H}_{0}$.

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4e SÉRIE - TOME 48-2015 - No 3
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Here, according to McMullen [25], a parameter $\lambda \in \mathbb{C}^{*}$ is called a cusp if the map $f_{\lambda}$ has a parabolic cycle. Geometrically, a cusp is a point where the bifurcation locus is cusp-shaped.

Let us recall that McMullen proved that cusps are dense on the Bers' boundary of Teichmüller space in [23]. The analogue of this result in the world of rational maps, as a conjecture posed by McMullen [25], would be, in the space of degree $d$ polynomials, the density of geometrically finite parabolics on the boundary of the hyperbolic component $\psi_{d}$ containing $z^{d}$. This conjecture is verified by P. Roesch [32] in the one-dimensional slice $z^{d}+c z^{d-1}(d \geq 3$ and $c \in \mathbb{C}$ ). It remains open in full generality. The result we prove here is, in the sprit, related to this conjecture. But it has to be thought of as a phenomenon in some one-dimensional slices of rational maps.

We would like to explain why we concentrate on the case $m=n \geq 3$ of the general McMullen family $z \mapsto z^{n}+\lambda z^{-m}, \lambda \in \mathbb{C}^{*}, m, n \geq 1$. This is because our proof rests on the technical Yoccoz puzzle theory. When $m=n \geq 3$, we can succeed in applying this theory to study both the dynamical plane and the parameter plane. However when $m=n=2$, it is impossible to find a non-degenerate critical annulus for the Yoccoz puzzle constructed in [28]. The existence of a non-degenerate critical annulus is technically necessary in this theory. In the general case when $m \neq n \geq 1,0$ and $\infty$ have different folds of symmetries. To the best of our knowledge, whether the Yoccoz puzzle structure exists in this case is unknown.

### 1.1. Overview of the paper

Let us recall some definitions and give the basic notions to precise our results before to go to the proof.

For any $\lambda \in \mathbb{C}^{*}$, the map $f_{\lambda}$ has a fixed point at $\infty$ which is superattracting since the derivative is 0 . The immediate attracting basin of $\infty$ denoted by $B_{\lambda}$ is the set of points converging under iteration to $\infty$ and lying in the connected component of $\infty$. The set of critical points of $f_{\lambda}$ is $\{0, \infty\} \cup C_{\lambda}$, where $C_{\lambda}=\left\{c \in \mathbb{C} ; c^{2 n}=\lambda\right\}$. Hence, there is only one free critical orbit (up to a sign) since besides $\infty$, there are only two critical values: $v_{\lambda}^{+}=2 \sqrt{\lambda}$ and $v_{\lambda}^{-}=-2 \sqrt{\lambda}$ (here, when restricted to the fundamental domain, $v_{\lambda}^{+}$and $v_{\lambda}^{-}$are welldefined, see Section 3).

A rational map is hyperbolic if all critical orbits are attracted by the attracting cycles (see [27, 24]). Hence, a McMullen map $f_{\lambda}$ is hyperbolic if the free critical orbit is attracted either by $\infty$ or by an attracting cycle in $\mathbb{C}$. Every hyperbolic component is isomorphic to either the unit disk $\mathbb{D}$ or $\mathbb{D}^{*}=\mathbb{D}-\{0\}$ (see Theorem 2.2). In particular, $\mathscr{H}_{0}$ is a topological punctured disk. Assuming Theorem 1.1, one gets a canonical parameterization $\nu: \mathbb{S} \rightarrow \partial \mathscr{H}_{0}$, where $\nu(\theta)$ is defined to be the landing point of the parameter ray $\mathscr{R}_{0}(\theta)$ (see Section 6) in $\mathscr{H}_{0}$. The complete characterization of $\partial \mathscr{H}_{0}$ we give is the following :

Theorem 1.3 (Characterization of $\partial \mathscr{H}_{0}$ and cusps). - We have

1. $\lambda \in \partial \mathscr{H}_{0}$ if and only if $\partial B_{\lambda}$ contains either $C_{\lambda}$ or a parabolic cycle.
2. $\nu(\theta)$ is a cusp if and only if $n^{p} \theta \equiv \theta \bmod \mathbb{Z}$ for some $p \geq 1$.

Theorem 1.2 is an immediate consequence of Theorem 1.3 since the set $\left\{\theta \mid n^{p} \theta \equiv \theta \bmod \mathbb{Z}, p \geq 1\right\}$ is a dense subset of the unit circle $\mathbb{S}$.

The main part of the paper is to prove Theorem 1.1. We briefly sketch the idea of the proof and the organization of the paper.

The proof rests on the dynamics and namely the fact that this family admits the "Yoccoz puzzle" structure (see [28]). The Yoccoz puzzle is induced by a kind of Jordan curve called "cut ray" that was first constructed by Bob Devaney [8].

To obtain results in the parameter plane, the idea is to use the rigidity of the maps. And for this we study deeply the dynamical planes. The idea is therefore different from the well known "parapuzzle techniques" (known to be a powerful tool to study the boundary of hyperbolic components, see [32, 30]). Until Section 7, we consider only the hyperbolic components of escape type, which are called escape domains, i.e., for which the critical orbits tend to $\infty$.

Precisely,

- In Section 2, we parameterize the escape domains.
- In Section 3, we recall (quickly) the construction of the cut rays, since it is necessary for the study of escape domains. The crucial fact regarding cut rays is that they move continuously in the Hausdorff topology with respect to the parameter.
- In Section 4, we give some characterizations of maps on $\partial \mathscr{H}_{0}$.
- In Section 5, we prove a rigidity result for the maps on $\partial \mathscr{H}_{0}$. We first construct a topological conjugacy between the maps with the same combinatorial information, using the Yoccoz puzzle techniques. Then we use the idea due to Kozlovski, Shen and van Strien [16], and a 'zero measure argument' following Lyubich, to get the rigidity result.
- In Section 6, we prove that $\partial \mathscr{H}_{0}$ is a Jordan curve and give some related consequences.
- In Section 7, we prove that the boundaries of all escape domains of level $k \geq 3$ are Jordan curves. These escape domains are called Sierpiński holes. The proof is based on three ingredients: the boundary regularity of $\mathcal{H}_{0}$, holomorphic motion and continuity of cut rays. We remark that our approach also applies to the boundary of the hyperbolic component containing the punctured neighborhood of the origin. This yields a different proof from Devaney's in [6].
- In Section 8, we show that every hyperbolic component is a Jordan domain by considering the ones which are not of "escape type".

We refer the reader to $[6,7,8,9,1,10,11,15,28,31,34]$ and the references therein for various related results on the dynamics of McMullen maps.

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## 2. Escape domains and parameterizations

There are two kinds of hyperbolic McMullen maps based on the behavior of the free critical orbit. If the orbit escapes to infinity, the corresponding hyperbolic component is called an escape domain. If the orbit tends to an attracting cycle other than $\infty$, the corresponding hyperbolic component is of renormalizable type.
$4^{\mathrm{e}}$ SÉRIE - TOME 48 - 2015 - No 3

In this section, we present some known facts about escape domains. We refer the reader to [10] for more background materials. The hyperbolic components of renormalizable type will be discussed in Section 8.

For any $\lambda \in \mathbb{C}^{*}$, the Julia set $J\left(f_{\lambda}\right)$ of $f_{\lambda}$ can be identified as the boundary of $\bigcup_{k \geq 0} f_{\lambda}^{-k}\left(B_{\lambda}\right)$. It satisfies $e^{\pi i / n} J\left(f_{\lambda}\right)=J\left(f_{\lambda}\right)$. The Fatou set $F\left(f_{\lambda}\right)$ of $f_{\lambda}$ is defined by $F\left(f_{\lambda}\right)=\widehat{\mathbb{C}}-J\left(f_{\lambda}\right)$. We denote by $T_{\lambda}$ the component of $f_{\lambda}^{-1}\left(B_{\lambda}\right)$ containing 0 . It is possible that $B_{\lambda}=T_{\lambda}$. In that case, the critical set $C_{\lambda} \subset B_{\lambda}$ and $J\left(f_{\lambda}\right)$ is a Cantor set (see Theorem 2.1).

For any $k \geq 0$, we define a parameter set $\mathscr{H}_{k}$ as follows:

$$
\mathscr{H}_{k}=\left\{\lambda \in \mathbb{C}^{*} ; k \text { is the first integer such that } f_{\lambda}^{k}\left(C_{\lambda}\right) \subset B_{\lambda}\right\} .
$$

A component of $\mathscr{H}_{k}$ is called an escape domain of level $k$. One may verify that $\mathscr{H}_{0}=\left\{\lambda \in \mathbb{C}^{*} ; v_{\lambda}^{+} \in B_{\lambda}\right\}, \mathscr{H}_{1}=\varnothing$ and $\mathscr{H}_{k}=\left\{\lambda \in \mathbb{C}^{*} ; f_{\lambda}^{k-2}\left(v_{\lambda}^{+}\right) \in T_{\lambda} \neq B_{\lambda}\right\}$ for $k \geq 2$. See Figure 1. The complement of the escape domains is called the non-escape locus $\mathcal{M}$. It can be written as

$$
\mathcal{M}=\left\{\lambda \in \mathbb{C}^{*} ; f_{\lambda}^{k}\left(v_{\lambda}^{+}\right) \text {does not tend to infinity as } k \rightarrow \infty\right\}
$$

The set $\mathcal{M}$ is invariant under the maps $z \mapsto \bar{z}$ and $z \mapsto e^{\frac{2 \pi i}{n-1}} z$.
Theorem 2.1 (Escape Trichotomy [10] and Connectivity [12]). - We have

1. If $\lambda \in \mathscr{H}_{0}$, then $J\left(f_{\lambda}\right)$ is a Cantor set.
2. If $\lambda \in \mathscr{H}_{2}$, then $J\left(f_{\lambda}\right)$ is a Cantor set of circles.
3. If $\lambda \in \mathscr{H}_{k}$ for some $k \geq 3$, then $J\left(f_{\lambda}\right)$ is a Sierpiński curve.
4. If $\lambda \in \mathcal{M}$, then Julia set $J\left(f_{\lambda}\right)$ is connected.

See Figure 2 for the four typical types of Julia sets.
Based on Theorem 2.1, we give some remarks on the escape domains. According to Devaney, $\mathscr{H}_{0}$ is called the Cantor set locus, $\mathscr{H}_{2}$ is called the McMullen domain, $\mathscr{H}_{k}$ with $k \geq 3$ is called the Sierpiński locus and each of its component is called a Sierpiński hole. Devaney showed that the boundary $\partial \mathscr{H}_{2}$ is a Jordan curve [6] and $\mathscr{H}_{k}$ with $k \geq 3$ consists of $(2 n)^{k-3}(n-1)$ disk components [7].

The Böttcher map $\phi_{\lambda}$ of $f_{\lambda}$ is defined in a neighborhood of $\infty$ by $\phi_{\lambda}(z)=\lim _{k \rightarrow \infty}\left(f_{\lambda}^{k}(z)\right)^{n^{-k}}$. It is unique if we require $\phi_{\lambda}^{\prime}(\infty)=1$. The map $\phi_{\lambda}$ satisfies $\phi_{\lambda}\left(f_{\lambda}(z)\right)=\phi_{\lambda}(z)^{n}$ and $\phi_{\lambda}\left(e^{\pi i / n} z\right)=e^{\pi i / n} \phi_{\lambda}(z)$. One may verify that near infinity,

$$
\phi_{\lambda}(z)=\sum_{k \geq 0} a_{k}(\lambda) z^{1-2 k n}, a_{0}(\lambda)=1, a_{1}(\lambda)=\lambda / n, \ldots
$$

If $\lambda \in \mathbb{C}^{*} \backslash \mathscr{H}_{0}$, then both $B_{\lambda}$ and $T_{\lambda}$ are simply connected. In that case, there is a unique Riemann mapping $\psi_{\lambda}: T_{\lambda} \rightarrow \mathbb{D}$, such that $\psi_{\lambda}(w)^{-n}=\phi_{\lambda}\left(f_{\lambda}(w)\right)$ for $w \in T_{\lambda}$ and $\psi_{\lambda}^{\prime}(0)=\sqrt[n]{\lambda}$. The external ray $R_{\lambda}(t)$ of angle $t$ in $B_{\lambda}$ is defined by $R_{\lambda}(t):=\phi_{\lambda}^{-1}\left((1,+\infty) e^{2 \pi i t}\right)$, the internal ray $R_{T_{\lambda}}(t)$ of angle $t$ in $T_{\lambda}$ is defined by $R_{T_{\lambda}}(t):=\psi_{\lambda}^{-1}\left((0,1) e^{2 \pi i t}\right)$.

Theorem 2.2 (Parameterization of escape domains, [7, 31, 34]). - We have

1. $\mathscr{H}_{0}$ is the unbounded component of $\mathbb{C}^{*}-\mathcal{M}$. The map $\Phi_{0}: \mathcal{H}_{0} \rightarrow \mathbb{C}-\overline{\mathbb{D}}$ defined by $\Phi_{0}(\lambda)=\phi_{\lambda}\left(v_{\lambda}^{+}\right)^{2}$ is a conformal isomorphism.


Figure 2. The Julia sets: a Cantor set (upper-left), a Cantor set of circles (upper-right), a Sierpiński curve (lower-left) and a connected set (lower-right).
2. $\mathscr{H}_{2}$ is the component of $\mathbb{C}^{*}-\mathcal{M}$ containing the punctured neighborhood of 0 . The holomorphic map $\Phi_{2}: \mathscr{H}_{2} \rightarrow \mathbb{C}-\overline{\mathbb{D}}$ defined via $\Phi_{2}(\lambda)^{n-2}=\phi_{\lambda}\left(f_{\lambda}\left(v_{\lambda}^{+}\right)\right)^{2}$ and $\lim _{\lambda \rightarrow 0} \lambda \Phi_{2}(\lambda)=2^{\frac{2 n}{2-n}}$, is a conformal isomorphism.
3. Let $\mathscr{H}$ be an escape domain of level $k \geq 3$. The map $\Phi_{\mathscr{H}}: \mathscr{H} \rightarrow \mathbb{D}$ defined by $\Phi_{\mathscr{H}}(\lambda)=\psi_{\lambda}\left(f_{\lambda}^{k-2}\left(v_{\lambda}^{+}\right)\right)$is a conformal isomorphism.
Both $\Phi_{0}$ and $\Phi_{2}$ satisfy $\Phi_{\epsilon}\left(e^{\frac{2 \pi i}{n-1}} \lambda\right)=e^{\frac{2 \pi i}{n-1}} \Phi_{\epsilon}(\lambda)$ and $\Phi_{\epsilon}(\bar{\lambda})=\overline{\Phi_{\epsilon}(\lambda)}$ for $\epsilon \in\{0,2\}$ and $\lambda \in \mathscr{H}_{\epsilon}$. Thus they take the form $\Phi_{\epsilon}(\lambda)=\lambda \Psi_{\epsilon}\left(\lambda^{n-1}\right)$, where $\Psi_{\epsilon}$ is a holomorphic function whose expansion has real coefficients.

Theorem 2.3 (Connectivity of $\mathcal{M}$ ). - The non-escape locus $\mathcal{M}$ is connected and has logarithmic capacity 1/4.

Proof. - By Theorem 2.2, each component of $\widehat{\mathbb{C}}-\mathcal{M}$ is a topological disk. So $\mathcal{M}$ is connected. The logarithmic capacity of $\mathcal{M}$ follows from the expansion of $\Phi_{0}$ near $\infty$ : $\Phi_{0}(\lambda)=4 \lambda+\theta\left(\lambda^{2-n}\right)$.

## 3. Cut rays in the dynamical plane

The topology of $\partial B_{\lambda}$ is considered in [28], where the authors showed
Theorem 3.1 ([28]). - For any $n \geq 3$ and any $\lambda \in \mathbb{C}^{*}$,

- $\partial B_{\lambda}$ is either a Cantor set or a Jordan curve. In the latter case, all Fatou components eventually mapped to $B_{\lambda}$ are Jordan domains.
- If $\partial B_{\lambda}$ is a Jordan curve containing neither a parabolic point nor the recurrent critical set $C_{\lambda}$, then $\partial B_{\lambda}$ is a quasi-circle.

Here, the critical set $C_{\lambda}$ is called recurrent if $C_{\lambda} \subset J\left(f_{\lambda}\right)$ and the set $\bigcup_{k \geq 1} f_{\lambda}^{k}\left(C_{\lambda}\right)$ has an accumulation point in $C_{\lambda}$. The proof of Theorem 3.1 is based on the Yoccoz puzzle theory. To apply this theory, we need to construct a kind of Jordan curve which cuts the Julia set into two connected parts. These curves are called cut rays. They play a crucial role in our study of the boundaries of escape domains. For this, we briefly sketch their construction here.

To begin, we identify the unit circle $\mathbb{S}=\mathbb{R} / \mathbb{Z}$ with $(0,1]$. We define a map $\tau: \mathbb{S} \rightarrow \mathbb{S}$ by $\tau(\theta)=n \theta \bmod 1$. Let $\Theta_{k}=\left(\frac{k}{2 n}, \frac{k+1}{2 n}\right]$ for $0 \leq k \leq n$ and $\Theta_{-k}=\left(\frac{k}{2 n}+\frac{1}{2}, \frac{k+1}{2 n}+\frac{1}{2}\right]$ for $1 \leq k \leq n-1$. Obviously, $(0,1]=\bigcup_{-n<j \leq n} \Theta_{j}$.

Let $\Theta$ be the set of all angles $\theta \in(0,1]$ whose orbits remain in $\bigcup_{k=1}^{n-1}\left(\Theta_{k} \cup \Theta_{-k}\right)$ under all iterations of $\tau$. One may verify that $\Theta$ is a Cantor set. Given an angle $\theta \in \Theta$, the itinerary of $\theta$ is a sequence of symbols $\left(s_{0}, s_{1}, s_{2}, \ldots\right) \in\{ \pm 1, \ldots, \pm(n-1)\}^{\mathbb{N}}$ such that $\tau^{k}(\theta) \in \Theta_{s_{k}}$ for all $k \geq 0$. The angle $\theta \in \Theta$ and its itinerary $\left(s_{0}, s_{1}, s_{2}, \ldots\right)$ satisfy the identity [28, Lemma 3.1]:

$$
\theta=\frac{1}{2}\left(\frac{\chi\left(s_{0}\right)}{n}+\sum_{k \geq 1} \frac{\left|s_{k}\right|}{n^{k+1}}\right),
$$

where $\chi\left(s_{0}\right)=s_{0}$ if $0 \leq s_{0} \leq n$ and $\chi\left(s_{0}\right)=n-s_{0}$ if $-(n-1) \leq s_{0} \leq-1$.
Note that $e^{\pi i /(n-1)} f_{\lambda}(z)=(-1)^{n} f_{e^{2 \pi i /(n-1)} \lambda}\left(e^{\pi i /(n-1)} z\right)$ for all $\lambda \in \mathbb{C}^{*}$. This implies that the fundamental domain of the parameter plane is

$$
\mathcal{F}_{0}=\left\{\lambda \in \mathbb{C}^{*} ; 0 \leq \arg \lambda<2 \pi /(n-1)\right\} .
$$

We denote the interior of $\mathcal{F}_{0}$ by

$$
\mathscr{F}:=\left\{\lambda \in \mathbb{C}^{*} ; 0<\arg \lambda<2 \pi /(n-1)\right\} .
$$

In our discussion, we assume $\lambda \in \mathcal{F}_{0}$ and let $O_{\lambda}=\bigcup_{k \geq 0} f_{\lambda}^{-k}(\infty)$ be the grand orbit of $\infty$. Let $c_{0}=c_{0}(\lambda)=\sqrt[2 n]{\lambda}$ be the critical point that lies on $\mathbb{R}^{+}:=[0,+\infty)$ when $\lambda \in \mathbb{R}^{+}$and varies analytically as $\lambda$ ranges over $\mathcal{F}$. Let $c_{k}(\lambda)=c_{0} e^{k \pi i / n}$ for $1 \leq k \leq 2 n-1$. The critical points $c_{k}$ with $k$ even are mapped to $v_{\lambda}^{+}$while the critical points $c_{k}$ with $k$ odd are mapped to $v_{\lambda}^{-}$.

Let $\ell_{k}=c_{k}[0,+\infty]$ be the closed straight line connecting 0 to $\infty$ and passing through $c_{k}$ for $0 \leq k \leq 2 n-1$. The closed sector bounded by $\ell_{k}$ and $\ell_{k+1}$ is denoted by $S_{k}^{\lambda}$ for $0 \leq k \leq n$. Define $S_{-k}^{\lambda}=-S_{k}^{\lambda}$ for $1 \leq k \leq n-1$. These sectors are arranged counterclockwise about the origin as $S_{0}^{\lambda}, S_{1}^{\lambda}, \ldots, S_{n}^{\lambda}, S_{-1}^{\lambda}, \ldots, S_{-(n-1)}^{\lambda}$. See [28, Figure 2].

The critical value $v_{\lambda}^{+}$always lies in $S_{0}^{\lambda}$ because $\arg c_{0} \leq \arg v_{\lambda}^{+}<\arg c_{1}$ for all $\lambda \in \mathcal{F}_{0}$. Correspondingly, the critical value $v_{\lambda}^{-}$lies in $S_{n}^{\lambda}$. The image of $\ell_{k}$ under $f_{\lambda}$ is a straight ray connecting one of the critical values to $\infty$; this ray is called a critical value ray. As a consequence, $f_{\lambda}$ maps the interior of each of the sectors of $S_{ \pm 1}^{\lambda}, \ldots, S_{ \pm(n-1)}^{\lambda}$ univalently onto a region $\Upsilon_{\lambda}$, which can be identified as the complex sphere $\widehat{\mathbb{C}}$ minus two critical value rays. For any $\epsilon \in\{ \pm 1, \ldots, \pm(n-1)\}$, let $\operatorname{int}\left(S_{\epsilon}^{\lambda}\right)$ be the interior of $S_{\epsilon}^{\lambda}$, the inverse of $f_{\lambda}: \operatorname{int}\left(S_{\epsilon}^{\lambda}\right) \rightarrow \Upsilon_{\lambda}$ is denoted by $h_{\epsilon}^{\lambda}: \Upsilon_{\lambda} \rightarrow \operatorname{int}\left(S_{\epsilon}^{\lambda}\right)$.

Theorem 3.2 (Cut ray, [8] [28]). - For any $\lambda \in \mathcal{F}$ and any angle $\theta \in \Theta$ with itinerary $\left(s_{0}, s_{1}, s_{2}, \ldots\right)$, the set

$$
\Omega_{\lambda}^{\theta}:=\bigcap_{k \geq 0} f_{\lambda}^{-k}\left(S_{s_{k}}^{\lambda} \cup S_{-s_{k}}^{\lambda}\right)
$$

is a Jordan curve intersecting the Julia set $J\left(f_{\lambda}\right)$ in a Cantor set.

Theorem 3.2 is originally proven for the parameters $\lambda \in \mathcal{F} \cap \mathcal{M}$ in [28]. The proof actually works for all $\lambda \in \mathscr{F}$ without any difference.

Here are some facts about the cut rays:

$$
\begin{gathered}
\Omega_{\lambda}^{\theta}=-\Omega_{\lambda}^{\theta} \text { and } \Omega_{\lambda}^{\theta}=\Omega_{\lambda}^{\theta+1 / 2} \\
R_{\lambda}(\theta) \cup R_{\lambda}\left(\theta+\frac{1}{2}\right) \subset \Omega_{\lambda}^{\theta} \cap F\left(f_{\lambda}\right) \subset \bigcup_{k \geq 0} f_{\lambda}^{-k}\left(B_{\lambda}\right)
\end{gathered}
$$

$0, \infty \in \Omega_{\lambda}^{\theta}$ and $\Omega_{\lambda}^{\theta} \backslash\{0, \infty\}$ is contained in the interior of $S_{s_{0}}^{\lambda} \cup S_{-s_{0}}^{\lambda} ; f_{\lambda}\left(\Omega_{\lambda}^{\theta}\right)=\Omega_{\lambda}^{\tau(\theta)}$ and $f_{\lambda}: \Omega_{\lambda}^{\theta} \rightarrow \Omega_{\lambda}^{\tau(\theta)}$ is a two-to-one map. We refer the reader to [28] for more details of the cut rays.


Figure 3. Some cut rays $\Omega_{\lambda}^{\theta}$ with $\theta=1 / 4,1 / 3,1 / 2 .(n=3)$

Now we give some new dynamical properties of the cut rays. These facts will be useful to study the parameter plane. We denote by $B(z, r)$ the open Euclidean disk centered at $z$ with radius $r$. For any $\lambda \in \mathbb{C}^{*} \backslash \mathscr{H}_{0}$, set $B_{\lambda}^{L}:=\left\{w \in B_{\lambda} ;\left|\phi_{\lambda}(w)\right|>L\right\}$ for $L \geq 1$.

Lemma 3.3 (Holomorphic motion of the cut rays). - Fix an angle $\theta \in \Theta$, the cut ray $\Omega_{\lambda}^{\theta}$ moves holomorphically with respect to $\lambda \in \mathscr{F}$.

Proof. - Fix a parameter $\lambda_{0} \in \mathcal{F}$. We will define a holomorphic motion $h: \mathcal{F} \times\left(\left(\Omega_{\lambda_{0}}^{\theta} \backslash O_{\lambda_{0}}\right) \cap F\left(f_{\lambda_{0}}\right)\right) \rightarrow \widehat{\mathbb{C}}$ with base point $\lambda_{0}$ as follows. For any $\lambda \in \mathcal{F}$, there is a number $L \geq 1$ (depending on $\lambda$ ) such that the Böttcher map $\phi_{u}: B_{u}^{L} \rightarrow\{\zeta \in \widehat{\mathbb{C}} ;|\zeta|>L\}$ is a conformal isomorphism, for $u \in\left\{\lambda, \lambda_{0}\right\}$.

If $z \in\left(\Omega_{\lambda_{0}}^{\theta} \backslash O_{\lambda_{0}}\right) \cap B_{\lambda_{0}}^{L}$, we define $h(\lambda, z)=\phi_{\lambda}^{-1} \circ \phi_{\lambda_{0}}(z)$. If $z \in\left(\Omega_{\lambda_{0}}^{\theta} \backslash O_{\lambda_{0}}\right) \cap$ $\left(F\left(f_{\lambda_{0}}\right) \backslash B_{\lambda_{0}}^{L}\right)$, we consider the itinerary of $z$, which is the unique sequence of symbols $\left(\epsilon_{0}, \epsilon_{1}, \epsilon_{2}, \ldots\right) \in\{ \pm 1, \ldots, \pm(n-1)\}^{\mathbb{N}}$ such that $f_{\lambda_{0}}^{k}(z) \in S_{\epsilon_{k}}^{\lambda_{0}}$ for all $k \geq 0$. Let $N \geq 1$ be the first integer such that $f_{\lambda_{0}}^{N}(z) \in B_{\lambda_{0}}^{L}$. We define $h(\lambda, z)=h_{\epsilon_{0}}^{\lambda} \circ \cdots \circ h_{\epsilon_{N-1}}^{\lambda} \circ \phi_{\lambda}^{-1} \circ \phi_{\lambda_{0}}\left(f_{\lambda_{0}}^{N}(z)\right)$. In this way, we get a well-defined map $h: \mathscr{F} \times\left(\left(\Omega_{\lambda_{0}}^{\theta} \backslash O_{\lambda_{0}}\right) \cap F\left(f_{\lambda_{0}}\right)\right) \rightarrow \mathbb{\mathbb { C }}$. Since both $\phi_{\lambda}$ and $h_{\epsilon_{j}}^{\lambda}$ are holomorphic with respect to $\lambda \in \mathscr{F}$, one may verify that the map $h$ is a holomorphic motion parameterized by $\mathcal{F}$, with base point $\lambda_{0}$ (namely, $h\left(\lambda_{0}, z\right) \equiv z$ ). Moreover, for any $\lambda \in \mathcal{F}$, we have $h\left(\lambda,\left(\Omega_{\lambda_{0}}^{\theta} \backslash O_{\lambda_{0}}\right) \cap F\left(f_{\lambda_{0}}\right)\right)=\left(\Omega_{\lambda}^{\theta} \backslash O_{\lambda}\right) \cap F\left(f_{\lambda}\right)$.

Note that for any $\lambda \in \mathcal{F}$, the closure of $\left(\Omega_{\lambda}^{\theta} \backslash O_{\lambda}\right) \cap F\left(f_{\lambda}\right)$ is $\Omega_{\lambda}^{\theta}$. By the $\lambda$-Lemma (see [21] or [24]), there is a holomorphic motion $H: \mathscr{F} \times \Omega_{\lambda_{0}}^{\theta} \rightarrow \widehat{\mathbb{C}}$ extending $h$ and for any $\lambda \in \mathcal{F}$, one has $H\left(\lambda, \Omega_{\lambda_{0}}^{\theta}\right)=\Omega_{\lambda}^{\theta}$. That is to say, the cut ray $\Omega_{\lambda}^{\theta}$ moves holomorphically when $\lambda$ ranges over $\mathcal{F}$.

Let $\Theta_{p e r}$ be a subset of $\Theta \backslash\{1,1 / 2\}$, consisting of all periodic angles under the map $\tau$. One may verify that $\Theta_{p e r}$ is a dense subset of $\Theta$.

Theorem 3.4 (Cut rays with real parameters). - For any $\lambda \in(0,+\infty)$ and any angle $\theta \in \Theta_{\text {per }}$ with itinerary $\left(s_{0}, s_{1}, s_{2}, \ldots\right)$, the set

$$
\Omega_{\lambda}^{\theta}:=\bigcap_{k \geq 0} f_{\lambda}^{-k}\left(\left(S_{s_{k}}^{\lambda} \cup S_{-s_{k}}^{\lambda}\right) \backslash \mathbb{R}^{*}\right)
$$

is a Jordan curve intersecting the Julia set $J\left(f_{\lambda}\right)$ in a Cantor set, where $\mathbb{R}^{*}=\mathbb{R} \backslash\{0\}$. Moreover, if $\mathcal{F}_{0} \ni \lambda_{j} \rightarrow \lambda \in(0,+\infty)$, then $\Omega_{\lambda_{j}}^{\theta} \rightarrow \Omega_{\lambda}^{\theta}$ in Hausdorff topology.

Here is a remark. If $\lambda \in \mathcal{F}$, then $\bigcap_{k \geq 0} f_{\lambda}^{-k}\left(S_{s_{k}}^{\lambda} \cup S_{-s_{k}}^{\lambda}\right)=\bigcap_{k \geq 0} f_{\lambda}^{-k}\left(\left(S_{S_{k}}^{\lambda} \cup S_{-s_{k}}^{\lambda}\right) \backslash \mathbb{R}^{*}\right)$, so the latter is also a reasonable definition of cut rays. However, if $\lambda \in(0,+\infty)$, the set $\bigcap_{k \geq 0} f_{\lambda}^{-k}\left(S_{s_{k}}^{\lambda} \cup S_{-s_{k}}^{\lambda}\right)$ is not a Jordan curve in general.

The proof of Theorem 3.4 is essentially the same as that of Proposition 3.9 in [28]. We would like to mention the idea of the proof here. Let $Y_{\lambda}=\widehat{\mathbb{C}} \backslash\left(\left[-\infty, v_{\lambda}^{+}\right] \cup\left[v_{\lambda}^{+},+\infty\right] \cup \overline{B_{\lambda}^{L}}\right)$ for some large $L>1$ and $p$ be the period of $\theta$. The itinerary of $\theta$ satisfies $s_{p+k}=s_{k}$ for all $k \geq 0$. Since $\theta \neq 1,1 / 2$, one of $s_{k}$ will be in the set $\{ \pm 1, \ldots, \pm(n-2)\}$ and for any $k \geq 0$ and any $\left(\epsilon_{1}, \ldots, \epsilon_{p}\right)=\left( \pm s_{k}, \cdots, \pm s_{k+p-1}\right)$, the set $h_{\epsilon_{1}}^{\lambda} \circ \cdots \circ h_{\epsilon_{p}}^{\lambda}\left(Y_{\lambda}\right)$ is compactly contained in $Y_{\lambda}$ (one should note that if $\theta=1$ or $1 / 2$, then $h_{1}^{\lambda} \circ \cdots \circ h_{1}^{\lambda}\left(Y_{\lambda}\right)$ is not compactly contained in $Y_{\lambda}$ ). Similar to the proof of Proposition 3.9 in [28], one can construct two sequences of Jordan curves converging to the boundaries of the two components of $\widehat{\mathbb{C}}-\Omega_{\lambda}^{\theta}$. In this way $\Omega_{\lambda}^{\theta}$ is locally connected. One can show that its two complementary components share the same boundary, so $\Omega_{\lambda}^{\theta}$ is a Jordan curve.

Using the same proof as Lemma 3.3, one can show that if $\theta \in \Theta_{\text {per }}$, then the cut ray $\Omega_{\lambda}^{\theta}$ is a holomorphic motion in a neighborhood of the real and positive axis. This yields the continuity of cut rays. We omit the details.

Proposition 3.5 (Preimages of cut ray, [28], Prop 3.5). - For any $\lambda \in \mathscr{F}_{0}$ and any $\theta \in \Theta_{\text {per }}$, suppose that $\left(\Omega_{\lambda}^{\theta}-\{0, \infty\}\right) \cap\left(\bigcup_{1 \leq k \leq N} f_{\lambda}^{k}\left(C_{\lambda}\right)\right)=\varnothing$ for some $N \geq 1$. Then, for any $\alpha \in \bigcup_{0 \leq k \leq N} \tau^{-k}(\theta)$, there is a unique Jordan curve $\Omega_{\lambda}^{\alpha}$ (or $\Omega_{\lambda}^{\alpha+1 / 2}$ ) containing 0 and $\infty$, such that $f_{\lambda}\left(\Omega_{\lambda}^{\alpha}\right)=\Omega_{\lambda}^{\tau(\alpha)}$ and $R_{\lambda}(\alpha) \cup R_{\lambda}(\alpha+1 / 2) \subset \Omega_{\lambda}^{\alpha} \cap B_{\lambda}$.

The Jordan curve $\Omega_{\lambda}^{\alpha}$ defined in Proposition 3.5 is also called a cut ray. We remark that the statement of Proposition 3.5 is slightly different from Prop. 3.5 in [28], but their proofs are the same.

Remark 3.6. - The cut ray $\Omega_{\lambda}^{\alpha}$ defined by Proposition 3.5 satisfies the following property:
There is a neighborhood $\mathcal{U}$ of $\lambda$, such that for all $u \in \mathscr{U} \cap \mathscr{F}_{0},\left(\Omega_{u}^{\theta}-\{0, \infty\}\right) \cap$ $\left(\bigcup_{1 \leq k \leq N} f_{u}^{k}\left(C_{u}\right)\right)=\varnothing$ (this implies the cut ray $\Omega_{u}^{\alpha}$ exists). By Lemma 3.3 and Theorem 3.4, the cut ray $\Omega_{u}^{\alpha}$ moves continuously with respect to $u \in U \cap \mathscr{F}_{0}$.

Lemma 3.7. - For any $\lambda \in \mathcal{F}_{0}$ and any two different external rays $R_{\lambda}\left(t_{1}\right)$ and $R_{\lambda}\left(t_{2}\right)$, there is a cut ray $\Omega_{\lambda}^{\alpha}$ with $\alpha \in \bigcup_{k \geq 0} \tau^{-k}\left(\Theta_{\text {per }}\right)$ separating them.

Proof. - Since $\Theta_{\text {per }}$ is an infinite set, we can find an angle $\theta \in \Theta_{\text {per }}$ such that $\left(\Omega_{\lambda}^{\theta}-\{0, \infty\}\right) \cap\left(\bigcup_{k \geq 1} f_{\lambda}^{k}\left(C_{\lambda}\right)\right)=\varnothing$. The preimages $\bigcup_{k \geq 0} \tau^{-k}(\theta)$ of $\theta$ are dense in the unit circle, so there is $\alpha \in \bigcup_{k \geq 0} \tau^{-k}(\theta)$ lying in between $t_{1}$ and $t_{2}$. Then $R_{\lambda}\left(t_{1}\right)$ and $R_{\lambda}\left(t_{2}\right)$ are contained in different components of $\widehat{\mathbb{C}}-\Omega_{\lambda}^{\alpha}$.

## 4. Maps on $\partial \mathcal{H}_{0}$

In this section, we give some a priori characterizations of the maps on $\partial \mathscr{H}_{0}$. We begin with a dynamical result for our purpose. To prove Theorem 3.1 in [28], we reduce the situation to the following:

Theorem 4.1 (Backward contraction on $\left.\partial B_{\lambda},[28]\right)$. - Suppose that $\lambda \in \mathbb{C}^{*} \backslash \mathscr{H}_{0}$ and $\partial B_{\lambda}$ contains neither a parabolic point nor the recurrent critical set $C_{\lambda}$, then $f_{\lambda}$ satisfies the following property on $\partial B_{\lambda}$ : there exist three constants $\delta_{0}>0, C>0$ and $0<\rho<1$ such that for any $0<\delta<\delta_{0}$, any $z \in \partial B_{\lambda}$, any integer $k \geq 0$ and any component $U_{k}(z)$ of $f_{\lambda}^{-k}(B(z, \delta))$ that intersects with $\partial B_{\lambda}, U_{k}(z)$ is simply connected with Euclidean diameter $\operatorname{diam}\left(U_{k}(z)\right) \leq C \delta \rho^{k}$.

We refer the reader to [28] for a detailed proof based on the Yoccoz puzzle theory (to obtain Theorem 4.1, one should combine two results in [28]: Theorem 1.2 in Section 7.5 and Proposition 6.1 in Section 6). The proof of Theorem 4.1 rests on the crucial fact that the puzzle pieces around each point on $\partial B_{\lambda}$ shrink to a single point.

Lemma 4.2. - Suppose that $J\left(f_{\lambda}\right)$ is not a Cantor set. If $\partial B_{\lambda}$ contains neither a critical point nor a parabolic cycle, then there exist an integer $k \geq 1$ and two topological disks $U_{\lambda}, V_{\lambda}$ with $\overline{B_{\lambda}} \subset V_{\lambda} \subset U_{\lambda}$, such that $f_{\lambda}^{k}: V_{\lambda} \rightarrow U_{\lambda}$ is a polynomial-like map of degree $n^{k}$ with only one critical point $\infty$.
$4^{\mathrm{e}}$ SÉRIE - TOME 48 - 2015 - No 3

Proof. - The map $f_{\lambda}$ satisfies the assumptions in Theorem 4.1. This guarantees the existence of three constants $\delta_{0}, C, \rho$.

Let $N_{\delta}$ be $\delta$-neighborhood of $\partial B_{\lambda}$, defined as the set of all points whose Euclidean distance to $\partial B_{\lambda}$ is smaller than $\delta$. We choose an integer $\ell>0$ and a number $\delta<\delta_{0}$ such that $C \rho^{\ell}<1$ and $\left(\bigcup_{0 \leq j<\ell} f_{\lambda}^{-j}\left(C_{\lambda}\right)\right) \cap N_{\delta}=\varnothing$.

Given a Jordan curve $\gamma$, we define its partial distance to $\partial B_{\lambda}$ by $\varpi(\gamma):=\max _{z \in \gamma} d\left(z, \partial B_{\lambda}\right)$, where $d(\cdot, \cdot)$ is Euclidean distance. We choose a Jordan curve $\gamma_{0} \subset \widehat{\mathbb{C}} \backslash \bar{B}_{\lambda}$ with $\varpi\left(\gamma_{0}\right)<\delta$. The annulus between $\gamma_{0}$ and $\partial B_{\lambda}$ is denoted by $A_{0}$. Since $\left(\bigcup_{0 \leq j<\ell} f_{\lambda}^{-j}\left(C_{\lambda}\right)\right) \cap N_{\delta}=\varnothing$, there is an annular component of $f_{\lambda}^{-\ell}\left(A_{0}\right)$, say $A_{1}$, with $\partial B_{\lambda}$ as one of its boundary components. The other boundary curve is denoted by $\gamma_{1}$. Theorem 4.1 implies $\varpi\left(\gamma_{1}\right) \leq \varpi\left(\gamma_{0}\right) C \rho^{\ell}<\delta$. Continuing inductively, for any $k \geq 1$, there is an annular component of $f_{\lambda}^{-\ell}\left(A_{k-1}\right)$, say $A_{k}$, whose boundary curves are $\partial B_{\lambda}$ and $\gamma_{k}$. Then we have

$$
\varpi\left(\gamma_{k}\right) / \varpi\left(\gamma_{0}\right) \leq C \rho^{k \ell} .
$$

So we can choose $k_{0}>0$ such that $\varpi\left(\gamma_{k_{0}}\right)<\min _{z \in \gamma_{0}} d\left(z, \partial B_{\lambda}\right)$. Let $V_{\lambda}$ be the unbounded component of $\widehat{\mathbb{C}}-\gamma_{k_{0}}$ and $U_{\lambda}$ be the unbounded component of $\widehat{\mathbb{C}}-\gamma_{0}$. Then $f_{\lambda}^{k_{0} \ell}: V_{\lambda} \rightarrow U_{\lambda}$ is a polynomial-like map of degree $n^{k_{0} \ell}$, with only one critical point $\infty$. It is actually quasiconformally conjugate to the power map $z \mapsto z^{n^{k_{0} \ell}}$.

We give here an alternative proof of Lemma 4.2 which uses Mañe' Lemma. This proof does not need the precise estimate of Theorem 4.1.

Proof of Lemma 4.2. - From [28] it follows that $\partial B_{\lambda}$ is a Jordan curve. Therefore, $\overline{B_{\lambda}}$ is full, i.e., $\widehat{\mathbb{C}} \backslash \overline{B_{\lambda}}$ is connected. Since there is no critical point on the boundary $\partial B_{\lambda}$, then $\overline{B_{\lambda}} \cap \overline{T_{\lambda}}=\varnothing$. Therefore there exists an open disk $U$ containing $\overline{B_{\lambda}}$, avoiding $\overline{T_{\lambda}}$ and the critical set. Let $\phi$ be a Riemann map from $\widehat{\mathbb{C}} \backslash \overline{B_{\lambda}}$ to $\widehat{\mathbb{C}} \backslash \overline{\mathbb{D}}$. The map $\phi \circ f_{\lambda} \circ \phi^{-1}$ defined on $\phi\left(U \backslash \overline{B_{\lambda}}\right)$ extends by the Schwarz reflexion principle to a map noted by $\tilde{f}$ on a neighborhood of the unit circle. Let $g$ denote the restriction of $\tilde{f}$ to the unit circle. It is a covering of degree $n$, the degree of the map $f_{\lambda}$ on $U \backslash \overline{B_{\lambda}}$.

Let us prove by contradiction that $g$ has no non-repelling periodic point. In the case of existence of an attracting or parabolic cycle of $g$ on $\partial \mathbb{D}$, choose a point $a$ in the cycle, let $m$ be its period and let $\mathscr{G}$ be its immediate basin. Let $\mathscr{B}=\phi^{-1}(\mathscr{G})=\phi^{-1}(\mathscr{G} \cap(\mathbb{C} \backslash \overline{\mathbb{D}}))$. Then $\mathscr{B}$ is bounded in $\mathbb{C}$ and stable by $f_{\lambda}^{m p}$. It is therefore contained in a periodic component $\mathscr{B}^{\prime}$ of the Fatou set.

Since every point of $\mathscr{C}$ tends to $a$ under iteration of $\tilde{f}^{m p}$, every point of $\mathscr{B}$ tends to $\partial B_{\lambda}$ under iteration of $f_{\lambda}^{m p}$. Thus in particular the component $\mathcal{B}$ is not part of a cycle of an attracting point nor a cycle of Siegel disks, because every point in a Siegel disk has an orbit that stays bounded away from the Julia set. It is therefore in a component of an immediate parabolic basin which contradicts the assumption that there is no parabolic point on $\partial B_{\lambda}$.

Let us cite Theorem A in [19] (see also [20]).
Theorem 4.3 (Mañé). - Let $N=S^{1}$ (the circle) or $N=[0,1]$. If f is a $C^{2}$ map from $N$ to $N$ and $\Lambda \subset N$ is a compact invariant set that does not contain critical points, sinks or nonhyperbolic periodic points, then either $\Lambda=N=S^{1}$ and $f$ is topologically equivalent to a rotation or $\Lambda$ is a hyperbolic set.

By the preceding discussion, we can apply this to $g$ on $\partial \mathbb{D}$ with $\Lambda=N=\partial \mathbb{D}$. So $g$, and therefore $\widetilde{f}$, is hyperbolic on $\partial \mathbb{D}$, which means that there is a continuous function $\rho(z)>0$ on the unit circle such that $\tilde{f}$ is uniformly expanding with respect to the metric $\rho(z)|d z|$, i.e., $\rho(\widetilde{f}(z))\left|\widetilde{f}^{\prime}(z)\right|>\kappa \rho(z)$ for some $\kappa>1$. Consider for some $\epsilon>0$ the domain

$$
V=\left\{z=r e^{i \theta}| | \log r \mid \leq \epsilon \rho\left(e^{i \theta}\right)\right\} .
$$

Let $U$ be the connected component of $\tilde{f}^{-1}(V)$ containing the unit circle. For $\epsilon$ small enough, $U$ is compactly contained in $V$.

Let $U^{\prime}=\overline{B_{\lambda}} \cup \phi^{-1}(U)$ and $V^{\prime}=\overline{B_{\lambda}} \cup \phi^{-1}(V)$. Then $U^{\prime}$ and $V^{\prime}$ are open, connected, simply connected, $U^{\prime}$ is a connected component of $f_{\lambda}^{-p}\left(V^{\prime}\right)$ and $U^{\prime}$ is compactly contained in $V^{\prime}$. Therefore the restriction of $f_{\lambda}^{p}: U^{\prime} \rightarrow V^{\prime}$ is a polynomial-like map. This restriction is also unicritical.

Lemma 4.4. - Suppose $\lambda \in \partial \mathscr{H}_{0}$, then $\partial B_{\lambda}$ contains either the critical set $C_{\lambda}$ or a parabolic cycle of $f_{\lambda}$.

Proof. - If $\partial B_{\lambda}$ contains neither the critical set $C_{\lambda}$ nor a parabolic cycle, then it follows from Lemma 4.2 that there exist an integer $k \geq 1$ and two topological disks $U_{\lambda}, V_{\lambda}$ with $\overline{B_{\lambda}} \subset V_{\lambda} \subset U_{\lambda}$, such that $f_{\lambda}^{k}: V_{\lambda} \rightarrow U_{\lambda}$ is a polynomial-like map of degree $n^{k}$ with only one critical point $\infty$. We may assume that $\overline{U_{\lambda}}$ has no intersection with $\bigcup_{0 \leq j<k} f_{\lambda}^{-j}\left(C_{\lambda}\right)$.

Then there is a neighborhood of $\mathcal{U}$ of $\lambda$, such that for all $u \in \mathcal{U}$, the set $\bigcup_{0 \leq j<k} f_{u}^{-j}\left(C_{u}\right)$ has no intersection with $\overline{U_{\lambda}}$, thus the component $V_{u}$ of $f_{u}^{-k}\left(U_{\lambda}\right)$ that contains $\infty$ is a disk. Since $\partial V_{u}$ moves holomorphically with respect to $u \in \mathcal{U}$, we may shrink $\mathcal{U}$ a little bit so that for all $u \in \mathcal{U}, \partial V_{u}$ is contained in $U_{\lambda}$. Set $U_{u}=U_{\lambda}$. In this way, we get a polynomial-like map $f_{u}^{k}: V_{u} \rightarrow U_{u}$ with only one critical point $\infty$, for all $u \in \mathcal{U}$. As a consequence, the Julia set $J\left(f_{u}\right)$ is not a Cantor set for $u \in \mathcal{U}$.

But this is impossible since $\lambda \in \partial \mathscr{H}_{0}$.
Given a parameter $\lambda \in \mathcal{F}$, if $C_{\lambda} \subset \partial B_{\lambda}$, then there is a unique external ray $R_{\lambda}(t)$ landing at $v_{\lambda}^{+}$. We define $\theta(\lambda)=t$. Note that $C_{\lambda} \subset \partial B_{\lambda}$ if and only if $v_{\lambda}^{+} \in \partial B_{\lambda}$.

Lemma 4.5. - If $\lambda \in \mathscr{F}$ and $v_{\lambda}^{+} \in \partial B_{\lambda}$, then $0<\theta(\lambda)<\frac{1}{2(n-1)}$.
Proof. - If $\lambda \in \mathcal{F}$, then $v_{\lambda}^{+}$is contained in the interior of the closed sector $S_{0}^{\lambda}$. Note that $\Omega_{\lambda}^{1} \subset S_{n-1}^{\lambda} \cup S_{-(n-1)}^{\lambda}$ and $\Omega_{\lambda}^{\frac{1}{2(n-1)}} \subset S_{1}^{\lambda} \cup S_{-1}^{\lambda}$, we have $0<\theta(\lambda)<\frac{1}{2(n-1)}$.

Lemma 4.6 ([28], Proposition 7.5). - If $\partial B_{\lambda}$ contains a parabolic cycle, then the following hold:

1. There is a symbol $\epsilon \in\{ \pm 1\}$, an integer $p \geq 1$, a critical point $c \in C_{\lambda}$ and two topological disks $U$ and $V$ containing $c$, such that $\epsilon f_{\lambda}^{p}: U \rightarrow V$ is a quadratic-like map, hybrid equivalent to the polynomial $z \mapsto z^{2}+1 / 4$.
2. Let $K$ be the filled Julia set of $\epsilon f_{\lambda}^{p}: U \rightarrow V$, then for any $j \geq 0$, the intersection $f_{\lambda}^{j}(K) \cap \partial B_{\lambda}$ is a singleton.

Based on Lemma 4.6, let $K^{+} \in\left\{f_{\lambda}(K),-f_{\lambda}(K)\right\}$ be the set containing $v_{\lambda}^{+}$, and $\beta_{\lambda}$ be the intersection point of $K^{+}$and $\partial B_{\lambda}$.
$4^{\mathrm{e}}$ SÉRIE - TOME 48 - 2015 - No 3

REMARK 4.7. - If $n$ is odd, since $f_{\lambda}$ is an odd function, $\beta_{\lambda}$ is necessarily a parabolic point; if $n$ is even, either $\beta_{\lambda}$ or $-\beta_{\lambda}$ is a parabolic point. Thus $\theta$ satisfies either $\tau^{p}(\theta) \equiv \theta$ or $\tau^{p}(\theta) \equiv \theta+\frac{1}{2}$ for some $p \geq 1$.

For any $t \in[0,1)$, the parameter ray $\mathcal{R}_{0}(t)$ of angle $t$ in $\mathcal{H}_{0}$ is defined by $\mathcal{R}_{0}(t):=\Phi_{0}^{-1}\left((1,+\infty) e^{2 \pi i t}\right)$. Its impression $\chi_{t}$ is defined by

$$
\left.\chi_{t}:=\bigcap_{k \geq 1} \overline{\Phi_{0}^{-1}\left(\left\{r e^{2 \pi i \theta} ; 1<r<1+1 / k,|\theta-t|<1 / k\right\}\right.}\right)
$$

The set $\chi_{t}$ is a connected and compact subset of $\partial \mathscr{H}_{0}$. It satisfies

$$
\chi_{t+\frac{1}{n-1}}=e^{2 \pi i /(n-1)} \chi_{t},\left\{\bar{\lambda} ; \lambda \in \chi_{t}\right\}=\chi_{1-t}
$$

Lemma 4.8. - Let $t \in\left[0, \frac{1}{n-1}\right)$ and $\lambda \in \chi_{t} \cap \mathcal{F}_{0}$.

1. If $\lambda$ is not a cusp, then the external ray $R_{\lambda}(t / 2)$ lands at $v_{\lambda}^{+}$.
2. If $\lambda$ is a cusp, then the external ray $R_{\lambda}(t / 2)$ lands at $\beta_{\lambda}$.

Proof. - For any parameter $\lambda \in \chi_{t} \cap \mathcal{F}_{0}$, it follows from Lemma 4.4 that either $C_{\lambda} \subset \partial B_{\lambda}$ or $\partial B_{\lambda}$ contains a parabolic cycle. Since $\partial B_{\lambda}$ is a Jordan curve (Theorem 3.1), there is an external ray $R_{\lambda}\left(t^{\prime}\right)$ landing at $v_{\lambda}^{+}$(if $\lambda$ is not a cusp) or $\beta_{\lambda}$ (if $\lambda$ is a cusp).

If $t^{\prime} \notin\{t / 2,(1+t) / 2\}$, then there exist two cut rays $\Omega_{\lambda}^{\alpha}$ and $\Omega_{\lambda}^{\beta}$ with $\alpha, \beta \in \bigcup_{k \geq 0} \tau^{-k}\left(\Theta_{p e r}\right)$ (Lemma 3.7) such that the connected set $R_{\lambda}\left(t^{\prime}\right) \cup\left\{v_{\lambda}^{+}\right\}$(if $\lambda$ is not a cusp) or $R_{\lambda}\left(t^{\prime}\right) \cup K^{+}$ (if $\lambda$ is a cusp), and the external rays $R_{\lambda}(t / 2), R_{\lambda}((t+1) / 2)$ are contained in three different components of $\widehat{\mathbb{C}} \backslash\left(\Omega_{\lambda}^{\alpha} \cup \Omega_{\lambda}^{\beta}\right)$. See Figure 4. Since the critical value $v_{u}^{+}=2 \sqrt{u}$ and the


Figure 4. Two cut rays $\Omega_{\lambda}^{\alpha}$ and $\Omega_{\lambda}^{\beta}$ separate the external rays $R_{\lambda}\left(t^{\prime}\right), R_{\lambda}(t / 2), R_{\lambda}((t+1) / 2)$ in case that $\lambda$ is not a cusp.
cut rays $\Omega_{u}^{\alpha}, \Omega_{u}^{\beta}$ move continuously with respect to the parameter $u \in \mathcal{F}_{0}$ (Lemma 3.3 and Remark 3.6), there is a neighborhood $V$ of $\lambda$ such that for all $u \in \mathcal{V} \cap \mathcal{F}_{0}$,

- $R_{u}\left(t^{\prime}\right)$ and $v_{u}^{+}$are contained in the same component of $\widehat{\mathbb{C}} \backslash\left(\Omega_{u}^{\alpha} \cup \Omega_{u}^{\beta}\right)$.
- The external rays $R_{u}\left(t^{\prime}\right), R_{u}(t / 2), R_{u}((t+1) / 2)$ are contained in three different components of $\widehat{\mathbb{C}} \backslash\left(\Omega_{u}^{\alpha} \cup \Omega_{u}^{\beta}\right)$.

By shrinking $V$ a little bit, we see that there is a small number $\varepsilon>0$ such that $\arg \Phi_{0}(u)=2 \arg \phi_{u}\left(v_{u}^{+}\right) \notin(t-\varepsilon, t+\varepsilon)$ for all $u \in \mathcal{V} \cap \mathcal{F}_{0} \cap \mathscr{H}_{0}$. This is a contradiction since $\lambda \in \chi_{t}$.

So either $t^{\prime}=t / 2$ or $t^{\prime}=(1+t) / 2$. To finish, we show that the latter is impossible. If $\lambda \in(0,+\infty)$, then $\partial B_{\lambda}$ contains a cusp and $t^{\prime}=0$. If $\lambda \in \mathcal{F}$, then there is a component $V$ of $\widehat{\mathbb{C}} \backslash\left(\Omega_{\lambda}^{1} \cup \Omega_{\lambda}^{\frac{1}{2(n-1)}}\right)$ such that $v_{\lambda}^{+} \cup R_{\lambda}\left(t^{\prime}\right) \subset \bar{V}$. In this case, we have $0 \leq t \leq \frac{1}{2(n-1)}$.

So $t^{\prime}=t / 2$.

## 5. A rigidity result

The main result of this section is the following:

Theorem 5.1. - Given two parameters $\lambda_{1}, \lambda_{2} \in \mathcal{F}$, if $v_{\lambda_{i}}^{+} \in \partial B_{\lambda_{i}}(i=1,2)$ and $\theta\left(\lambda_{1}\right)=\theta\left(\lambda_{2}\right)$, then $\lambda_{1}=\lambda_{2}$.

Recall that given a parameter $\lambda \in \mathcal{F}$ with $C_{\lambda} \subset \partial B_{\lambda}$, the angle $\theta(\lambda)$ is defined such that the external ray $R_{\lambda}(\theta(\lambda))$ lands at $v_{\lambda}^{+}$. Theorem 5.1 is a crucial step in proving that $\partial \mathscr{H}_{0}$ is a Jordan curve in the next section.

When $\theta\left(\lambda_{1}\right)$ is a rational number, the proof is based on Thurston's Theorem [13], as follows:

Proof of Theorem 5.1 when $\theta\left(\lambda_{1}\right)$ is a rational number. - In this case, both $f_{\lambda_{1}}$ and $f_{\lambda_{2}}$ are postcritically finite. We define a homeomorphism $\psi: \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$ such that $\left.\psi\right|_{B_{\lambda_{1}}}=\phi_{\lambda_{2}}^{-1} \circ \phi_{\lambda_{1}}$. Then there is a homeomorphism $\varphi: \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$ satisfying $\psi \circ f_{\lambda_{1}}=f_{\lambda_{2}} \circ \varphi$ and $\left.\varphi\right|_{B_{\lambda_{1}}}=\left.\psi\right|_{B_{\lambda_{1}}}$. (In fact, $\varphi$ and $\psi$ can be made quasiconformal because $\partial B_{\lambda_{1}}$ and $\partial B_{\lambda_{2}}$ are quasi-circles, see Theorem 3.1.) The condition $\theta\left(\lambda_{1}\right)=\theta\left(\lambda_{2}\right)$ implies that $\varphi$ and $\psi$ are isotopic relative to the postcritical set $P\left(f_{\lambda_{1}}\right):=\{\infty\} \cup \bigcup_{k \geq 1} f_{\lambda_{1}}^{k}\left(C_{\lambda_{1}}\right)$. Thus $f_{\lambda_{1}}$ and $f_{\lambda_{2}}$ are combinatorially equivalent. It follows from Thurston's theorem (see [13]) that $f_{\lambda_{1}}$ and $f_{\lambda_{2}}$ are conjugate via a Möbius transformation. This Möbius map necessarily takes the form $\gamma(z)=a z$ with $a^{n-1}=1$ and $\lambda_{2}=a^{2} \lambda_{1}$. The condition $\lambda_{1}, \lambda_{2} \in \mathscr{F}$ implies $\lambda_{1}=\lambda_{2}$.

When $\theta\left(\lambda_{1}\right)$ is an irrational number, The proof of Proposition 5.1 involves the Yoccoz puzzle theory. We first recall the Yoccoz puzzle construction in [28].
$4^{\mathrm{e}}$ SÉRIE - TOME 48 - 2015 - No 3

### 5.1. The Yoccoz puzzle

Given a parameter $\lambda \in \mathcal{M} \cap \mathcal{F}$, we define a graph $G_{\lambda}\left(\theta_{1}, \ldots, \theta_{N}\right)$ by

$$
G_{\lambda}\left(\theta_{1}, \ldots, \theta_{N}\right)=\partial B_{\lambda}^{L} \cup\left(\left(\widehat{\mathbb{C}} \backslash B_{\lambda}^{L}\right) \cap \bigcup_{k \geq 0}\left(\Omega_{\lambda}^{\tau^{k}\left(\theta_{1}\right)} \cup \cdots \cup \Omega_{\lambda}^{\tau^{k}\left(\theta_{N}\right)}\right)\right)
$$

where $L>1$ is a fixed number and $\theta_{1}, \ldots, \theta_{N} \in \Theta$ are $\tau$-periodic angles. The angles $\theta_{1}, \ldots, \theta_{N}$ are chosen so that the free critical orbit $\bigcup_{k \geq 1} f_{\lambda}^{k}\left(C_{\lambda}\right)$ avoids the graph. The puzzle pieces of depth $d \geq 0$ are defined to be all the connected components of $f_{\lambda}^{-d}\left(\left(\widehat{\mathbb{C}} \backslash B_{\lambda}^{L}\right) \backslash G_{\lambda}\left(\theta_{1}, \ldots, \theta_{N}\right)\right)$. For any point $z \in J\left(f_{\lambda}\right)$ whose orbit avoids the graph, the puzzle piece of depth $d$ containing $z$ is denoted by $P_{d}^{\lambda}(z)$. We say the graph $G_{\lambda}\left(\theta_{1}, \ldots, \theta_{N}\right)$ is admissible if there exists a non-degenerate critical annulus $P_{d}^{\lambda}(c) \backslash \overline{P_{d+1}^{\lambda}(c)}$ (or $P_{0}^{\lambda}\left(f_{\lambda}^{d}(c)\right) \backslash \overline{P_{1}^{\lambda}\left(f_{\lambda}^{d}(c)\right)}$ ) for some $c \in C_{\lambda}$ and some $d \geq 1$.

Lemma 5.2 ([28], Proposition 4.1). - Suppose $\lambda \in \mathcal{M} \cap \mathscr{F}$ and the map $f_{\lambda}$ is postcritically infinite, then there exists an admissible graph $G_{\lambda}\left(\theta_{1}, \ldots, \theta_{N}\right)$.

For $c \in C_{\lambda}$, the tableau $T(c)$ is defined as the two-dimensional array $\left(P_{d}^{\lambda}\left(f_{\lambda}^{l}(c)\right)\right)_{d, l \geq 0}$. We say $T(c)$ is periodic if there is an integer $p \geq 1$ such that $P_{d}^{\lambda}\left(f_{\lambda}^{p}(c)\right)=P_{d}^{\lambda}(c)$ for all $d \geq 0$.

Lemma 5.3 ([28], Lemma 5.2 and Propositions 7.2 and 7.3). - Suppose $\lambda \in \mathcal{M} \cap \mathscr{F}$ and the graph $G_{\lambda}\left(\theta_{1}, \ldots, \theta_{N}\right)$ is admissible.

1. If $T(c)$ is periodic for some $c \in C_{\lambda}$, then $f_{\lambda}$ is either renormalizable or $*$-renormalizable. Let $K$ be the small filled Julia set of this (*-)renormalization, then $K \cap \partial B_{\lambda}$ contains at most one point.
2. If none of $T(c)$ with $c \in C_{\lambda}$ is periodic, then for any sequence of shrinking puzzle pieces $P_{0}^{\lambda} \supset P_{1}^{\lambda} \supset P_{2}^{\lambda} \supset \cdots$, the intersection $\bigcap_{d \geq 0} \overline{P_{d}^{\lambda}}$ is a singleton.

See Section 8 for the definitions of renormalizable or $*$-renormalizable.
To prove Proposition 5.1, we need the following result:
Theorem 5.4 (Lebesgue measure). - If $f_{\lambda}^{k}\left(v_{\lambda}^{+}\right) \in \partial B_{\lambda}$ for some $k \geq 0$, then the Lebesgue measure of $J\left(f_{\lambda}\right)$ is zero.

The proof of Theorem 5.4 is based on the Yoccoz puzzle theory following Lyubich [18]. For this, we put the proof in the appendix.

By Lemmas 5.2, 5.3 and assuming Theorem 5.4, the proof of Theorem 5.1 goes as follows: we first construct a topological conjugacy between $f_{\lambda_{1}}$ and $f_{\lambda_{2}}$ (Section 5.2), then we will show that this topological conjugacy is actually quasiconformal. To this end, we will need a shape bound result for puzzle pieces around every point on the Julia set (Section 5.4). Finally, we will apply a QC-criterion lemma by Kozlovski, Shen and van Strien [16, Lemma 12.1] (Sections 5.3) and Theorem 5.4 to get our result.

### 5.2. Constructing a topological conjugacy

We assume $\theta\left(\lambda_{1}\right)$ is an irrational number. In that case, both $f_{\lambda_{1}}$ and $f_{\lambda_{2}}$ are postcritically infinite. We will construct a topological conjugacy between $f_{\lambda_{1}}$ and $f_{\lambda_{2}}$ with the help of the Yoccoz puzzle.

First, by Lemma 5.2, we can find an admissible graph $G_{\lambda_{1}}\left(\theta_{1}, \ldots, \theta_{N}\right)$ for $f_{\lambda_{1}}$. That is to say, there exist $c \in C_{\lambda_{1}}$ and $k \geq 1$ such that the annulus $P_{0}^{\lambda_{1}}\left(f_{\lambda_{1}}^{k}(c)\right) \backslash \overline{P_{1}^{\lambda_{1}}\left(f_{\lambda_{1}}^{k+1}(c)\right)}$ is non-degenerate.

Define a graph

$$
\Gamma_{0}=\bigcup_{k \geq 0} \bigcup_{1 \leq j \leq N} \Omega_{\lambda_{1}}^{\tau^{k}\left(\theta_{j}\right)}
$$

It is known from Theorem 3.1 that $\partial B_{\lambda_{1}}$ and $\partial B_{\lambda_{2}}$ are Jordan curves. This allows us to construct a homeomorphism $\psi_{0}$ from the dynamical plane of $f_{\lambda_{1}}$ to that of $f_{\lambda_{2}}$ satisfying

- $\left.\psi_{0}\right|_{B_{\lambda_{1}}}=\left.\phi_{\lambda_{2}}^{-1} \circ \phi_{\lambda_{1}}\right|_{B_{\lambda_{1}}}$.
- $\psi_{0}\left(\Omega_{\lambda_{1}}^{\alpha}\right)=\Omega_{\lambda_{2}}^{\alpha}, \forall \Omega_{\lambda_{1}}^{\alpha} \subset \Gamma_{0}$.


Figure 5. Partition and labeling.

We then construct a sequence of homeomorphisms $\psi_{j}$ in the way that
(a) $f_{\lambda_{2}} \circ \psi_{j+1}=\psi_{j} \circ f_{\lambda_{1}}$ for all $j \geq 0$,
(b) $\psi_{j}\left(\Omega_{\lambda_{1}}^{\alpha}\right)=\Omega_{\lambda_{2}}^{\alpha}$ for any cut ray $\Omega_{\lambda_{1}}^{\alpha} \subset f_{\lambda_{1}}^{-j}\left(\Gamma_{0}\right)$.

The construction is inductively as follows. For $\lambda \in\left\{\lambda_{1}, \lambda_{2}\right\}$, any $d \geq 0$ and any $z \in C_{\lambda} \cup\left\{v_{\lambda}^{+}, v_{\lambda}^{-}\right\}$, let $Q_{d}^{\lambda}(z)$ be the component of $\overline{\mathbb{C}} \backslash f_{\lambda}^{-d}\left(\Gamma_{0}\right)$ containing $z$. The domain $\Gamma_{d}^{\lambda}:=\overline{\mathbb{C}} \backslash\left(\overline{Q_{d}^{\lambda}\left(v_{\lambda}^{+}\right) \cup Q_{d}^{\lambda}\left(v_{\lambda}^{-}\right)}\right)$either is empty or consists of one or two topological disks. Each component of $f_{\lambda}^{-1}\left(\overline{\mathbb{C}} \backslash \Gamma_{d}^{\lambda}\right)$ is a disk. Let $Q_{d+1, j}^{\lambda}$ be its component lying in between $Q_{d+1}^{\lambda}\left(c_{j}(\lambda)\right)$ and $Q_{d+1}^{\lambda}\left(c_{j+1}(\lambda)\right)$ for $0 \leq j<2 n, c_{2 n}(\lambda)=c_{0}(\lambda)$. See Figure 5. Note that the map $\left.f_{\lambda}\right|_{Q_{d+1, j}^{\lambda}}: Q_{d+1, j}^{\lambda} \rightarrow \Gamma_{d}^{\lambda}$ is a conformal isomorphism.

Suppose that $\psi_{0}, \psi_{1}, \ldots, \psi_{d}$ are already defined and satisfy (a), (b). The assumption $\theta\left(\lambda_{1}\right)=\theta\left(\lambda_{2}\right)$ implies that the piece $Q_{d}^{\lambda_{1}}\left(v_{\lambda_{1}}^{+}\right)$is bounded by $\Omega_{\lambda_{1}}^{\alpha_{1}}, \ldots, \Omega_{\lambda_{1}}^{\alpha_{s}} \subset f_{\lambda_{1}}^{-d}\left(\Gamma_{0}\right)$ if and only if $Q_{d}^{\lambda_{1}}\left(v_{\lambda_{1}}^{+}\right)$is bounded by $\Omega_{\lambda_{2}}^{\alpha_{1}}=\psi_{d}\left(\Omega_{\lambda_{1}}^{\alpha_{1}}\right), \ldots, \Omega_{\lambda_{2}}^{\alpha_{s}}=\psi_{d}\left(\Omega_{\lambda_{1}}^{\alpha_{s}}\right) \subset f_{\lambda_{2}}^{-d}\left(\Gamma_{0}\right)$. This fact is very important. It enables us to get a lift $\psi_{d+1}$ of $\psi_{d}$. Actually, we can define $\psi_{d+1}$ piece by piece. Set $\left.\psi_{d+1}\right|_{Q_{d+1, j}^{\lambda_{1}}}=\left(\left.f_{\lambda_{2}}\right|_{Q_{d+1, j}^{\lambda_{2}}}\right)^{-1} \circ \psi_{d} \circ\left(\left.f_{\lambda_{1}}\right|_{Q_{d+1, j}^{\lambda_{1}}}\right)$. We then define $\left.\psi_{d+1}\right|_{Q_{d+1}^{\lambda_{1}}\left(c_{j}\right)}$ so that it coincides with $\left.\psi_{d+1}\right|_{Q_{d+1, j}^{\lambda_{1}}}$ in their common boundary and that the following diagram commutes

$$
\begin{aligned}
& \overline{Q_{d+1}^{\lambda_{1}}\left(c_{j}\left(\lambda_{1}\right)\right)} \xrightarrow{f_{\lambda_{1}}} \overline{Q_{d}^{\lambda_{1}}\left(f_{\lambda_{1}}\left(c_{j}\left(\lambda_{1}\right)\right)\right)} \\
& \overline{\psi_{d+1}} \downarrow \\
& \overline{Q_{d+1}^{\lambda_{2}}\left(c_{j}\left(\lambda_{2}\right)\right)} \xrightarrow[f_{\lambda_{2}}]{ } \overline{Q_{d}^{\lambda_{2}}\left(f_{\lambda_{2}}\left(c_{j}\left(\lambda_{2}\right)\right)\right) .}
\end{aligned}
$$

One may verify that $\psi_{d+1}$ is well defined and satisfies $f_{\lambda_{2}} \circ \psi_{d+1}=\psi_{d} \circ f_{\lambda_{1}}$. By induction assumption, $\psi_{d}$ preserves the $d$-th preimages of $\Gamma_{0}$. Then the condition $\theta\left(\lambda_{1}\right)=\theta\left(\lambda_{2}\right)$ and the construction of $\psi_{d+1}$ implies that $\psi_{d+1}\left(\Omega_{\lambda_{1}}^{\alpha}\right)=\Omega_{\lambda_{2}}^{\alpha}$ for any cut ray $\Omega_{\lambda_{1}}^{\alpha} \subset f_{\lambda_{1}}^{-d-1}\left(\Gamma_{0}\right)$.

In this way, we get a sequence of homeomorphism $\psi_{j}, j \geq 0$. The construction implies that

1 . For any $d \geq 0$,

$$
\left.\psi_{d+1}\right|_{f_{\lambda_{1}}^{-d}\left(B_{\lambda_{1}}\right)}=\left.\psi_{d}\right|_{f_{\lambda_{1}}^{-d}\left(B_{\lambda_{1}}\right)} .
$$

2. The graph $G_{\lambda_{2}}\left(\theta_{1}, \ldots, \theta_{N}\right)$ is also admissible for $f_{\lambda_{2}}$ and the annulus $P_{0}^{\lambda_{2}}\left(f_{\lambda_{2}}^{k}(c)\right) \backslash \overline{P_{1}^{\lambda_{2}}\left(f_{\lambda_{2}}^{k+1}(c)\right)}$ is non-degenerate.

By Lemma 5.3, the $\psi_{j}$ converges to a continuous and one-to-one map $\psi_{\infty}: \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$, holomorphic in the Fatou set $F\left(f_{\lambda_{1}}\right)=\cup_{k} f_{\lambda_{1}}^{-k}\left(B_{\lambda_{1}}\right)$ and satisfies $f_{\lambda_{2}} \circ \psi_{\infty}=\psi_{\infty} \circ f_{\lambda_{1}}$. This is a topological conjugacy, as required.

### 5.3. Rigidity

In order to improve the quality of $\psi_{\infty}$ constructed in the previous section, we first introduce the QC-Criterion result in [16].

For a homeomorphism $\phi: \Omega \rightarrow \tilde{\Omega}$ and for $x \in \Omega$, let

$$
\underline{H}(\phi, x)=\liminf _{r \rightarrow 0} \frac{\sup _{|y-x|=r}|\phi(y)-\phi(x)|}{\inf _{|y-x|=r}|\phi(y)-\phi(x)|} \in[1, \infty] .
$$

Lemma 5.5. - Let $\phi: \Omega \rightarrow \tilde{\Omega}$ be a homeomorphism between two Jordan domains, $k \in(0,1)$ be a constant. Let $X$ be a subset of $\Omega$ such that both $X$ and $\phi(X)$ have zero Lebesgue measures. Assume the following hold:

1. $|\bar{\partial} \phi| \leq k|\partial \phi|$ a.e on $\Omega \backslash X$.
2. For each $x \in X, \underline{H}(\phi, x)<\infty$.

Then $\phi$ is a $K$-quasiconformal map, where $K$ depends only on $k$.

Lemma 5.5 is a simplified version of [16, Lemma 12.1]. See therein for a detailed proof based on a standard extremal length argument.

Let $U \subsetneq \mathbb{C}$ be a simply connected planar domain and $z \in U$. The shape of $U$ about $z$ is defined by:

$$
\operatorname{Shape}(U, z)=\sup _{x \in \partial U}|x-z| / \inf _{x \in \partial U}|x-z| .
$$

Remark 5.6. - In Lemma 5.5, Condition 2 can be replaced by: There is a constant $M>0$ such that for each $x \in X$, there exist a number $M_{x}>0$ and a sequence of shrinking topological disks $D_{1} \supset D_{2} \supset \cdots$ containing $x$, such that
(a) $\cap \overline{D_{j}}=\{x\}$;
(b) $\operatorname{Shape}\left(D_{j}, x\right) \leq M, \operatorname{Shape}\left(\phi\left(D_{j}\right), \phi(x)\right) \leq M_{x}, \forall j \geq 1$.

The original proof of [16, Lemma 12.1] goes through without essential differences.
By Lemma 5.5 and Remark 5.6 , to show that $\psi_{\infty}$ is a quasiconformal map it suffices to show that

Lemma 5.7. - There is a constant $M>0$ such that for any $z \in J\left(f_{\lambda_{1}}\right)$, there exist a number $M_{z}>0$ and a sequence of topological disks $D_{1} \supset D_{2} \supset \cdots$ containing $z$, such that
(a) $\cap \overline{D_{j}}=\{z\}$;
(b) $\operatorname{Shape}\left(D_{j}, z\right) \leq M$, $\operatorname{Shape}\left(\psi_{\infty}\left(D_{j}\right), \psi_{\infty}(z)\right) \leq M_{z}, \forall j \geq 1$.

Assuming Lemma 5.7, we can give the following
Proof of Theorem 5.1 when $\theta\left(\lambda_{1}\right)$ is an irrational number. - In this case, with the help of the Yoccoz puzzle, we have constructed a topological conjugacy $\psi_{\infty}$, holomorphic in the Fatou set of $f_{\lambda_{1}}$. By Lemmas 5.5 and $5.7, \psi_{\infty}$ is a quasiconformal map.

By Theorem 5.4, the Lebesgue measure of $J\left(f_{\lambda_{1}}\right)$ is zero, so $\psi_{\infty}$ is a Möbius map of the form $\psi_{\infty}(z)=a z$. One may verify that $a^{n-1}=1$ and $\lambda_{2}=a^{2} \lambda_{1}$. The condition $\lambda_{1}, \lambda_{2} \in \mathcal{F}$ implies that $\lambda_{1}=\lambda_{2}$.

### 5.4. Proof of Lemma 5.7

Recall that for $\lambda \in\left\{\lambda_{1}, \lambda_{2}\right\}$, we have found an admissible graph $G_{\lambda}\left(\theta_{1}, \ldots, \theta_{N}\right)$ and had a Yoccoz puzzle construction induced by this graph.

We first define the Yoccoz $\tau_{Y}$-function as follows. We choose some $c \in C_{\lambda}$. For each $d \geq 1$, we define $\tau_{Y}(d)$ to be the biggest integer $k \in[0, d-1]$ such that the puzzle piece $f_{\lambda}^{d-k}\left(P_{d}^{\lambda}(c)\right)$ contains some critical point in $C_{\lambda}$, we set $\tau_{Y}(d)=-1$ if no such integer exists. Since $f_{\lambda}\left(e^{\pi i / n} z\right)=-f_{\lambda}(z)$, by the symmetry of puzzle pieces (namely, $P_{d}^{\lambda}\left(e^{\pi i / n} z\right)=e^{\pi i / n} P_{d}^{\lambda}(z)$ for all $d \geq 1$, see [28, Lemma 4.1]), we see that the Yoccoz $\tau_{Y}$-function is well-defined (independent of the choice of $c \in C_{\lambda}$ ). Moreover, it satisfies $\tau_{Y}(d+1) \leq \tau_{Y}(d)+1$.

We say that the critical set $C_{\lambda}$ is non-recurrent if $\tau_{Y}(d)$ is uniformly bounded for all $d \geq 1$; recurrent if $\lim \sup \tau_{Y}(d)=\infty$ (this definition is in fact consistent with the definition in Section 3). There are two cases when $C_{\lambda}$ is recurrent. We say $C_{\lambda}$ is reluctantly recurrent if $\lim \inf \tau_{Y}(d)<\infty ;$ persistently recurrent if $\lim \inf \tau_{Y}(d)=\infty$.

For $\lambda \in\left\{\lambda_{1}, \lambda_{2}\right\}$ and $z \in J\left(f_{\lambda}\right)$, let $\omega(z)$ be the $\omega$-limit set of $z$, defined as

$$
\left\{w \in J\left(f_{\lambda}\right) ; \text { there exist } n_{k} \rightarrow \infty \text { such that } f_{\lambda}^{n_{k}}(z) \rightarrow w\right\} .
$$

Recall that $\Gamma_{0}^{\lambda}$ is defined as

$$
\Gamma_{0}^{\lambda}=\bigcup_{k \geq 0} \bigcup_{1 \leq j \leq N} \Omega_{\lambda}^{\tau^{k}\left(\theta_{j}\right)}
$$

We then decompose $J\left(f_{\lambda}\right)$ into three disjoint sets $J_{0}^{\lambda}, J_{1}^{\lambda}, J_{2}^{\lambda}$, where

$$
\begin{aligned}
& J_{0}^{\lambda}=J\left(f_{\lambda}\right) \cap\left(\bigcup_{k \geq 0} f_{\lambda}^{-k}\left(\Gamma_{0}^{\lambda}\right)\right), \\
& J_{1}^{\lambda}=\left\{z \in J\left(f_{\lambda}\right) \backslash J_{0}^{\lambda} ; C_{\lambda} \cap \omega(z)=\varnothing\right\}, \\
& J_{2}^{\lambda}=\left\{z \in J\left(f_{\lambda}\right) \backslash J_{0}^{\lambda} ; C_{\lambda} \cap \omega(z) \neq \varnothing\right\} .
\end{aligned}
$$

Our task is to verify Lemma 5.7 in these three cases.
Case 1: $z \in J_{0}^{\lambda}$. - Let us first look at the points $z \in J\left(f_{\lambda}\right) \cap \Gamma_{0}^{\lambda}$. The idea of the proof is based on the expanding property of $f_{\lambda}$ on $J\left(f_{\lambda}\right) \cap \Gamma_{0}^{\lambda}$. Recall that $S_{k}^{\lambda}$ with $k= \pm 1, \ldots, \pm(n-1)$ are the closed sectors defined in Section 3, let $\widehat{B}_{\lambda}=\left\{w \in B_{\lambda} ;\left|\phi_{\lambda}(w)\right|>2\right\}$ be a subset of $B_{\lambda}$ and $\widehat{S}_{k}^{\lambda}=S_{k}^{\lambda} \backslash f_{\lambda}^{-1}\left(\widehat{B}_{\lambda}\right)$. For each $z \in J\left(f_{\lambda}\right) \cap \Gamma_{0}^{\lambda}$ and for all $k \geq 0$, let $U_{0}\left(f_{\lambda}^{k}(z)\right)$ be the $\widehat{S}_{j}^{\lambda}$ containing $f_{\lambda}^{k}(z)$ (note that $U_{0}\left(f_{\lambda}^{k}(z)\right)$ is always well-defined because $f_{\lambda}^{k}(z)$ cannot sit on the boundary of two sectors), and then let $U_{k}(z)$ be the component of $f_{\lambda}^{-1}\left(U_{0}\left(f_{\lambda}^{k}(z)\right)\right)$ containing $z$. The following facts are immediate:

1. For all $z \in J\left(f_{\lambda}\right) \cap \Gamma_{0}^{\lambda}$, the compact set $U_{k+1}(z)$ is contained in the interior of $U_{k}(z)$.
2. There are two constants $C_{1}, C_{2}>0$ such that for all $z \in J\left(f_{\lambda}\right) \cap \Gamma_{0}^{\lambda}$, we have Shape $\left(U_{1}(z), z\right) \leq C_{1}$ and $\bmod \left(U_{0}(z) \backslash U_{1}(z)\right) \geq C_{2}$.

Note that for all $k \geq 0$, the map $f_{\lambda}^{k-1}: U_{k}(z) \rightarrow U_{1}\left(f_{\lambda}^{k}(z)\right)$ is a conformal map. By the shape distortion [28, Lemma 6.1], we have

$$
\operatorname{Shape}\left(U_{k}(z), z\right) \leq C_{3} \operatorname{Shape}\left(U_{1}\left(f_{\lambda}^{k}(z)\right), f_{\lambda}^{k}(z)\right) \leq C_{3} C_{1},
$$

where $C_{3}$ depends on $C_{2}$.
By pulling back via $f_{\lambda}$ and the shape distortion, one sees that the above inequality holds for all $z \in J_{0}^{\lambda}$.

In particular, the above discussion holds for $\lambda=\lambda_{1}$. In order to verify Lemma 5.7, it remains to show that $\operatorname{Shape}\left(\psi_{\infty}\left(U_{k}(z)\right), \psi_{\infty}(z)\right)$ is uniformly bounded above for all $k$ and all $z \in J_{0}^{\lambda_{1}}$. By the uniform continuity of $\psi_{\infty}$, there are two constants $C_{1}^{\prime}, C_{2}^{\prime}>0$ such that for any $z \in J\left(f_{\lambda_{1}}\right) \cap \Gamma_{0}^{\lambda_{1}}$, Shape $\left(\psi_{\infty}\left(U_{1}(z)\right), \psi_{\infty}(z)\right) \leq C_{1}^{\prime}$ and $\bmod \left(\psi_{\infty}\left(U_{0}(z) \backslash U_{1}(z)\right)\right) \geq C_{2}^{\prime}$. Because $f_{\lambda_{2}}^{k-1}: \psi_{\infty}\left(U_{k}(z)\right) \rightarrow \psi_{\infty}\left(U_{0}\left(f_{\lambda_{1}}^{k}(z)\right)\right)$ is a conformal map, by the shape distortion [28, Lemma 6.1], we have

$$
\operatorname{Shape}\left(\psi_{\infty}\left(U_{k}(z)\right), \psi_{\infty}(z)\right) \leq C_{3}^{\prime} \operatorname{Shape}\left(\psi_{\infty}\left(U_{1}\left(f_{\lambda_{1}}^{k}(z)\right)\right), \psi_{\infty}\left(f_{\lambda_{1}}^{k}(z)\right)\right) \leq C_{3}^{\prime} C_{1}^{\prime}
$$

where $C_{3}^{\prime}$ depends on $C_{2}^{\prime}$. By pulling back via $f_{\lambda_{2}}$ and the shape distortion, we see that the above inequality holds for all $z \in J_{0}^{\lambda_{1}}$.

Case 2: $z \in J_{1}^{\lambda_{1}}$. - In this case, by Lemma 5.3, $C_{\lambda_{1}} \cap \omega(z)=\varnothing$ implies that there exists an integer $n_{0}=n_{0}(z)$ such that $f_{\lambda_{1}}^{k}(z) \notin \bigcup_{c \in C_{\lambda_{1}}} P_{n_{0}}^{\lambda_{1}}(c)$ for all $k \geq 1$. There also exists a sequence of integers $k_{j} \rightarrow \infty$ such that $f_{\lambda_{1}}^{k_{j}}(z) \rightarrow z^{*} \in J\left(f_{\lambda_{1}}\right) \backslash J_{0}^{\lambda_{1}}$. By passing to a subsequence, we assume $f_{\lambda_{1}}^{k_{j}}(z) \in P_{n_{0}}^{\lambda_{1}}\left(z^{*}\right)$ for all $j$. We have $\operatorname{deg}\left(f_{\lambda_{1}}^{k_{j}}: P_{n_{0}+k_{j}}^{\lambda_{1}}(z) \rightarrow P_{n_{0}}^{\lambda_{1}}\left(z^{*}\right)\right) \leq 2$. Take a small number $r>0$ such that $\bmod \left(P_{n_{0}}^{\lambda_{1}}\left(z^{*}\right) \backslash \overline{D\left(z^{*}, r\right)}\right) \geq 1$, where $D\left(z^{*}, r\right)$ is the open Euclidean disk centered at $z^{*}$ with radius $r$. For large $j$, we have $f_{\lambda_{1}}^{k_{j}}(z) \in D\left(z^{*}, r / 2\right)$. Let $D_{j}(z)$ be the component of $f_{\lambda_{1}}^{-k_{j}}\left(D\left(z^{*}, r\right)\right)$ containing $z$. Then by the shape distortion [28, Lemma 6.1], there is a universal constant $C$ such that for all large $j$,

$$
\operatorname{Shape}\left(D_{j}(z), z\right) \leq C \operatorname{Shape}\left(D\left(z^{*}, r\right), f_{\lambda_{1}}^{k_{j}}(z)\right) \leq 3 C .
$$

Let $m=\bmod \left(\psi_{\infty}\left(P_{n_{0}}^{\lambda_{1}}\left(z^{*}\right) \backslash \overline{D\left(z^{*}, r\right)}\right)\right)=\bmod \left(P_{n_{0}}^{\lambda_{2}}\left(\psi_{\infty}\left(z^{*}\right)\right) \backslash \overline{\psi_{\infty}\left(D\left(z^{*}, r\right)\right)}\right)$. Again by the shape distortion [28, Lemma 6.1], there is a constant $C_{1}$ depending on $m$ such that for all large $j$,

$$
\operatorname{Shape}\left(\psi_{\infty}\left(D_{j}(z)\right), \psi_{\infty}(z)\right) \leq C_{1} \operatorname{Shape}\left(\psi_{\infty}\left(D\left(z^{*}, r\right)\right), \psi_{\infty}\left(f_{\lambda_{1}}^{k_{j}}(z)\right)\right) \leq C_{1} C_{2},
$$

where $C_{2}$ depends on $z$ but is independent of $j$.
Case 3: $z \in J_{2}^{\lambda}$ and $c \in C_{\lambda} \cap \omega(z), \lambda \in\left\{\lambda_{1}, \lambda_{2}\right\}$. - There are three subcases:
Case 3.1: $C_{\lambda}$ is non-recurrent. - First there are three puzzle pieces $P_{n_{2}}^{\lambda} \Subset P_{n_{1}}^{\lambda} \Subset P_{n_{0}}^{\lambda}$ of depths $n_{2}, n_{1}, n_{0}$ and a sequence of integers $\left\{k_{j}\right\}_{j}$ such that $f_{\lambda}^{k_{j}}(c) \in P_{n_{2}}^{\lambda}$ and $\operatorname{deg}\left(f_{\lambda}^{k_{j}}: P_{n_{0}+k_{j}}^{\lambda}(c) \rightarrow P_{n_{0}}^{\lambda}\right)=2$ for all $j$. For each $j$, let $l_{j} \geq 0$ be the smallest integer such that $f_{\lambda}^{l_{j}}(z) \in P_{n_{2}+k_{j}}^{\lambda}(c)$. Then

$$
\operatorname{deg}\left(f_{\lambda}^{l_{j}}: P_{n_{0}+k_{j}+l_{j}}^{\lambda}(z) \rightarrow P_{n_{0}+k_{j}}^{\lambda}(c)\right) \leq C 2^{n_{2}-n_{0}}
$$

for all $j$, where $C=2^{\operatorname{Card}\left(C_{\lambda}\right)}=2^{2 n}$. We have that for all $j$,

$$
\operatorname{deg}\left(f_{\lambda}^{k_{j}+l_{j}}: P_{n_{0}+k_{j}+l_{j}}^{\lambda}(z) \rightarrow P_{n_{0}}^{\lambda}\right) \leq C 2^{n_{2}-n_{0}+1}
$$

and $f_{\lambda}^{k_{j}+l_{j}}(z) \in P_{n_{2}}^{\lambda}$. By the shape distortion [28, Lemma 6.1],

$$
\operatorname{Shape}\left(P_{n_{1}+k_{j}+l_{j}}^{\lambda}(z), z\right) \leq M_{2}<+\infty .
$$

This constant $M_{2}$ is independent of the points $z \in J_{2}^{\lambda}$ and $j$.
Case 3.2: $C_{\lambda}$ is reluctantly recurrent. - In this case, there exist $c^{\prime} \in C_{\lambda}\left(c^{\prime}\right.$ may be $\left.c\right)$ and integers $n_{0}$ and $\left\{k_{j}\right\}_{j}$ such that $f_{\lambda}^{k_{j}}$ maps $P_{n_{0}+k_{j}}^{\lambda}(c)$ to $P_{n_{0}}^{\lambda}\left(c^{\prime}\right)$ of degree two. Take puzzle pieces $P_{n_{2}}^{\lambda}\left(c^{\prime}\right) \Subset P_{n_{1}}^{\lambda}\left(c^{\prime}\right) \Subset P_{n_{0}}^{\lambda}\left(c^{\prime}\right)$. For each $j$, let $l_{j} \geq 0$ be the smallest integer such that $f_{\lambda}^{l_{j}}\left(f_{\lambda}^{k_{j}}(c)\right) \in P_{n_{2}}^{\lambda}\left(c^{\prime}\right)$, then for all $j$,

$$
\operatorname{deg}\left(f_{\lambda}^{k_{j}+l_{j}}: P_{n_{0}+k_{j}+l_{j}}^{\lambda}(c) \rightarrow P_{n_{0}}^{\lambda}\left(c^{\prime}\right)\right) \leq 2^{n_{2}-n_{0}}
$$

For each $j$, let $m_{j}$ be the smallest integer such that $f_{\lambda}^{m_{j}}(z) \in P_{n_{2}+k_{j}+l_{j}}^{\lambda}(c)$. Then

$$
\operatorname{deg}\left(f_{\lambda}^{m_{j}}: P_{n_{0}+k_{j}+l_{j}+m_{j}}^{\lambda}(z) \rightarrow P_{n_{0}+k_{j}+l_{j}}^{\lambda}(c)\right) \leq C 2^{n_{2}-n_{0}},
$$

where $C=2^{\operatorname{Card}\left(C_{\lambda}\right)}=2^{2 n}$. We have that for all $j$,

$$
\operatorname{deg}\left(f_{\lambda}^{k_{j}+l_{j}+m_{j}}: P_{n_{0}+k_{j}+l_{j}+m_{j}}^{\lambda}(z) \rightarrow P_{n_{0}}^{\lambda}\left(c^{\prime}\right)\right) \leq C 4^{n_{2}-n_{0}} .
$$

Therefore, by the shape distortion [28, Lemma 6.1],

$$
\operatorname{Shape}\left(P_{n_{1}+k_{j}+l_{j}+m_{j}}^{\lambda}(z), z\right) \leq M_{3}<+\infty .
$$

This constant $M_{3}$ is independent of the points $z \in J_{2}^{\lambda}$ and $j$.
Case 3.3: $C_{\lambda}$ is persistently recurrent. - In this case, the following crucial property holds:
There is a number $m>0$ and a sequence of puzzle pieces containing $c$ :

$$
K_{1}^{\prime}(c) \supset K_{1}(c) \supset \widetilde{K}_{1}(c) \supset K_{2}^{\prime}(c) \supset K_{2}(c) \supset \widetilde{K}_{2}(c) \supset \cdots
$$

such that $\bigcap \overline{K_{j}(c)}=\{c\}$ and for all $j \geq 1$,

$$
\bmod \left(K_{j}^{\prime}(c) \backslash \overline{K_{j}(c)}\right), \bmod \left(K_{j}(c) \backslash \overline{\widetilde{K}_{j}(c)}\right) \geq m
$$

and

$$
\left(K_{j}^{\prime}(c) \backslash \overline{K_{j}(c)}\right) \cap\left(\bigcup_{k \geq 0} f_{\lambda}^{k}\left(C_{\lambda}\right)\right)=\left(K_{j}(c) \backslash \overline{\widetilde{K}_{j}(c)}\right) \cap\left(\bigcup_{k \geq 0} f_{\lambda}^{k}\left(C_{\lambda}\right)\right)=\varnothing .
$$

This construction of the sequence of puzzle pieces is due to Kozlovski, Shen and van Strien in [16, Section 8]. The proof of the complex bounds of the moduli is given by Kozlovskivan Strien [17] and Qiu-Yin [29] independently. In literature, this sequence of puzzle pieces is called KSS nest or enhanced nest. See [16] for its original construction.

In our case, there is essentially only one free critical orbit, the construction is simplified. Following [35], we will sketch the construction of a sequence of puzzle pieces containing $c$ inductively:

$$
I_{1}(c) \supsetneqq K_{1}^{\prime}(c) \supsetneqq K_{1}(c) \supsetneqq \cdots I_{j}(c) \supsetneqq K_{j}^{\prime}(c) \supsetneqq K_{j}(c) \supsetneqq I_{j+1}(c) \supsetneqq \cdots .
$$

Let $I_{1}(c)=P_{L}^{\lambda}(c)$ be a critical puzzle piece, the number $L$ is chosen so that the annulus $P_{L}^{\lambda}(c) \backslash \overline{P_{L+1}^{\lambda}(c)}$ is non-degenerate. Assume that $I_{j}(c)=P_{k_{j}}^{\lambda}(c)$ is defined. It follows from [28, Lemma 4.7] that the set

$$
S_{j}=\left\{k \geq L: \tau_{Y}(k)=k_{j}, \tau_{Y}(k+1)=k_{j}+1\right\}
$$

consists of at least two elements. Since the critical set $C_{\lambda}$ is persistently recurrent, the set $S_{j}$ consists of at most finite elements. Let $a_{j}=\min S_{j}$ and $b_{j}=\max S_{j}$. Set $K_{j}^{\prime}(c)=P_{b_{j}}^{\lambda}(c)$ and $K_{j}(c)=P_{t_{j}}^{\lambda}(c)$ with $t_{j}=b_{j}+a_{j}-k_{j}$. Let $s_{j}>b_{j}-k_{j}$ be the largest integer such that $P_{a_{j}}^{\lambda}\left(f_{\lambda}^{k}(c)\right) \cap C_{\lambda}=\varnothing$ for all $b_{j}-k_{j}<k<s_{j}$, and set $\widetilde{K}_{j}(c)=P_{s_{j}}^{\lambda}(c)$. Now let

$$
Z_{j}=\left\{k \geq L: \tau_{Y}(k)=b_{j}+a_{j}-k_{j}, \tau_{Y}(k+1)=b_{j}+a_{j}-k_{j}+1\right\} .
$$

Again, the set $Z_{j}$ is a finite set containing at least two elements. Let $k_{j+1}=\max Z_{j}$ and $I_{j+1}(c)=P_{k_{j+1}}^{\lambda}(c)$. The sequence of puzzle pieces is then constructed inductively.

In this case, by [36, Proposition 1], there is a constant $M>0$ such that the $\operatorname{Shape}\left(K_{j}(c), c\right) \leq M$ for all $j$. For each $j$, let $m_{j}$ be the smallest integer such that $f_{\lambda}^{m_{j}}(z) \in \widetilde{K}_{j}(c)=P_{s_{j}}^{\lambda}(c)$. It follows that

$$
\operatorname{deg}\left(f_{\lambda}^{m_{j}}: P_{m_{j}+s_{j}}^{\lambda}(z) \rightarrow P_{s_{j}}^{\lambda}(c)\right) \leq 2^{\operatorname{Card}\left(C_{\lambda}\right)} .
$$

The property

$$
\left(K_{j}^{\prime}(c) \backslash \overline{K_{j}(c)}\right) \cap\left(\bigcup_{k \geq 0} f_{\lambda}^{k}\left(C_{\lambda}\right)\right)=\left(K_{j}(c) \backslash \overline{\widetilde{K}_{j}(c)}\right) \cap\left(\bigcup_{k \geq 0} f_{\lambda}^{k}\left(C_{\lambda}\right)\right)=\varnothing
$$

implies that

$$
\operatorname{deg}\left(f_{\lambda}^{m_{j}}: P_{m_{j}+b_{j}}^{\lambda}(z) \rightarrow P_{b_{j}}^{\lambda}(c)\right)=\operatorname{deg}\left(f_{\lambda}^{m_{j}}: P_{m_{j}+t_{j}}^{\lambda}(z) \rightarrow P_{t_{j}}^{\lambda}(c)\right) \leq 2^{\operatorname{Card}\left(C_{\lambda}\right)} .
$$

By the complex bounds for moduli and the shape distortion, we have

$$
\operatorname{Shape}\left(P_{t_{j}+m_{j}}^{\lambda}(z), z\right) \leq M_{4}<+\infty .
$$

This constant $M_{4}$ is independent of the points $z \in J_{2}^{\lambda}$ and $j$.
This completes the proof of Lemma 5.7 in all possible cases.

## 6. $\partial \mathscr{H}_{0}$ is a Jordan curve

In this section, we prove that $\partial \mathscr{H}_{0}$ is a Jordan curve and give some consequences.
Theorem 6.1. - $\partial \mathscr{H}_{0}$ is a Jordan curve.
Proof. - We first show that $\chi_{0}$ is a singleton. To do this, first note that the parameter ray $\mathscr{R}_{0}(0)$ is contained in the real and positive axis. So $\chi_{0}$ contains at least one positive number. We define $g_{\lambda}(z)=z^{n}\left(f_{\lambda}(z)-z\right)=z^{2 n}-z^{n+1}+\lambda$ for $\lambda, z>0$. The positive critical point of $g_{\lambda}$ is $z_{*}=\left(\frac{n+1}{2 n}\right)^{\frac{1}{n-1}}$ and for all $z>z_{*}$, we have $g_{\lambda}^{\prime}(z)>0$.

Let $\lambda_{*}$ solve $g_{\lambda_{*}}\left(z_{*}\right)=0$, then $\lambda_{*}=\frac{n-1}{2 n}\left(\frac{n+1}{2 n}\right)^{\frac{n+1}{n-1}}$. For any $\lambda>\lambda_{*}$, we have $g_{\lambda}>0$. In this case, let us look at the graph of the real function $f_{\lambda}(z)$ with $\lambda, z>0$, we see that for any $\lambda>\lambda_{*}, z>0$, we have $f_{\lambda}^{k}(z) \rightarrow \infty$ as $k \rightarrow \infty$. This implies $(0,+\infty) \subset B_{\lambda}$. In particular, $v_{\lambda}^{+} \in B_{\lambda}$. Thus $\left(\lambda_{*},+\infty\right) \subset \mathcal{R}_{0}(0)$. On the other hand, we have $f_{\lambda_{*}}\left(z_{*}\right)=z_{*}$ and $f_{\lambda_{*}}^{\prime}\left(z_{*}\right)=1$. So $\lambda_{*}$ is a cusp and $\lambda_{*} \in \chi_{0}$. Moreover, by elementary properties of real functions, there is a small number $\epsilon>0$ such that for all $\lambda \in\left(\lambda_{*}-\epsilon, \lambda_{*}\right)$, the map $f_{\lambda}$ has an attracting cycle. So $\left(\lambda_{*}-\epsilon, \lambda_{*}\right)$ is contained in a hyperbolic component (of renormalizable type, see Section 8) and $\left(\lambda_{*}-\epsilon, \lambda_{*}\right) \cap \chi_{0}=\varnothing$. If $\chi_{0} \backslash\left\{\lambda_{*}\right\} \neq \varnothing$, then there is $\lambda \in \chi_{0} \cap \mathcal{F}$ which is not a cusp. By Lemma 4.5, we have $0<\theta(\lambda)<\frac{1}{2(n-1)}$. However by Lemma 4.8, we have $\theta(\lambda)=0$. This leads to a contradiction.

In the following, we assume $t \in\left(0, \frac{1}{n-1}\right)$. Take two parameters $\lambda_{1}, \lambda_{2} \in \chi_{t} \cap \mathcal{F}$ which are not cusps, it follows from Lemma 4.8 that $\theta\left(\lambda_{1}\right)=\theta\left(\lambda_{2}\right)=t / 2$. By Theorem 5.1 we have $\lambda_{1}=\lambda_{2}$. Since there are countably many cusps, the impression $\chi_{t}$ is necessarily a singleton. So $\partial \mathscr{H}_{0}$ is locally connected.

If there are two different angles $t_{1}, t_{2} \in\left[0, \frac{1}{n-1}\right)$ with $\chi_{t_{1}}=\chi_{t_{2}}=\{\lambda\}$, then by Lemma 4.8, the external rays $R_{\lambda}\left(t_{1} / 2\right)$ and $R_{\lambda}\left(t_{2} / 2\right)$ land at the same point on $\partial B_{\lambda}$. But this is a contradiction since $\partial B_{\lambda}$ is a Jordan curve (Theorem 3.1).

Theorem 6.1 has several consequences. First, one gets a canonical parameterization $\nu: \mathbb{S} \rightarrow \partial \mathscr{H}_{0}$, where $\nu(\theta)$ is defined to be the landing point of the parameter ray $\mathscr{R}_{0}(\theta)$ (namely, $\nu(\theta):=\lim _{r \rightarrow 1^{+}} \Phi_{0}^{-1}\left(r e^{2 \pi i \theta}\right)$ ).

Theorem 6.2. $-\nu(\theta)$ is a cusp if and only if $\theta$ is $\tau$-periodic.

Proof. - By Theorem 6.1, we see that $\nu(0)=\frac{n-1}{2 n}\left(\frac{n+1}{2 n}\right)^{\frac{n+1}{n-1}}$ is a cusp.
Note that $\nu\left(\theta+\frac{1}{n-1}\right)=e^{2 \pi i /(n-1)} \nu(\theta)$ and $(-1)^{n} f_{\nu\left(\theta+\frac{1}{n-1}\right)}\left(e^{\pi i /(n-1)} z\right)=e^{\pi i /(n-1)} f_{\nu(\theta)}(z)$, thus $\nu(\theta)$ is a cusp if and only if $\nu\left(\theta+\frac{1}{n-1}\right)$ is a cusp. For this, we assume $\theta \in\left(0, \frac{1}{n-1}\right)$.

If $\nu(\theta)$ is a cusp, then by Lemma 4.8, the external ray $R_{\nu(\theta)}(\theta / 2)$ lands at $\beta_{\nu(\theta)}$. By Remark 4.7, $\frac{\theta}{2}$ satisfies either $\tau^{p}\left(\frac{\theta}{2}\right) \equiv \frac{\theta}{2}$ or $\tau^{p}\left(\frac{\theta}{2}\right) \equiv \frac{\theta}{2}+\frac{1}{2}$ for some $p \geq 1$. In either case, $\theta$ is $\tau$-periodic.

Conversely, we assume $\theta$ is $\tau$-periodic. If $\nu(\theta)$ is not a cusp, then by Lemma 4.8, the external ray $R_{\nu(\theta)}\left(\frac{\theta}{2}\right)$ lands at $v_{\nu(\theta)}^{+}$. Note that $\frac{\theta}{2}$ satisfies either $\tau^{p}\left(\frac{\theta}{2}\right)=\frac{\theta}{2}$ or $\tau^{p}\left(\frac{\theta}{2}\right)=\frac{\theta}{2}+\frac{1}{2}$ for some $p \geq 1$. We have that either $f_{\nu(\theta)}^{p}\left(v_{\nu(\theta)}^{+}\right)=v_{\nu(\theta)}^{+}$or $f_{\nu(\theta)}^{p}\left(v_{\nu(\theta)}^{+}\right)=v_{\nu(\theta)}^{-}$. In the former case, we get a periodic critical point $c \in f_{\nu(\theta)}^{-1}\left(v_{\nu(\theta)}^{+}\right)$; in the latter case, we get a periodic critical point $c \in f_{\nu(\theta)}^{-1}\left(v_{\nu(\theta)}^{-}\right)$. These critical points will be in the Fatou set. But this contradicts $v_{\nu(\theta)}^{+} \in \partial B_{\nu(\theta)}$.

Remark 6.3. - As a consequence of Lemma 4.8 and Theorem 6.2,

1. if $\theta$ is $\tau$-periodic, then $\nu(\theta)$ is a cusp;
2. if $\theta$ is rational but not $\tau$-periodic, then $f_{\nu(\theta)}$ is postcritically finite;
3. if $\theta$ is irrational, then $f_{\nu(\theta)}$ is postcritically infinite.

In the last two cases, one has $C_{\nu(\theta)} \subset \partial B_{\nu(\theta)}$. Moreover, by Borel's Normal Number Theorem [4], for almost all $\theta \in(0,1]$, we have $\overline{\bigcup_{k \geq 1} f_{\nu(\theta)}^{k}\left(C_{\nu(\theta)}\right)}=\partial B_{\nu(\theta)}$.

Proposition 6.4. $-\operatorname{Set} \partial B_{0}=\mathbb{S}$ and $V=\mathbb{C} \backslash \overline{\mathcal{H}_{0}}$, then there is a holomorphic motion $H: V \times \mathbb{S} \rightarrow \mathbb{C}$ parameterized by $V$ and with base point 0 such that $H(\lambda, \mathbb{S})=\partial B_{\lambda}$ for all $\lambda \in V$.

Proof. - We first prove that every repelling periodic point of $f_{0}(z)=z^{n}$ moves holomorphically in $\mathscr{H}_{2} \cup\{0\}$. Let $z_{0} \in \mathbb{S}=J\left(f_{0}\right)$ be such a point with period $k$. For small $\lambda$, the $\operatorname{map} f_{\lambda}$ is a perturbation of $f_{0}$. By Implicit Function Theorem, there is a neighborhood $U_{0}$ of 0 such that $z_{0}$ becomes a repelling point $z_{\lambda}$ of $f_{\lambda}$ with the same period $k$, for all $\lambda \in \mathcal{U}_{0}$. On the other hand, for all $\lambda \in \mathcal{H}_{2}$, each repelling cycle of $f_{\lambda}$ moves holomorphically throughout $\mathscr{H}_{2}$ (see [24], Theorem 4.2).

Since $\mathscr{H}_{2} \cup\{0\}$ is simply connected, by the Monodromy Theorem [2], there is a holomorphic map $Z_{z_{0}}: \mathscr{H}_{2} \cup\{0\} \rightarrow \mathbb{C}$ such that $Z_{z_{0}}(\lambda)=z_{\lambda}$ for $\lambda \in \mathcal{U}_{0}$. Let $\operatorname{Per}\left(f_{0}\right)$ be all repelling periodic points of $f_{0}$. One may verify that the map $y: \mathcal{H}_{2} \cup\{0\} \times \operatorname{Per}\left(f_{0}\right) \rightarrow \mathbb{C}$ defined by $y(\lambda, z)=Z_{z}(\lambda)$ is a holomorphic motion. Note that $\mathbb{S}=\overline{\operatorname{Per}\left(f_{0}\right)}$, by $\lambda$-Lemma (see [21] or [24]), there is an extension of $y$, say $Y: \mathscr{H}_{2} \cup\{0\} \times \mathbb{S} \rightarrow \mathbb{C}$. It is obvious that $Y(\lambda, \mathbb{S})$ is a connected component of $J\left(f_{\lambda}\right)$.

Now, we show $Y(\lambda, \mathbb{S})=\partial B_{\lambda}$ for all $\lambda \in \mathscr{H}_{2} \cup\{0\}$. By the uniqueness of the holomorphic motion of hyperbolic Julia sets, it suffices to show $Y(\lambda, \mathbb{S})=\partial B_{\lambda}$ for a small and real parameter $\lambda \in(0, \epsilon)$, where $\epsilon>0$. To see this, note that when $\lambda \in(0, \epsilon)$ the fixed point $p_{0}=1$ of $f_{0}$ becomes the repelling fixed point $p_{\lambda}$ of $f_{\lambda}$, which is real and close to 1 . The map $f_{\lambda}$ has exactly two real and positive fixed points. One is $p_{\lambda}$ and the other is $p_{\lambda}^{*}$, which is near 0 . It is obvious that $p_{\lambda}$ is the landing point of the zero external ray of $f_{\lambda}$. So $Y(\lambda, 1)=p_{\lambda} \in \partial B_{\lambda}$. This implies $Y(\lambda, \mathbb{S})=\partial B_{\lambda}$ for all $\lambda \in(0, \epsilon)$.

By the above argument and Carathéodory convergence theorem, the following map defines a holomorphic motion of $\widehat{\mathbb{C}} \backslash \overline{\mathbb{D}}$ parameterized by $V$ :

$$
h: V \times(\widehat{\mathbb{C}} \backslash \overline{\mathbb{D}}) \rightarrow \widehat{\mathbb{C}} \quad h(u, z):=\phi_{u}^{-1}(z) \text { if } u \in \mathcal{U} \backslash\{0\} \text { and } h(0, z):=z .
$$

By Slodkowski's Theorem (see [14] or [33]), there is a holomorphic motion $H: V \times \widehat{\mathbb{C}}$ extending $h$ and for any $v \in \mathcal{V}$, we have $H(v, \mathbb{S})=\partial B_{v}$.

Theorem 6.5. $-\lambda \in \partial \mathscr{H}_{0}$ if and only if $\partial B_{\lambda}$ contains either the critical set $C_{\lambda}$ or a parabolic cycle of $f_{\lambda}$.

Proof. - By Lemma 4.4, it suffices to prove the 'if' part.
We first assume that $f_{\lambda}$ has a parabolic cycle on $\partial B_{\lambda}$. By Lemma 4.6, the Julia set $J\left(f_{\lambda}\right)$ contains a quasiconformal copy of the Julia set of $z \mapsto z^{2}+1 / 4$. So the boundary $\partial B_{\lambda}$ is not a quasi-circle. It follows from Proposition 6.4 that for all $v \in \mathbb{C} \backslash \overline{\mathcal{H}_{0}}, \partial B_{v}$ is a quasi-circle. Thus $\lambda \in \partial \mathscr{H}_{0}$.

Now assume $\lambda \in \mathscr{F}$ and $C_{\lambda} \subset \partial B_{\lambda}$. Recall that $\theta(\lambda)$ is defined such that $R_{\lambda}(\theta(\lambda))$ lands at $v_{\lambda}^{+}$. By Lemma 4.5, we have $0<\theta(\lambda)<\frac{1}{2(n-1)}$.

Similar to the proof of Theorem 6.2, we conclude that $2 \theta(\lambda)$ is not $\tau$-periodic. Then $\lambda^{\prime}=\nu(2 \theta(\lambda)) \in \mathscr{F}$ is not a cusp (Theorem 6.2). It satisfies $\theta\left(\lambda^{\prime}\right)=\theta(\lambda)$. It follows from Proposition 5.1 that $\lambda=\lambda^{\prime} \in \partial \mathscr{H}_{0}$.

## 7. Sierpiński holes are Jordan domains

Besides $\mathscr{H}_{0}$, there are two kinds of escape domains: the McMullen domain $\mathscr{H}_{2}$ and the Sierpiński locus $\mathscr{H}_{k}, k \geq 3$. In [6], Devaney showed that the boundary $\partial \mathscr{H}_{2}$ is a Jordan curve by constructing a sequence of analytic curves converging to it. In this section, we will show that the boundary of every Sierpiński hole is a Jordan curve. We remark that our approach also applies to $\partial \mathscr{H}_{2}$. This will yield a different proof from Devaney's. An interesting fact is that our proof relies on the boundary regularity of $\mathscr{H}_{0}$.

Let $\mathscr{H}$ be an escape domain of level $k \geq 3$. It has no intersection with $\mathbb{R}^{+}:=(0,+\infty)$. (In fact, by elementary properties of real functions, one may verify that there is a positive parameter $\lambda^{*} \in(0, \nu(0))$ such that $\left(0, \lambda^{*}\right) \subset \mathscr{H}_{2}$ and for all $\lambda \in\left[\lambda^{*}, \nu(0)\right]$, the critical orbit of $f_{\lambda}$ remains bounded, in that case, $f_{\lambda}$ is renormalizable, see [28] Lemma 7.5.)

The relation $e^{\pi i /(n-1)} f_{\lambda}(z)=(-1)^{n} f_{e^{2 \pi i /(n-1)}}\left(e^{\pi i /(n-1)} z\right)$ implies that $e^{2 \pi /(n-1)} \mathscr{H}_{k}=\mathscr{H}_{k}$. So we may assume $\mathscr{H} \subset \mathcal{F}$. The relation $\overline{f_{\lambda}(\bar{z})}=f_{\bar{\lambda}}(z)$ implies that $\mathscr{H}_{k}$ is symmetric about the real axis. We may assume further: either $\mathscr{H}$ is symmetric about $\left\{\lambda \in \mathbb{C}^{*} ; \arg \lambda=\frac{\pi}{n-1}\right\}$ or $\mathscr{H} \subset\left\{\lambda \in \mathbb{C}^{*} ; 0<\arg \lambda<\frac{\pi}{n-1}\right\}$.

The parameter ray $\mathcal{R}_{\mathscr{H}}(t)$ of angle $t \in(0,1]$ in $\mathscr{H}$ is defined by $\mathcal{R}_{\mathscr{H}}(t):=\Phi_{\mathscr{H}}^{-1}\left((0,1) e^{2 \pi i t}\right)$, its impression $\chi_{\mathscr{H}}(t)$ is defined by

$$
\left.\chi_{\mathscr{H}}(t):=\bigcap_{j \geq 1} \overline{\Phi_{\mathscr{H}}^{-1}\left(\left\{r e^{2 \pi i \theta} ; 1-1 / j<r<1,|\theta-t|<1 / j\right\}\right.}\right) .
$$

When $\lambda$ ranges over $\overline{\mathcal{H}}$, the preimages $f_{\lambda}^{2-k}(0)$ move continuously and $f_{\lambda}^{k-2}$ maps each component of $f_{\lambda}^{2-k}\left(T_{\lambda}\right)$ conformally onto $T_{\lambda}$. Let $U_{\lambda}$ be the component of $f_{\lambda}^{2-k}\left(T_{\lambda}\right)$ containing $v_{\lambda}^{+}$and $g_{\lambda}$ be the inverse of $\left.f_{\lambda}^{k-2}\right|_{U_{\lambda}}$. Both $g_{\lambda}(0)$ and $U_{\lambda}$ move continuously
for $\lambda \in \overline{\mathscr{H}}$ (and holomorphically in $\mathscr{H}$ ). The internal ray $R_{U_{\lambda}}(t)$ of angle $t$ in $U_{\lambda}$ is defined by $R_{U_{\lambda}}(t):=g_{\lambda}\left(R_{T_{\lambda}}(t)\right)$.

Lemma 7.1. - For any integer $p \geq 0$, the set $f_{\lambda}^{-p}\left(\bar{B}_{\lambda}\right)$ moves continuously (in the Hausdorff topology) with respect to $\lambda \in \mathbb{C}^{*} \backslash \overline{\mathcal{H}_{0}}$.

Proof. - It is an immediate consequence of Proposition 6.4.
Lemma 7.2. - For any $t \in[0,1)$ and any $\lambda \in \chi_{\mathscr{H}}(t) \backslash \partial \mathscr{H}_{0}$, we have $v_{\lambda}^{+} \in \partial U_{\lambda}$ and the internal ray $R_{U_{\lambda}}(t)$ lands at $v_{\lambda}^{+}$.

Proof. - It follows from Lemma 7.1 that the closure of the external ray $R_{\lambda}(t)$ moves continuously (in the Hausdorff topology) for $\lambda \in \overline{\mathscr{H}} \backslash \partial \mathscr{H}_{0}$. Note that pulling back $\overline{R_{\lambda}(t)}$ via $f_{\lambda}^{p}$ preserves the continuity.

Proposition 7.3. - For any $t \in[0,1)$, the impression $\chi_{\mathcal{t}}(t)$ is either a singleton or contained in $\partial \mathscr{H}_{0}$.

Proof. - If not, there exist $t \in[0,1)$ and a connected and compact subset $\mathcal{E}$ of $\chi_{\mathscr{H}}(t) \backslash \partial \mathscr{H}_{0}$ containing at least two points. In fact, the set $\mathscr{E}$ can be chosen as follows: we first take a point $\lambda_{0} \in \mathcal{X}_{\mathscr{H}}(t) \backslash \partial \mathscr{H}_{0}$, then let $\mathcal{E}$ be a connected component of $\left\{\lambda:\left|\lambda-\lambda_{0}\right| \leq r\right\} \cap\left(\chi_{\mathscr{H}}(t) \backslash \partial \mathscr{H}_{0}\right)$ where $r>0$ is a very small number. By Lemma 7.2, the internal ray $R_{U_{\lambda}}(t)$ lands at $v_{\lambda}^{+}$for $\lambda \in \mathcal{E}$. One may verify that for any $\lambda \in \mathcal{E}$, we have $f_{\lambda}^{k-2}\left(v_{\lambda}^{+}\right) \notin \partial B_{\lambda}$ and $f_{\lambda}^{k-1}\left(v_{\lambda}^{+}\right) \in \partial B_{\lambda}$. There is disk neighborhood $\mathscr{D} \subset \mathbb{C}^{*} \backslash \partial \mathscr{H}_{0}$ of $\mathcal{E}$ such that for all $\lambda \in \mathscr{D}, f_{\lambda}^{k-2}\left(v_{\lambda}^{+}\right) \notin \overline{B_{\lambda}}$.

Take two different parameters $\lambda_{1}, \lambda_{2} \in \mathcal{E}$ with $\left|\arg \lambda_{1}-\arg \lambda_{2}\right|<\frac{2 \pi}{n-1} \quad$ and let $J=\left\{f_{\lambda_{1}}^{j}\left(v_{\lambda_{1}}^{\varepsilon}\right) ; 0 \leq j \leq k-2, \varepsilon= \pm\right\} \cup \overline{B_{\lambda_{1}}}$. We define a continuous map $h: \mathscr{D} \times J \rightarrow \widehat{\mathbb{C}}$ in the following way:

1. $h\left(\lambda_{1}, z\right)=z$ for all $z \in J$;
2. $h(\lambda, z)=\phi_{\lambda}^{-1} \circ \phi_{\lambda_{1}}(z)$ for all $(\lambda, z) \in \mathscr{D} \times \overline{B_{\lambda_{1}}}$;
3. for any $\lambda \in \mathscr{D}$, we define $h\left(\lambda, f_{\lambda_{1}}^{j}\left(v_{\lambda_{1}}^{\varepsilon}\right)\right)=f_{\lambda}^{j}\left(v_{\lambda}^{\varepsilon}\right)$ for $0 \leq j \leq k-2$ and $\varepsilon \in\{ \pm\}$.

The map $h$ is a holomorphic motion parameterized by $\mathscr{D}$, with base point $\lambda_{1}$. By Slodkowski's theorem [33], there is a holomorphic motion $H: \mathscr{D} \times \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$ extending $h$. We consider the restriction $H_{0}=\left.H\right|_{\mathscr{E} \times \widehat{\mathbb{C}}}$ of $H$. Note that for any $\lambda \in \mathcal{E}$, the map $H_{0}(\lambda, \cdot)$ preserves the postcritical relation. So there is unique continuous map $H_{1}: \mathcal{E} \times \widehat{\mathbb{C}} \rightarrow \widehat{\mathbb{C}}$ such that $H_{1}\left(\lambda_{1}, \cdot\right) \equiv \mathrm{id}$ and the following diagram commutes:


Set $\psi_{0}=H_{0}\left(\lambda_{2}, \cdot\right)$ and $\psi_{1}=H_{1}\left(\lambda_{2}, \cdot\right)$. Both $\psi_{0}$ and $\psi_{1}$ are quasiconformal maps satisfying $f_{\lambda_{2}} \circ \psi_{1}=\psi_{0} \circ f_{\lambda_{1}}$. One may verify that $\psi_{0}$ and $\psi_{1}$ are homotopic relative to $P\left(f_{\lambda_{1}}\right) \cup \overline{B_{\lambda_{1}}}$. To see this, note that $H_{1}(\lambda, \cdot)^{-1} \circ H_{0}(\lambda, \cdot)$ is continuous in $\lambda \in \mathscr{E}$ with $H_{1}\left(\lambda_{1}, \cdot\right)^{-1} \circ H_{0}\left(\lambda_{1}, \cdot\right)=i d$, this is to say the map $\psi_{1}^{-1} \circ \psi_{0}=H_{1}\left(\lambda_{2}, \cdot\right)^{-1} \circ H_{0}\left(\lambda_{2}, \cdot\right)$ is homotopic to the identity map relative to $P\left(f_{\lambda_{1}}\right) \cup \overline{B_{\lambda_{1}}}$.

Then there is a sequence of quasi-conformal maps $\psi_{j}$ such that
(a) $f_{\lambda_{2}} \circ \psi_{j+1}=\psi_{j} \circ f_{\lambda_{1}}$ for all $j \geq 0$,
(b) $\psi_{j+1}$ and $\psi_{j}$ are homotopic relative to $f_{\lambda_{1}}^{-j}\left(P\left(f_{\lambda_{1}}\right) \cup \overline{B_{\lambda_{1}}}\right)$.

The maps $\psi_{j}$ form a normal family since their dilations are uniformly bounded above. Let $\psi_{\infty}$ be the limit map of $\psi_{j}$. It is holomorphic in the Fatou set $F\left(f_{\lambda_{1}}\right)=\cup_{k} f_{\lambda_{1}}^{-k}\left(B_{\lambda_{1}}\right)$ and satisfies $f_{\lambda_{2}} \circ \psi_{\infty}=\psi_{\infty} \circ f_{\lambda_{1}}$ in $F\left(f_{\lambda_{1}}\right)$. By continuity, $f_{\lambda_{2}} \circ \psi_{\infty}=\psi_{\infty} \circ f_{\lambda_{1}}$ in $\widehat{\mathbb{C}}$.

By Theorem 5.4, the Lebesgue measures of $J\left(f_{\lambda_{1}}\right)$ and $J\left(f_{\lambda_{2}}\right)$ are zero. Thus $\psi_{\infty}$ is a Möbius map. It takes the form $\psi_{\infty}(z)=a z$ where $a^{n-1}=1$ and $\lambda_{2}=a^{2} \lambda_{1}$. The condition $\left|\arg \lambda_{1}-\arg \lambda_{2}\right|<\frac{2 \pi}{n-1}$ implies $\lambda_{1}=\lambda_{2}$. But this is a contradiction.

So the set $\chi_{\mathcal{H}}(t) \backslash \partial \mathscr{H}_{0}$ is either empty or a singleton. This implies that the impression $\chi_{\mathscr{H}}(t)$ is either a singleton or contained in $\partial \mathscr{H}_{0}$.

## Proposition 7.4. - The boundary $\partial \mathscr{H}$ is locally connected.

Proof. - It follows from Lemma 7.3 that for any $t$, the impression $\chi_{\mathscr{H}}(t)$ is either a singleton or contained in $\partial \mathscr{H}_{0}$. In the latter case, for any $\lambda \in \chi_{\mathcal{H}}(t) \cap \mathscr{F}$ which is not a cusp, it follows from Lemma 4.8 that there is an external ray $R_{\lambda}(\alpha)$ landing at $v_{\lambda}^{+}$.

We claim that $n t=n^{k-1} \alpha \bmod 1$. If not, then by Lemma 3.7, there is a cut ray $\Omega_{\lambda}^{\beta}$ separating $R_{\lambda}(n t)$ and $R_{\lambda}\left(n^{k-1} \alpha\right)$. By stability of cut rays, there exists a neighborhood $\mathcal{U}$ of $\lambda$ and $\varepsilon>0$ such that $Z_{n t, \varepsilon}^{u}$ and $Z_{n^{k-1} \alpha, \varepsilon}^{u}$ are contained in different components of $\overline{\mathbb{C}}-\Omega_{u}^{\beta}$ for all $u \in \mathcal{U} \cap\left(\overline{\mathbb{C}}-\mathcal{H}_{0}\right)$, where

$$
Z_{t, \varepsilon}^{u}:=\phi_{u}^{-1}\left(\left\{r e^{2 \pi i \theta} ; r>1,|\theta-t|<\varepsilon\right\}\right) .
$$

Moreover, by shrinking $\mathcal{U}$ a little bit, we see that $f_{u}^{k-1}\left(v_{u}^{+}\right) \in Z_{n^{k-1} \alpha, \varepsilon}^{u}$ for all $u \in \mathscr{H} \cap \mathcal{U}$. Then there is a cut ray $\Omega_{u}^{\eta} \subset f_{u}^{1-k}\left(\Omega_{u}^{\beta}\right)$, separating $v_{u}^{+}$and $\bigcup_{|\theta-t|<\varepsilon / n} R_{U_{u}}(\theta)$ for all $u \in \mathscr{H} \cap \mathcal{U}$. However, by the definition of $\chi_{\mathscr{H}}(t)$, when $k$ is large so that $1 / k<\varepsilon / n$, there is $\lambda_{k} \in \mathcal{U} \cap \Phi_{\mathscr{H}}^{-1}\left(\left\{r e^{2 \pi i \theta} ; 1-1 / k<r<1,|\theta-t|<1 / k\right\}\right)$. So we have $v_{\lambda_{k}}^{+} \subset \bigcup_{|\theta-t|<1 / k} R_{U_{\lambda_{k}}}(\theta) \subset \bigcup_{|\theta-t|<\varepsilon / n} R_{U_{\lambda_{k}}}(\theta)$. But this is a contradiction. This completes the proof of the claim.

Thus each $\lambda \in \chi_{\mathscr{H}}(t)$ is either a cusp or contained in $\left\{\nu(\alpha) ; n t=n^{k-1} \alpha\right\}$ (a finite set). The connectivity of $\chi_{\mathscr{H}}(t)$ implies that it is a singleton.

Theorem 7.5. - The boundary $\partial \mathscr{H}$ is a Jordan curve.
Proof. - If not, then there exist a parameter $\lambda \in \partial \mathscr{H}$ with $0 \leq \arg \lambda \leq \frac{\pi}{n-1}$ (by assumption of $\mathscr{H}$ ) and two different angles $t_{1}, t_{2}$ such that $\chi_{\mathscr{H}}\left(t_{1}\right)=\chi_{\mathcal{H}}\left(t_{2}\right)=\{\lambda\}$.

By Lemma 3.7, there is a cut ray $\Omega_{\lambda}^{\alpha}$ separating the internal rays $R_{U_{\lambda}}\left(t_{1}\right)$ and $R_{U_{\lambda}}\left(t_{2}\right)$. Suppose that $v_{\lambda}^{+}$and $R_{U_{\lambda}}\left(t_{1}\right)$ are contained in the same component of $\widehat{\mathbb{C}}-\Omega_{\lambda}^{\alpha}$. By the stability of cut rays, there is a neighborhood $\mathcal{U}$ of $\lambda$ such that for any $u \in \mathscr{U} \cap \mathcal{H}$, the set $\left\{v_{u}^{+}\right\} \cup R_{U_{u}}\left(t_{1}\right)$ and the internal ray $R_{U_{u}}\left(t_{2}\right)$ are contained in different components of $\widehat{\mathbb{C}}-\Omega_{u}^{\alpha}$. But this contradicts the assumption that $\mathcal{X}_{\mathscr{H}}\left(t_{2}\right)=\{\lambda\}$.

## 8. Hyperbolic components of renormalizable type

In this section, we study the hyperbolic components of renormalizable type.
We begin with a definition. We say a McMullen map $f_{\lambda}$ is renormalizable (resp. *-renormalizable) at $c \in C_{\lambda}$ if there exist an integer $p \geq 1$ and two disks $U$ and $V$ containing $c$, such that $f_{\lambda}^{p}: U \rightarrow V$ (resp. $-f_{\lambda}^{p}: U \rightarrow V$ ) is a quadratic-like map whose Julia set is connected. The triple $\left(f_{\lambda}^{p}, U, V\right)$ (resp. $\left(-f_{\lambda}^{p}, U, V\right)$ ) is called the renormalization (resp. $*$-renormalization) of $f_{\lambda}$.

Let $\mathscr{B}$ be a hyperbolic component of renormalizable type. For any $\lambda \in \mathscr{B}$, the map $f_{\lambda}$ has an attracting cycle in $\mathbb{C}$, say $z_{\lambda} \mapsto f_{\lambda}\left(z_{\lambda}\right) \mapsto \cdots \mapsto f_{\lambda}^{p}\left(z_{\lambda}\right)=z_{\lambda}$, where $p$ is the period. We may assume that the attracting cycle is suitably chosen and labeled so that $z_{\lambda}$ is holomorphic with respect to $\lambda \in \mathscr{B}$.

Lemma 8.1 ([28], Prop 5.4). - If $\lambda \in \mathcal{B}$, then $f_{\lambda}$ is either renormalizable or $*$-renormalizable. Moreover,

1. if $f_{\lambda}$ is renormalizable and $n$ is odd, then $f_{\lambda}$ has exactly two attracting cycles in $\mathbb{C}$;
2. if $f_{\lambda}$ is $*$-renormalizable and $n$ is odd, then $p$ is even, $f_{\lambda}^{p / 2}\left(z_{\lambda}\right)=-z_{\lambda}$ and $f_{\lambda}$ has exactly one attracting cycle in $\mathbb{C}$;
3. if $n$ is even, then $f_{\lambda}$ has exactly one attracting cycle in $\mathbb{C}$ and there is a unique $c \in C_{\lambda}$, such that $f_{\lambda}$ is renormalizable at $c$.

The terminology 'hyperbolic component of renormalizable type' comes from Lemma 8.1. Let $\rho(\lambda)=\left(f_{\lambda}^{p}\right)^{\prime}\left(z_{\lambda}\right)$ be the multiplier of the attracting cycle of $f_{\lambda}$ for $\lambda \in \mathcal{B}$. Based on Lemma 8.1, we set $(\epsilon, k)=(-1, p / 2)$ if $n$ is odd and $f_{\lambda}$ is $*$-renormalizable, and $(\epsilon, k)=(1, p)$ in the other cases. We define a map $\kappa: \mathcal{B} \rightarrow \mathbb{D}$ by $\kappa(\lambda)=\left(\epsilon f_{\lambda}^{k}\right)^{\prime}\left(z_{\lambda}\right)$. Note that either $\rho=\kappa^{2}$ or $\rho=\kappa$.

The main result of this section is:
Theorem 8.2. - The map $\kappa: \mathscr{B} \rightarrow \mathbb{D}$ is a conformal map. It can be extended continuously to a homeomorphism from $\overline{\mathcal{B}}$ to $\overline{\mathbb{D}}$.

Proof. - Note that $\kappa(\lambda)$ is the multiplier of the map $g_{\lambda}=\epsilon f_{\lambda}^{k}$ at its fixed point $z_{\lambda}$. By the Implicit Function Theorem, if $\mathscr{B} \ni \lambda_{n} \rightarrow \partial \mathscr{B}$, then $\left|\kappa\left(\lambda_{n}\right)\right| \rightarrow 1$, so the map $\kappa: \mathcal{B} \rightarrow \mathbb{D}$ is proper.

In the following, we will show that $\kappa$ is actually a covering map. To this end, we will construct a local inverse map of $\kappa$ by means of quasiconformal surgery. The idea is similar to the quadratic case [5].

Fix $\lambda_{0} \in \mathcal{B}$ and set $\kappa_{0}=\kappa\left(\lambda_{0}\right)$. We may relabel $z_{\lambda_{0}}$ so that the immediate attracting basin $A_{0}$ of $z_{\lambda_{0}}$ contains a critical point $c \in C_{\lambda_{0}}$. Note that $\epsilon f_{\lambda}^{k}\left(A_{0}\right)=A_{0}$ and there is a conformal map $\phi: A_{0} \rightarrow \mathbb{D}$ such that $\phi\left(z_{\lambda_{0}}\right)=0$ and the following diagram commutes:

where $B_{\zeta}$ is the Blaschke product defined by $B_{\zeta}(z)=z \frac{z+\zeta}{1+\bar{\zeta} z}$. Obviously $z=0$ is an attracting fixed point of $B_{\kappa_{0}}$ with multiplier $B_{\zeta}^{\prime}(0)=\zeta$. Then there is a neighborhood $\mathscr{U}$ of $\kappa_{0}$ and a continuous family of quasiregular maps $\widetilde{B}: U \times \mathbb{D} \rightarrow \mathbb{D}$ such that $\widetilde{B}\left(\kappa_{0}, \cdot\right)=B_{\kappa_{0}}(\cdot)$ and $\widetilde{B}(\zeta, z)=B_{\kappa_{0}}(z)$ for $\varepsilon<|z|<1\left(\varepsilon>0\right.$ is a small number), $\widetilde{B}(\zeta, z)=B_{\zeta}(z)$ for $|z|<\varepsilon / 2$ and $\widetilde{B}(\zeta, \cdot)$ is quasi-regular elsewhere.

Then we get a continuous family $\left\{G_{\zeta}\right\}_{\zeta \in \mathcal{U}}$ of quasiregular maps:

$$
G_{\zeta}(z)= \begin{cases}(-1)^{q}\left(\left.f_{\lambda_{0}}^{k-1}\right|_{f_{\lambda_{0}}\left(A_{0}\right)}\right)^{-1}\left(\epsilon \phi^{-1} \widetilde{B}\left(\zeta, \phi\left(e^{-q \pi i / n} z\right)\right)\right), & z \in e^{q \pi i / n} A_{0}, 0 \leq q<2 n \\ f_{\lambda_{0}}(z), & z \in \widehat{\mathbb{C}} \backslash \bigcup_{0 \leq q<2 n} e^{q \pi i / n} A_{0}\end{cases}
$$

We can construct a $G_{\zeta}$-invariant complex structure $\sigma_{\zeta}$ such that

- $\sigma_{\kappa_{0}}$ is the standard complex structure $\sigma_{0}$ on $\widehat{\mathbb{C}}$.
- $\sigma_{\zeta}$ is continuous with respect to $\zeta \in \mathcal{U}$.
- $\sigma_{\zeta}$ is invariant under the maps $z \mapsto e^{2 \pi i / n} z$ and $z \mapsto-z$.
- $\sigma_{\zeta}$ is the standard complex structure near the attracting cycle and outside $\bigcup_{k \geq 0} f_{\lambda_{0}}^{-k}\left(\bigcup_{0 \leq q<2 n} e^{q \pi i / n} A_{0}\right)$.

The Beltrami coefficient $\mu_{\zeta}$ of $\sigma_{\zeta}$ satisfies $\left\|\mu_{\zeta}\right\|<1$. By the Measurable Riemann Mapping Theorem [3], there is a continuous family of quasiconformal maps $\psi_{\zeta}$ fixing $0, \infty$ and normalized so that $\psi_{\zeta}^{\prime}(\infty)=1$. The map $\psi_{\zeta}$ satisfies $\psi_{\zeta}\left(e^{2 \pi i / n} z\right)=e^{2 \pi i / n} \psi_{\zeta}(z)$ and $\psi_{\zeta}(-z)=-\psi_{\zeta}(z)$. Then $F_{\zeta}=\psi_{\zeta} \circ G_{\zeta} \circ \psi_{\zeta}^{-1}$ is a rational map of the form $z^{-n}\left(z^{2 n}+\sum_{0 \leq k<2 n} b_{k}(\zeta) z^{k}\right)$. The symmetry $F_{\zeta}\left(e^{2 \pi i / n} z\right)=F_{\zeta}(z)$ implies $F_{\zeta}(z)=z^{n}+b_{0}(\zeta) z^{-n}+b_{n}(\zeta)$. Since the two free critical values of $F_{\zeta}$ satisfy $\psi_{\zeta}\left(v_{\lambda_{0}}^{+}\right)+\psi_{\zeta}\left(v_{\lambda_{0}}^{-}\right)=0$, we have $b_{n}(\zeta)=0$. So $F_{\zeta}=f_{b_{0}(\zeta)}$. The coefficient $b_{0}: \mathcal{U} \rightarrow \mathcal{B}$ is continuous with $\kappa\left(b_{0}(\zeta)\right)=\zeta$. So $b_{0}$ is the local inverse of $\kappa$. This implies $\kappa$ is a covering map. Since $\mathbb{D}$ is simply connected, $\kappa$ is actually a conformal map.

The map $\kappa$ has a continuation to the boundary $\partial \mathscr{B}$. By the Implicit Function Theorem, the boundary $\partial \mathscr{B}$ is an analytic curve except at $\kappa^{-1}(1)$. So $\partial \mathscr{B}$ is locally connected. Since for any $\lambda \in \partial \mathscr{B}$, the multiplier $e^{2 \pi i t}$ of the non-repelling cycle of $f_{\lambda}$ is uniquely determined by its angle $t \in \mathbb{S}$, the boundary $\partial \mathscr{B}$ is a Jordan curve.

Remark 8.3. - By Theorem 8.2, the multiplier map $\rho: \mathcal{B} \rightarrow \mathbb{D}$ is a double cover if and only if $n$ is odd and $f_{\lambda}$ is $*$-renormalizable. For example, when $n=3$, let $\mathcal{B}_{+}$(resp. $\mathcal{B}_{-}$) be the cardioid of the 'largest' baby Mandelbrot set intersecting the positive (resp. negative) real axis, then

1. $\rho: \mathscr{B}_{+} \rightarrow \mathbb{D}$ is a conformal map and near the center $\frac{1}{8}$ of $\mathscr{B}_{+}$,

$$
\rho(\lambda)=24\left(\lambda-\frac{1}{8}\right)+(216+156 \sqrt{2})\left(\lambda-\frac{1}{8}\right)^{2}+\Theta\left(\left(\lambda-\frac{1}{8}\right)^{3}\right) .
$$

2. $\rho: \mathcal{B}_{-} \rightarrow \mathbb{D}$ is a double covering and near the center $-\frac{1}{8}$ of $\mathcal{B}_{-}$,

$$
\rho(\lambda)=576\left(\lambda+\frac{1}{8}\right)^{2}+\theta\left(\left(\lambda+\frac{1}{8}\right)^{3}\right) .
$$

As a concluding remark, we would like to mention that and using a result of McMullen [26], we can conclude that the parameter plane of $f_{\lambda}$ contains many quasiconformal copies of the Mandelbrot set. In fact, every hyperbolic component $\mathcal{B}$ of renormalizable type is a quasiconformal image of a hyperbolic component of the Mandelbrot set, as we can see in Figure 1.

## 9. Appendix: Lebesgue measure

In this appendix, we shall prove Theorem 5.4 based on the Yoccoz puzzle theory. See Sections 5.1 and 5.4 for basic introductions. The notations also follow there.

THEOREM 9.1 (Lebesgue measure). - Suppose that $\lambda \in \mathcal{M} \cap \mathcal{F}$ and the graph $G_{\lambda}\left(\theta_{1}, \ldots, \theta_{N}\right)$ is admissible. If none of $T(c)$ with $c \in C_{\lambda}$ is periodic, then the Lebesgue measure of $J\left(f_{\lambda}\right)$ is zero.

Proof of Theorem 5.4 assuming Theorem 9.1. - We assume $\lambda \in \mathcal{F}_{0}$ (note that when $\lambda$ is real and positive, the map $f_{\lambda}$ is postcritically finite). It is known that if $f_{\lambda}$ is postcritically finite, then the Lebesgue measure of $J\left(f_{\lambda}\right)$ is zero. So we assume further $\lambda \in \mathcal{F}$ and $f_{\lambda}$ is postcritically infinite. By Lemma 5.2, there is an admissible graph $G_{\lambda}\left(\theta_{1}, \ldots, \theta_{N}\right)$. By the assumption $f_{\lambda}^{k}\left(v_{\lambda}^{+}\right) \in \partial B_{\lambda}$, none of $T(c)$ with $c \in C_{\lambda}$ is periodic. It follows from Theorem 9.1 that the Lebesgue measure of $J\left(f_{\lambda}\right)$ is zero.

In this section, we will actually prove Theorem 9.1 following Lyubich [18].
For $k \geq 0$, let $\mathscr{P}_{k}$ be the collection of all puzzle pieces of depth $k$. We first note that

Lemma 9.2. - We have $d_{k}=\max \left\{\operatorname{diam}(P) ; P \in \mathscr{P}_{k}\right\} \rightarrow 0$ as $k \rightarrow \infty$.

Proof. - If not, then there exist $\varepsilon>0$ and a sequence of puzzle pieces $P_{n_{k}} \in \mathscr{P}_{n_{k}}$ with $n_{1}<n_{2}<\cdots$ and $\operatorname{diam}\left(P_{n_{k}}\right) \geq \varepsilon$. There is $P_{n_{1}}^{*} \in \mathscr{P}_{n_{1}}$ such that $I_{1}=\left\{n_{k} ; P_{n_{k}} \subset P_{n_{1}}^{*}\right\}$ is an infinite set. For $k>1$, we define $P_{n_{k}}^{*}$ and $I_{k}$ inductively as follows: $P_{n_{k-1}}^{*} \supset P_{n_{k}}^{*} \in \mathscr{P}_{n_{k}}$ and the set $I_{k}=\left\{j \in I_{k-1} ; P_{j} \subset P_{n_{k}}^{*}\right\}$ is an infinite set. Then $P_{n_{k}}^{*}, k \geq 1$, is a sequence of shrinking puzzle pieces with $\operatorname{diam}\left(P_{n_{k}}^{*}\right) \geq \varepsilon$. This contradicts the fact that $\bigcap_{k} \overline{P_{n_{k}}^{*}}$ consists of a single point (see Lemma 5.3).

Lemma 9.3. - Let $U, V$ be two planar disks with $V \Subset U \neq \mathbb{C}, x \in V$. Suppose that $\operatorname{Shape}(U, x) \leq C, \quad \operatorname{Shape}(V, x) \leq C, \quad \bmod (U-\bar{V}) \leq m, \quad$ then there is a constant $\delta=\delta(C, m) \in(0,1)$, such that $\operatorname{area}(V) \geq \delta \operatorname{area}(U)$.

Here, the notation $V \Subset U$ means that $V$ is compactly contained in $U$, i.e., $\bar{V} \subset U$. The proof of Lemma 9.3 is based on the Koebe distortion theorem. We leave it to the reader as an exercise.

Lemma 9.4. - Let $f$ be a rational map with Julia set $J(f) \neq \widehat{\mathbb{C}}$. Let $z \in J(f)$, if there exist a number $\epsilon>0$, a sequence of integers $0 \leq n_{1}<n_{2}<\cdots$ and a constant $N>0$ such that

1. for any $k \geq 0$, the component $U_{k}(z)$ of $f^{-n_{k}}\left(B\left(f^{n_{k}}(z), \epsilon\right)\right)$ that contains $z$ is a disk;
2. $\operatorname{deg}\left(\left.f^{n_{k}}\right|_{U_{k}(z)}\right) \leq N$ for all $k \geq 1$.

Then $z$ is not a Lebesgue density point of $J(f)$.

Proof. - By passing to a subsequence, we assume $f^{n_{k}}(z) \rightarrow w \in J(f)$ as $k \rightarrow \infty$. We may assume further $z, w \neq \infty$ by a suitable change of coordinate. Choose $\epsilon_{0}<\epsilon$, when $k$ is large, we have $f^{n_{k}}(z) \in B\left(w, \epsilon_{0} / 2\right) \subset B\left(w, \epsilon_{0}\right) \subset B\left(f^{n_{k}}(z), \epsilon\right)$. Let $V_{k}(z)$ be the component of $f^{-n_{k}}\left(B\left(w, \epsilon_{0} / 2\right)\right)$ that contains $z$. Then $V_{k}(z)$ is a disk and $\operatorname{deg}\left(\left.f^{n_{k}}\right|_{V_{k}(z)}\right) \leq N$. By shape distortion (see [28], Lemma 6.1), the shape of $V_{k}(z)$ about $z$ is bounded above by some constant depending on $N$. We then show that $\operatorname{diam}\left(V_{k}(z)\right) \rightarrow 0$ as $k \rightarrow \infty$. In fact, if not, again by choosing a subsequence, we assume $V_{k}(z)$ contains a round disk $B(z, \rho)$ for some $\rho>0$. Then for any large $k$, the image $f^{n_{k}}(B(z, \rho))$ is contained in $B\left(w, \epsilon_{0} / 2\right)$. But this contradicts the fact that $J(f) \subset f^{n_{k}}(B(z, \rho))$ for large $k$ (see [27]).

Since $J(f) \neq \widehat{\mathbb{C}}$, there is a round disk $B(\zeta, r) \Subset B\left(w, \epsilon_{0} / 2\right) \cap F(f)$; here $F(f)$ is the Fatou set of $f$. Take a component $D_{k}$ of $f^{-n_{k}}(B(\zeta, r))$ in $V_{k}(z)$ and $p \in f^{-n_{k}}(\zeta) \cap D_{k}$, then by shape distortion (see [28], Lemma 6.1), there is a constant $C>0$ such that $\operatorname{Shape}\left(D_{k}, p\right) \leq C \operatorname{Shape}(B(\zeta, r), \zeta)=C \operatorname{Shape}\left(V_{k}(z), p\right) \leq C \operatorname{Shape}\left(B\left(w, \epsilon_{0} / 2\right), \zeta\right) \leq C \epsilon_{0} / r$. Moreover, $\bmod \left(V_{k}(z) \backslash \overline{D_{k}}\right) \leq \bmod \left(B\left(w, \epsilon_{0} / 2\right) \backslash \overline{B(\zeta, r)}\right)$. It follows from Lemma 9.3 that there is a constant $\delta$ with area $\left(D_{k}\right) \geq \delta \operatorname{area}\left(V_{k}(z)\right)$. So

$$
\operatorname{area}\left(J(f) \cap V_{k}(z)\right) \leq \operatorname{area}\left(V_{k}(z)-D_{k}\right) \leq(1-\delta) \operatorname{area}\left(V_{k}(z)\right) .
$$

This implies $z$ is not a Lebesgue density point.
Proposition 9.5. - If $C_{\lambda}$ is not persistently recurrent, then the Lebesgue measure of the Julia set $J\left(f_{\lambda}\right)$ is zero.

Proof. - Let $P\left(f_{\lambda}\right)=\overline{\bigcup_{k \geq 1} f_{\lambda}^{k}\left(C_{\lambda}\right)}$. It is known ([24], Theorem 3.9) that for almost all $z \in J\left(f_{\lambda}\right)$, the spherical distance $d_{\widehat{\mathbb{C}}}\left(f_{\lambda}^{n}(z), P\left(f_{\lambda}\right)\right) \rightarrow 0$ as $n \rightarrow \infty$.

Suppose that the critical set $C_{\lambda}$ is recurrent but not persistently recurrent; then there is a positive integer $L$ such that the set $\left\{k ; \tau_{Y}(k) \leq L\right\}$ is infinite. The recurrence of $C_{\lambda}$ implies that there is $d \geq L$ such that the annulus $P_{d}^{\lambda}(c) \backslash \overline{P_{d+1}^{\lambda}(c)}$ for some (hence all) $c \in C_{\lambda}$ is non-degenerate. Since $\tau_{Y}(k+1) \leq \tau_{Y}(k)+1$, the set $\Lambda=\left\{k ; \tau_{Y}(k)=d, \tau_{Y}(k+1)=d+1\right\}$ is infinite. Moreover $\left\{f_{\lambda}^{k-d}(c) ; k \in \Lambda\right\} \subset \bigcup_{\zeta \in C_{\lambda}} P_{d+1}(\zeta) \Subset \bigcup_{\zeta \in C_{\lambda}} P_{d}(\zeta)$. So any critical point $c \in C_{\lambda}$ satisfies the conditions in Lemma 9.4, thus it is not a Lebesgue density point. We consider a point $z \in J\left(f_{\lambda}\right) \backslash C_{\lambda}$ with $\lim d_{\widehat{\mathbb{C}}}\left(f_{\lambda}^{n}(z), P\left(f_{\lambda}\right)\right)=0$. We may assume that the forward orbit of $z$ does not meet the graph $G_{\lambda}\left(\theta_{1}, \ldots, \theta_{N}\right)$ (for else $z$ is not a Lebesgue density point by Lemma 9.4). In that case, for each $k \in \Lambda$, there is $n_{k}>k$ and $c^{\prime} \in C_{\lambda}$ such that $f_{\lambda}^{n_{k}-k-1}\left(P_{n_{k}}^{\lambda}(z)\right)=P_{k+1}^{\lambda}\left(c^{\prime}\right)$ and $f_{\lambda}^{j}\left(P_{n_{k}}^{\lambda}(z)\right), 0 \leq j<n_{k}-k-1$, meets no critical point. One can easily verify that $z$ satisfies the conditions in Lemma 9.4, and is not a Lebesgue density point of $J\left(f_{\lambda}\right)$.

If the critical set $C_{\lambda}$ is not recurrent, one can verify that each point $z \in J\left(f_{\lambda}\right)$ satisfies the condition in Lemma 9.4. Thus $J\left(f_{\lambda}\right)$ carries no Lebesgue density point. The proof is similar to, but easier than the previous argument. We omit the details.

We say a holomorphic map $g: \mathbf{U} \rightarrow \mathbf{V}$ is a repelling system if $\mathbf{U} \Subset \mathbf{V}$, the boundary $\partial \mathbf{U}$ avoids the critical orbit of $g$ and both $\mathbf{U}$ and $\mathbf{V}$ consist of finitely many disk components. The filled Julia set of $g$ is defined by $K(g)=\bigcap_{k \geq 1} g^{-k}(\mathbf{V})$, it can be an empty set.

Theorem 9.6. - If the critical set $C_{\lambda}$ is persistently recurrent, then there is a repelling system $g: \mathbf{U} \rightarrow \mathbf{V}$ such that


Figure 6. A repelling system $g: \mathbf{U} \rightarrow \mathbf{V}$, where $\mathbf{U}$ is the union of all shadow disks and $\mathbf{V}$ is the union of six larger disks.

1. each component of $\mathbf{U}$ and $\mathbf{V}$ is a puzzle piece;
2. for each component $U_{i}$ of $\mathbf{U},\left.g\right|_{U_{i}}=f_{\lambda}^{l_{i}}$ for some $l_{i}$.
3. $C_{\lambda} \subset K(g)$.

Moreover, the Lebesgue measure of $J\left(f_{\lambda}\right)-\bigcup_{k \geq 0} f_{\lambda}^{-k}(K(g))$ is zero.
Proof. - Since the graph $G_{\lambda}\left(\theta_{1}, \ldots, \theta_{N}\right)$ is admissible, we can find a non-degenerate critical annulus $P_{d}^{\lambda}(c) \backslash \overline{P_{d+1}^{\lambda}(c)}$ for some $d \geq 1$. Set $\mathbf{V}=\bigcup_{c \in C_{\lambda}} P_{d+1}^{\lambda}(c)$. Then $f_{\lambda}^{j}(\partial \mathbf{V}) \cap \overline{\mathbf{V}}=\varnothing$ for all $j \geq 1$. For any $j \geq 1$, either $f_{\lambda}^{j}\left(C_{\lambda}\right) \subset \mathbf{V}$ or $f_{\lambda}^{j}\left(C_{\lambda}\right) \cap \mathbf{V}=\varnothing$. Let $1 \leq n_{1}<n_{2}<\cdots$ be all the integers such that $f_{\lambda}^{n_{i}}\left(C_{\lambda}\right) \subset \mathbf{V}$. Let $l_{i}=n_{i+1}-n_{i}$ (set $n_{0}=0$ ) for $i \geq 0$, we pull back $\mathbf{V}$ along the orbit $\left\{f_{\lambda}^{j}\left(C_{\lambda}\right)\right\}_{j=n_{i}}^{n_{i+1}}$ and get $\mathbf{V}_{i}$. Namely, $\mathbf{V}_{i}$ is the union of all components of $f_{\lambda}^{-l_{i}}(\mathbf{V})$ intersecting with $f_{\lambda}^{n_{i}}\left(C_{\lambda}\right)$. For any $i$, the intermediate pieces $f_{\lambda}^{k}\left(\mathbf{V}_{i}\right), 0<k<l_{i}$ lie outside $\mathbf{V}$ and for any component $V$ of $\mathbf{V}_{i}$, the $\left.\operatorname{map} f_{\lambda}^{l_{i}}\right|_{V}$ is either univalent or a double covering.

Since $C_{\lambda}$ is persistently recurrent, the set $\left\{k ; \tau_{Y}(k) \leq d+1\right\}$ is finite and there are only finitely may different $\mathbf{V}_{i}$ 's. Moreover, if $\mathbf{V}_{i} \neq \mathbf{V}_{j}$, then $\mathbf{V}_{i} \cap \mathbf{V}_{j}=\varnothing$ (in fact, $\overline{\mathbf{V}}_{i} \cap \overline{\mathbf{V}}_{j}=\varnothing$ ).

Let $\mathbf{U}=\bigcup_{i} \mathbf{V}_{i}$ and define $\left.g\right|_{\mathbf{V}_{i}}=f_{\lambda}^{l_{i}}$. Then $\mathbf{U} \Subset \mathbf{V}$ follows from the fact that $f_{\lambda}^{l_{i}}\left(\partial \mathbf{V}_{i} \cap \partial \mathbf{V}\right) \subset \partial \mathbf{V} \cap f_{\lambda}^{l_{i}}(\partial \mathbf{V})=\varnothing$.

It follows from $\bigcup_{i \geq 0} f_{\lambda}^{n_{i}}\left(C_{\lambda}\right)=\bigcup_{k \geq 0} g^{k}\left(C_{\lambda}\right) \subset \mathbf{V}$ that $C_{\lambda} \subset K(g)$.
Similar to the proof of Proposition 9.5, we need only consider a point $z \in J\left(f_{\lambda}\right)$ with $d_{\widehat{\mathbb{C}}}\left(f_{\lambda}^{n}(z), P\left(f_{\lambda}\right)\right) \rightarrow 0$ as $n \rightarrow \infty$. For such point, there is an integer $N>0$ such that for all $n \geq N, f_{\lambda}^{n}(z) \in \mathbf{V}$ implies $f_{\lambda}^{n}(z) \in \mathbf{U}$. Note that there is $p \geq N$ such that $f_{\lambda}^{p}(z) \in \mathbf{V}$. Then for all $j \geq 1$, we have $g^{j}\left(f_{\lambda}^{p}(z)\right) \in \mathbf{V}$. It turns out that $f_{\lambda}^{p}(z) \in K(g)$. This implies $J\left(f_{\lambda}\right)-\bigcup_{k \geq 0} f_{\lambda}^{-k}(K(g))$ has zero Lebesgue measure.

Let $D \subset \mathbb{C}$ be a topological disk containing a compact subset $K$ (not necessarily connected), the modulus of $A=D-K$, denoted by $\mathbf{m}(A)$, is defined to be the extremal length of curves joining $\partial D$ and $\partial K$. It is equal to the reciprocal of Dirichlet integral of the harmonic
measure $u$ in $A$ (namely, $u$ is harmonic function in $A$ which tends to 0 at regular points of $\partial K$ and tends to 1 at regular points of $\partial D)$ :

$$
\mathbf{m}(A)=\left(\int_{A}|\nabla u|^{2} d x d y\right)^{-1}
$$

If we require further that $K$ consists of finitely many components, then we have the following area-modulus inequality (see [18]):

$$
\operatorname{area}(D) \geq \operatorname{area}(K)(1+4 \pi \mathbf{m}(A))
$$

Now we consider the repelling system $g: \mathbf{U} \rightarrow \mathbf{V}$ defined in Theorem 9.6. Set $\mathbf{V}^{0}=\mathbf{V}$ and consider the preimages $\mathbf{V}^{d}=g^{-d}(\mathbf{V})$ for $d \geq 1$. Note that $\mathbf{V}^{d+1} \Subset \mathbf{V}^{d}$. For any $z \in K(g)$ and $d \geq 0$, denote by $\mathbf{V}^{d}(z)$ the piece of level $d$ containing $z$. Let $\mathbf{A}^{d}(z)=\mathbf{V}^{d}(z)-\overline{\mathbf{V}^{d+1}}$, it is a multiconnected domain. One can verify that for any $d \geq 1$, if $\mathbf{V}^{d}(z)$ contains no critical point in $C_{\lambda}$, then $\mathbf{m}\left(\mathbf{A}^{d}(z)\right)=\mathbf{m}\left(\mathbf{A}^{d-1}\left(f_{\lambda}(z)\right)\right)$; if $\mathbf{V}^{d}(z)$ contains a critical point in $C_{\lambda}$, then $2 \mathbf{m}\left(\mathbf{A}^{d}(z)\right)=\mathbf{m}\left(\mathbf{A}^{d-1}\left(f_{\lambda}(z)\right)\right)$.

Using the same method as in [18], one can show that
Lemma 9.7. - For any $z \in K(g)$, we have $\sum_{d \geq 1} \mathbf{m}\left(\mathbf{A}^{d}(z)\right)=\infty$. It turns out that $K(g)$ is a Cantor set.

Now we have
Theorem 9.8. - Let $g: \mathbf{U} \rightarrow \mathbf{V}$ be the repelling system defined in Theorem 9.6, then the Lebesgue measure of $K(g)$ is zero.

Proof. - For any $d \geq 1$, let $\mathbf{V}^{d}\left(z_{1}\right), \ldots, \mathbf{V}^{d}\left(z_{k_{d}}\right)$ be all puzzle pieces of level $d$, where $z_{1}, \cdots, z_{k_{d}} \in K(g)$. We define

$$
M_{d}=\min _{1 \leq i \leq k_{d}} \sum_{0 \leq j<d} \mathbf{m}\left(\mathbf{A}^{j}\left(z_{i}\right)\right)
$$

By Lemma 9.7, we have $M_{d} \rightarrow \infty$ as $d \rightarrow \infty$. By area-modulus inequality, we have

$$
\operatorname{area}\left(\mathbf{V}^{d}\right) \leq \frac{\operatorname{area}(\mathbf{V})}{\min _{1 \leq i \leq k_{d}} \prod_{0 \leq j<d}\left(1+4 \pi \mathbf{m}\left(\mathbf{A}^{j}\left(z_{i}\right)\right)\right)} \leq \frac{\operatorname{area}(\mathbf{V})}{1+4 \pi M_{d}}
$$

This implies area $\left(\mathbf{V}^{d}\right) \rightarrow 0$ as $d \rightarrow \infty$.
Theorem 9.1 then follows from Proposition 9.5 and Theorems 9.6 and 9.8.

## BIBLIOGRAPHY

[1] Complex dynamics. Twenty-five years after the appearance of the Mandelbrot set, in Proceedings of the AMS-IMS-SIAM Joint Summer Research Conference held in Snowbird, UT, June 13-17, 2004 (R. L. Devaney, L. Keen, eds.), Contemporary Mathematics 396, Amer. Math. Soc., Providence, RI, 2006.
[2] L. V. Ahlfors, Complex analysis, third ed., McGraw-Hill Book Co., New York, 1978.
[3] L. V. Ahlfors, L. Bers, Riemann's mapping theorem for variable metrics, Ann. of Math. 72 (1960), 385-404.
[4] E. Borel, Les probabilités dénombrables et leurs applications arithmétiques, Rendiconti del Circolo Matematico di Palermo 27 (1909), 247-271.
[5] L. Carleson, T. W. Gamelin, Complex dynamics, Universitext: Tracts in Mathematics, Springer, New York, 1993.
[6] R. L. Devaney, Structure of the McMullen domain in the parameter planes for rational maps, Fund. Math. 185 (2005), 267-285.
[7] R. L. Devaney, The McMullen domain: satellite Mandelbrot sets and Sierpiński holes, Conform. Geom. Dyn. 11 (2007), 164-190.
[8] R. L. Devaney, Intertwined internal rays in Julia sets of rational maps, Fund. Math. 206 (2009), 139-159.
[9] R. L. Devaney, Singular perturbations of complex polynomials, Bull. Amer. Math. Soc. (N.S.) 50 (2013), 391-429.
[10] R. L. Devaney, D. M. Look, D. Uminsky, The escape trichotomy for singularly perturbed rational maps, Indiana Univ. Math. J. 54 (2005), 1621-1634.
[11] R. L. Devaney, K. M. Pilgrim, Dynamic classification of escape time Sierpiński curve Julia sets, Fund. Math. 202 (2009), 181-198.
[12] R. L. Devaney, E. Russell, Connectivity of Julia sets for singularly perturbed rational maps, in Chaos, CNN, Memristors and beyond: A Festschrift for Leon Chua (A. Adamatzky et al., eds.), World Scientific Publishing Co., 2013, 239-245.
[13] A. Douady, J. H. Hubbard, A proof of Thurston's topological characterization of rational functions, Acta Math. 171 (1993), 263-297.
[14] F. P. Gardiner, Y. Jiang, Z. Wang, Holomorphic motions and related topics, in Geometry of Riemann surfaces, London Math. Soc. Lecture Note Ser. 368, Cambridge Univ. Press, Cambridge, 2010, 156-193.
[15] P. Haïssinsky, K. M. Pilgrim, Quasisymmetrically inequivalent hyperbolic Julia sets, Rev. Mat. Iberoam. 28 (2012), 1025-1034.
[16] O. Kozlovski, W. Shen, S. van Strien, Rigidity for real polynomials, Ann. of Math. 165 (2007), 749-841.
[17] O. Kozlovski, S. van Strien, Local connectivity and quasi-conformal rigidity of non-renormalizable polynomials, Proc. Lond. Math. Soc. 99 (2009), 275-296.
[18] M. Lyubich, On the Lebesgue measure of the Julia set of a quadratic polynomial, preprint arXiv:math/9201285.
[19] R. Mañé, Hyperbolicity, sinks and measure in one-dimensional dynamics, Comm. Math. Phys. 100 (1985), 495-524.
[20] R. Mañé, Erratum: "Hyperbolicity, sinks and measure in one-dimensional dynamics" [Comm. Math. Phys. 100 (1985), 495-524], Comm. Math. Phys. 112 (1987), 721-724.
[21] R. Mañé, P. Sad, D. Sullivan, On the dynamics of rational maps, Ann. Sci. École Norm. Sup. 16 (1983), 193-217.
[22] C. T. McMullen, Automorphisms of rational maps, in Holomorphic functions and moduli, Vol. I (Berkeley, CA, 1986), Math. Sci. Res. Inst. Publ. 10, Springer, New York, 1988, 31-60.
[23] C. T. McMullen, Cusps are dense, Ann. of Math. 133 (1991), 217-247.
[24] C. T. McMullen, Complex dynamics and renormalization, Annals of Math. Studies 135, Princeton Univ. Press, Princeton, NJ, 1994.
[25] C. T. McMullen, Rational maps and Teichmüller space: analogies and open problems, in Linear and Complex Analysis Problem Book (V. P. Havin, N. K. Nikolskii, eds.), Lecture Notes in Math. 1574, 1994, 430-433.
[26] C. T. McMullen, The Mandelbrot set is universal, in The Mandelbrot set, theme and variations, London Math. Soc. Lecture Note Ser. 274, Cambridge Univ. Press, Cambridge, 2000, 1-17.
[27] J. Milnor, Dynamics in one complex variable, 2nd ed., Friedr. Vieweg \& Sohn, Braunschweig, 2000.
[28] W. Qiu, X. Wang, Y. Yin, Dynamics of McMullen maps, Adv. Math. 229 (2012), 2525-2577.
[29] W. Qiv, Y. Yin, Proof of the Branner-Hubbard conjecture on Cantor Julia sets, Sci. China Ser. A 52 (2009), 45-65.
[30] P. Roesch, Topologie locale des méthodes de Newton cubiques, thèse de doctorat, École normale supérieure de Lyon, 1997.
[31] P. Roesch, Captures for the family $F_{a}(z)=z^{2}+a / z^{2}$, in Dynamics on the Riemann Sphere, EMS, 2006.
[32] P. Roesch, Hyperbolic components of polynomials with a fixed critical point of maximal order, Ann. Sci. École Norm. Sup. 40 (2007), 901-949.
[33] Z. Slodkowski, Holomorphic motions and polynomial hulls, Proc. Amer. Math. Soc. 111 (1991), 347-355.
[34] N. Steinmetz, On the dynamics of the McMullen family $R(z)=z^{m}+\lambda / z^{l}$, Conform. Geom. Dyn. 10 (2006), 159-183.
$4^{\mathrm{e}}$ SÉRIE - TOME 48 - 2015 - No 3
[35] L. Tan, Y. Yin, The unicritical Branner-Hubbard conjecture, in Complex dynamics, A. K. Peters, Wellesley, MA, 2009, 215-227.
[36] Y. Yin, Y. Zhai, No invariant line fields on Cantor Julia sets, Forum Math. 22 (2010), 75-94.
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