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# INFINITE CONYOLUTION OF BERNOULLI  WROBLEMS IN THE DYNAMICS OF FRACTAL GEOMETRY 

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# INFINITE CONVOLUTION OF BERNOULLI MEASURES, PV NUMBERS AND RELATED PROBLEMS IN THE DYNAMICS OF FRACTAL GEOMETRY 

## by

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#### Abstract

The case of equality between the Minkowski and Hausdorff dimension for the graphs of Weierstrass-like functions, remains largely mysterious. However significant progresses have been obtained when the graph is self-affine: for instance the analysis of the limit Rademacher function by Przytycki \& Urbański shows how this question is concerned with the arithmetics of the Pisot-Vijayaraghavan (PV) numbers and an Erdős problem about the so-called Infinite Convolution Bernoulli Measures (ICBMs). A related question is to understand the fine multifractal/dynamical structures of the special ICBM associated with the golden number and called the Erdős measure: from this point of view, we answer a question of Sidorov and Vershik about the Gibbs nature of the invariant Erdös measure.

Résumé (Problèmes de dynamique liés à la géométrie fractale de certaines convolution d'une infinité de mesures de Bernoulli en base de Pisot-Viraraghavan)

Le cas d'égalité entre la dimension de Minkowski et la dimension de Hausdorff du graphe des fonctions $\grave{a} l a$ Weierstrass, demeure toujours mystérieux. Cependant, des progrès significatifs ont été réalisés pour certains graphes autoaffines: par exemple, pour les fonctions de Rademacher, Przytycki \& Urbański, montrent comment cette question est liée à l'arithmétique des nombres de Pisot-Vijayaraghavan (PV): on retrouve alors un problème d'Erdős sur certaines convolutions d'une infinité de mesures de Bernoulli (convolutions de Bernoulli). Une question attenante est de comprendre la structure multifractale de la mesure d'Erdós i.e. la convolution de Bernoulli associée au nombre d'or: de ce point de vue, nous répondons à une question posée par Sidorov et Vershik à propos de la nature gibbsienne de la mesure d'Erdős invariante.


2000 Mathematics Subject Classification. - 28A12, 11A67, 15A48.
Key words and phrases. - $\beta$-numeration, golden number, Bernoulli convolutions, Erdős measure, Gibbs and weak Gibbs measures, fractal geometry, Hausdorff dimension.

## Introduction

Let $\alpha, \beta>1$ be two real numbers with $\alpha / \beta>1$ and note $\Gamma$ the graph of the celebrated Weierstrass function (1872), W: $[0 ; 1] \rightarrow \mathbf{R}$ such that

$$
W(x)=\sum_{n=0}^{\infty} \frac{1}{\beta^{n}} \cos \left(2 \pi \alpha^{n} x\right)
$$

The function $W$ is a classical example of a continuous nowhere differentiable map, which makes its graph $\Gamma$ an interesting object of Fractal Geometry. We fix $\alpha=2$ and include in our presentation the class of the Weierstrass-like functions $\mathcal{W}_{\beta}:[0 ; 1] \rightarrow \mathbf{R}$ such that

$$
\mathcal{W}_{\beta}(x)=\sum_{n=0}^{\infty} \frac{1}{\beta^{n}} \varphi\left(2^{n} x\right)
$$

where $1<\beta<2$ and $\varphi: \mathbf{R} \rightarrow \mathbf{R}$ is a $\mathbf{Z}$-periodic function: in what follows $\Gamma_{\beta}$ stands for the graph of $\mathcal{W}_{\beta}$. When $\varphi$ is Lipschitz continuous, $\mathcal{W}_{\beta}$ turns out to be Hőlder continuous of exponent $\log \beta / \log 2$; the Hőlder exponent is then directly related to a classical upper-bound of the Minkowski dimension (box counting dimension), that is $\operatorname{dim}_{M} \Gamma_{\beta} \leq 2-\log \beta / \log 2$. However, even if a Weierstrass-like function need not be continuous, the Minkowski dimension provides a general upper-bound of the Hausdorff dimension, so that [29, 37]:

$$
\begin{equation*}
\operatorname{dim}_{H} \Gamma_{\beta} \leq 2-\frac{\log \beta}{\log 2} \tag{1}
\end{equation*}
$$

The case of equality in (1) is a delicate problem still open for the Weierstrass functions. Among partial results for the Weierstrass-like functions, two significant progresses are concerned with self-affine graphs, the first one by Przytycki \& Urbański [37] for the limit Rademacher functions, the second one by Ledrappier [25] for the Takagi functions. In both cases the problematic equality in (1) is proved to be related to an Erdös problem [2, 5, 11, 16, 18, 43] about the family of probability measures $\nu_{\beta}(1<\beta<2)$, and called Infinite Convolution of Bernoulli Measures (ICBMs): in Section 1, we focus our attention on the limit Rademacher functions to see how the arithmetic of $\beta$ is involved and more precisely how PV-numbers yields strict inequality in (1).

The multifractal analysis of the measures $\nu_{\beta}$ is supposed to give precious informations about the fractal properties of $\nu_{\beta}$ and $\Gamma_{\beta}$; however, the task is difficult and only few results have been established. Section 2 is devoted to the statement of the multifractal formalism satisfied by the Erdős measure, that is $\nu=\nu_{\beta}$ for $\beta=(1+\sqrt{5}) / 2$ $[10,23,26]$ (we refer to $[7,9,31]$ for multifractal analysis of $\nu_{\beta}$, for other PVnumbers $\beta$ ).

Following the seminal work of Alexander \& Yorke [2] about the so-called Fat Baker's Transformation, a detailed analysis of dynamical/ergodic properties of the Erdős measure is given by Sidorov \& Vershik [42]; in particular, it is proved


Figure 1. Graph $\Gamma$ of the Weierstrass function $W$ for $\alpha=2$ and $\beta=(1+$ $\sqrt{5}) / 2$ : here $\operatorname{dim}_{M} \Gamma=2-\log \beta / \log 2$. The exact value of the Hausdorff dimension $\operatorname{dim}_{H} \Gamma$ is still unknown.


Figure 2. Graph of the Takagi function $\mathcal{W}_{\beta}$ with $\beta=(1+\sqrt{5}) / 2$ and $\varphi(x)=1-\operatorname{dist}(x, \mathbf{Z})$. By a result of Ledrappier [25], $\operatorname{dim}_{M} \Gamma_{\beta}=\operatorname{dim}_{H} \Gamma_{\beta}$ when the ICBM $\nu_{\beta}$ has Hausdorff dimension 1, which [43] arises for almost all $\beta$.
[42, Corollary 1.9] that there exists a unique probability measure $\nu^{\prime}$ invariant w.r.t. multiplication by $\beta(\bmod 1)$ and equivalent to $\nu$. In Section 3 we present the result of Sidorov \& Vershik about the measure $\nu^{\prime}$, leading to the statement of our main result, Theorem 4.2, which answers a question in [42, Remark p. 222] about the Gibbs nature of $\nu^{\prime}$; the proof of Theorem 4.2, given in Section 4, makes use of Kac Recurrence Theorem, Abramov Formula and Thermodynamic Formalism for infinite alphabet shift spaces [28, 40, 44].

## 1. Limit Rademacher functions and ICBMs

Consider for any $i \in\{\mathbf{0}, \mathbf{1}\}$, the affine contraction $A_{i}: \mathbf{R}^{2} \rightarrow \mathbf{R}^{2}$ such that $A_{i}(x, y):=A(x, y)+i V$, where $A$ and $V$ are identified to their matricial form, that is:

$$
A:=\left(\begin{array}{cc}
1 / 2 & 0 \\
0 & 1 / \beta
\end{array}\right) \quad \text { and } \quad V=\binom{1 / 2}{1 / \beta^{2}}
$$

Let $\mathbf{A}: \Omega:=\{\mathbf{0}, \mathbf{1}\}^{\mathbf{N}} \rightarrow \mathbf{R}^{2}$ be the map such that $\mathbf{A}(\xi)$ is the limit of $A_{\xi_{0} \cdots \xi_{n-1}}(0,0)$, when $n$ goes to infinity that is, $\mathbf{A}(\xi)=(\mathbf{X}(\xi), \mathbf{Y}(\xi))$, where

$$
\mathbf{X}(\xi)=\sum_{k=0}^{\infty} \xi_{k} / 2^{k+1} \quad \text { and } \quad \mathbf{Y}(\xi)=(\beta-1) \sum_{k=0}^{\infty} \xi_{k} / \beta^{k+1}
$$

It is easily seen that $\mathbf{Y}(\xi)=\mathcal{W}_{\beta}(\mathbf{X}(\xi))$, where $\mathcal{W}_{\beta}$ is the Weierstrass-like function associated to the $\mathbf{Z}$-periodic function $\varphi: x \mapsto([2 x]-2[x]) /(1-1 / \beta)$ (the coefficient $1-1 / \beta$ is introduced for notational convenience). In particular, this means that $\mathbf{A}(\Omega)$ coincides with the graph $\Gamma_{\beta}$ of $\mathcal{W}_{\beta}$ or equivalently, that $\Gamma_{\beta}$ is the self-affine set such that $\Gamma_{\beta}=A_{\mathbf{0}}\left(\Gamma_{\beta}\right) \bigcup A_{\mathbf{1}}\left(\Gamma_{\beta}\right)$. There exists a natural expanding transformation $T: \Gamma_{\beta} \rightarrow \Gamma_{\beta}$, whose inverse branches are the affine contractions $A_{0}$ and $A_{1}$ meaning that $T\left(x, \mathcal{W}_{\beta}(x)\right)=A_{i}^{-1}\left(x, \mathcal{W}_{\beta}(x)\right)$, with $i=\mathbf{0}$ if $0 \leq x \leq 1 / 2$ and $i=\mathbf{1}$ otherwise. In order to study the fractal geometry of $\Gamma_{\beta}$, it is worth to consider the good positive measure supported by $\Gamma_{\beta}$ : here we define the distribution $\mu_{\beta}$ of the random variable $\mathbf{A}: \Omega \rightarrow \mathbf{R}^{2}$ when $\Omega$ is weighted by the uniform Bernoulli measure $\mathbb{P}$. The probability measure $\mu_{\beta}$ is $T$-ergodic (with metric entropy $\mathrm{h}_{\mu_{\beta}}(T)=\log 2$ ) and according to Hutchinson Theorem [15], [6, Thm. 2.8], it is the unique probability measure satisfying the self-affine equation

$$
\begin{equation*}
\mu_{\beta}=\frac{1}{2} \mu_{\beta} \circ A_{\mathbf{0}}^{-1}+\frac{1}{2} \mu_{\beta} \circ A_{\mathbf{1}}^{-1} \tag{2}
\end{equation*}
$$

Now, consider the orthogonal projections $\pi_{x}: \mathbf{R}^{2} \rightarrow \mathbf{R} \times\{0\}$ and $\pi_{y}: \mathbf{R}^{2} \rightarrow\{0\} \times \mathbf{R}$. (In what follows $\mathbf{R} \times\{0\}$ and $\{0\} \times \mathbf{R}$ are both identified to $\mathbf{R}$.) On the one hand, $\mu_{\beta} \circ \pi_{x}^{-1}$ is the distribution of the random variable $\mathbf{X}$ on $(\Omega, \mathbb{P})$ that is, the restriction

