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# CM STABILITY AND THE GENERALIZED FUTAKI INVARIANT II

by

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Dedicated to Jean-Michel Bismut on the occasion of his 60<sup>th</sup> birthday

Abstract. — The Mabuchi K-energy map is exhibited as a singular metric on the refined CM polarization of any equivariant family  $\mathbf{X} \xrightarrow{p} S$ . Consequently we show that the generalized Futaki invariant is the leading term in the asymptotics of the reduced K-energy of the generic fiber of the map p. Properness of the K-energy implies that the generalized Futaki invariant is strictly negative.

*Résumé* (CM-stabilité et invariant de Futaki généralisé II). — On interpréte la K-énergie de Mabuchi comme une métrique singulière sur la CM-polarisation raffinée d'une famille équivariante  $\mathbf{X} \xrightarrow{p} S$ . Nous montrons que l'invariant de Futaki généralisé est le terme principal de l'asymptotique de la K-énergie réduite de la fibre générique de l'application p. Si la K-énergie est propre, alors l'invariant de Futaki généralisé est strictement négatif.

## 1. Introduction

1.1. Statement of results. — Throughout this paper X and S denote smooth, proper complex projective varieties satisfying the following conditions.

- 1.  $\mathbf{X} \subset S \times \mathbb{P}^N$ ;  $\mathbb{P}^N$  denotes the complex projective space of *lines* in  $\mathbb{C}^{N+1}$ .
- 2.  $p := p_1 : \mathbf{X} \to S$  is flat of relative dimension n, degree d with Hilbert polynomial P.
- 3.  $L|_{\mathbf{X}_z}$  is very ample and the embedding  $\mathbf{X}_z := p_1^{-1}(z) \stackrel{L}{\hookrightarrow} \mathbb{P}^N$  is given by a complete linear system for  $z \in S$ .

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4. There is an action of  $G := SL(N + 1, \mathbb{C})$  on the data compatible with the projection and the standard action on  $\mathbb{P}^N$ .

It is well known that (1) and (3) imply that

(1.1) 
$$\mathbb{P}(p_{1*}L) \cong S \times \mathbb{P}^N.$$

Which in turn is equivalent to the existence of a line bundle  $\mathscr{A}$  on S such that

(1.2) 
$$p_{1*}L \cong \underbrace{\bigoplus \mathcal{O}}_{N+1}.$$

Below Chow( $\mathbf{X}/S$ ) denotes the Chow form of the family  $\mathbf{X}/S$ ,  $\mu$  is the coefficient of  $k^{n-1}$  in P(k), and  $\mathcal{M}_n$  is the coefficient of  $\binom{m}{n}$  in the CGKM expansion of det $(p_{1*}L^{\otimes m})$  for m >> 0. A complete discussion of these notions is given in "*CM Stability and the Generalized Futaki Invariant I*". We refer the reader to that paper for the basic definitions and constructions that are used in the present article.

We define an invertible sheaf on S as follows.

### **Definition 1** (The Refined CM polarization<sup>(1)</sup>). — We have

(1.3) 
$$\mathbb{L}_1(\mathbf{X}/S) := \{ \operatorname{Chow}(\mathbf{X}/S) \otimes \mathscr{C}^{d(n+1)} \}^{n(n+1)+\mu} \otimes \mathscr{M}_n^{-2(n+1)}$$

With the family  $p_1 : \mathbf{X} \to S$  fixed throughout, we will denote  $\mathbb{L}_1(\mathbf{X}/S)$  by  $\mathbb{L}_1$  in the remainder of the paper.

Our first result exhibits the Mabuchi energy as a singular Hermitian metric on  $\mathbb{L}_1$ .

**Theorem 1.** — Let || || be any smooth Hermitian metric on  $\mathbb{L}_1^{-1}$ .<sup>(2)</sup> Then there is a continuous function  $\Psi_S : S \setminus \Delta \to (-\infty, c)$  such that for all  $z \in S/\Delta$ 

(1.4) 
$$d(n+1)\nu_{\omega|_{\mathbf{X}_z}}(\varphi_{\sigma}) = \log\left(e^{(n+1)\Psi_S(\sigma z)}\frac{||\;||^2(\sigma z)}{||\;||^2(z)}\right).$$

Here c denotes a constant which depends only on the choice of background Kähler metrics on S and X,  $\Delta$  denotes the discriminant locus of the map  $p_1$ , and  $\omega|_{\mathbf{X}_z}$ denotes the restriction of the Fubini Study form of  $\mathbb{P}^N$  to the fiber  $\mathbf{X}_z$ .

**Remark 1.** — This should be compared with the main result in Section 8 of [17]. The principal contribution of our present work is the observation that the whole theory in Section 8 of [17] should be recast from the beginning with the sheaf  $\mathbb{L}_1$ .

Let  $X \hookrightarrow \mathbb{P}^N$  be an *n* dimensional projective variety with Hilbert polynomial *P*. Let  $Hilb_m(X)$  denote the *mth* Hilbert point of *X* (see **[12]** for further information ). If  $\lambda$  is a one parameter subgroup of *G* then it is known (see **[12]**) that the weight,

<sup>&</sup>lt;sup>(1)</sup> We use this terminology in order to distinguish this sheaf from one introduced by the second author in ([17]).

<sup>&</sup>lt;sup>(2)</sup>  $\mathbb{L}_1^{-1}$  denotes the dual of  $\mathbb{L}_1$ .

 $w_{\lambda}(m)$ , of  $Hilb_m(X)$  with respect to  $\lambda$  is a *polynomial* in m of degree at most n + 1. That is,

$$w_{\lambda}(m) = a_{n+1}(\lambda)m^{n+1} + a_n(\lambda)m^n + \cdots$$

Then the ratio may be expanded as follows.

$$\frac{w_{\lambda}(m)}{mP(m)} = F_0(\lambda) + F_1(\lambda)\frac{1}{m} + \dots + F_l(\lambda)\frac{1}{m^l} + \dots$$

**Definition 2** (Donaldson ([5])). —  $F_1(\lambda)$  is the generalized Futaki invariant of X with respect to  $\lambda$ .

In our previous paper we have shown the following.

*Theorem* (The weight of the Refined CM polarization). — i) There is a natural G linearization on the line bundle  $\mathbb{L}_1$ .

ii) Let  $\lambda$  be a one parameter subgroup of G. Let  $z \in \mathfrak{Hill}_{\mathbb{P}^N}^{P}(\mathbb{C})$ . Let  $w_{\lambda}(z)$  denote the weight of the restricted  $\mathbb{C}^*$  action (whose existence is asserted in i)) on  $\mathbb{L}_1^{-1}|_{z_0}$  where  $z_0 = \lambda(0)z$ . Then

(1.5) 
$$w_{\lambda}(z) = F_1(\lambda).$$

The main result of the paper is the following corollary of (1.4) and (1.5).

Corollary 1 (Algebraic asymptotics of the Mabuchi energy). — Let  $\varphi_{\lambda(t)}$  be the Bergman potential associated to an algebraic 1psg  $\lambda$  of G, and let  $z \in S \setminus \Delta$ . Then there is an asymptotic expansion

(1.6) 
$$d(n+1)\nu_{\omega|_{\mathbf{x}_{z}}}(\varphi_{\lambda(t)}) - \Psi_{S}(\lambda(t)) = F_{1}(\lambda)\log(|t|^{2}) + O(1) \text{ as } |t| \to 0.$$

Moreover  $\Psi_S(\lambda(t)) = \psi(\lambda) \log(|t|^2) + O(1)$  where  $\psi(\lambda) \in \mathbb{Q}_{\geq 0}$ . Moreover,  $\psi(\lambda) \in \mathbb{Q}_+$ if and only if  $\lambda(0)\mathbf{X}_z = \mathbf{X}_{\lambda(0)z}$  (the limit cycle<sup>(3)</sup> of  $\mathbf{X}_z$  under  $\lambda$ ) has a component of multiplicity greater than one. Here O(1) denotes any quantity which is bounded as  $|t| \to 0$ .

Moser iteration and a refined Sobolev inequality (see [11] and [7]) yield the following.

**Corollary 2.** — If  $\nu_{\omega|\mathbf{x}_z}$  is proper (bounded from below) then the generalized Futaki invariant of  $\mathbf{X}_z$  is strictly negative (nonnegative) for all  $\lambda \in G$ .

**Remark 2.** We call the left hand side of (1.6) the reduced K-Energy along  $\lambda$ . We also point out that while it is certainly the case that  $F_1(\lambda)$  may be defined for any subscheme of  $\mathbb{P}^N$  it evidently only controls the behavior of the K-Energy when  $\lambda(0)\mathbf{X}_z$  is reduced.

<sup>&</sup>lt;sup>(3)</sup> See [12] pg. 61.

**Remark 3.** — The precise constant d(n+1) in front of  $\nu_{\omega}$  is not really crucial, since what really matters is the sign of  $F_1(\lambda) + \psi(\lambda)$ . That  $\Psi_S(\lambda(t))$  has logarithmic singularities can be deduced from [13].

**Remark 4**. — We emphasize that we do not assume the limit cycle is smooth.

#### 2. Background and Motivation

Let  $(X, \omega)$  be a compact Kähler manifold ( $\omega$  not necessarily a Hodge class) and  $P(X, \omega) := \{\varphi \in C^{\infty}(X) : \omega_{\varphi} := \omega + \frac{\sqrt{-1}}{2\pi} \partial \overline{\partial} \varphi > 0\}$  the space of Kähler potentials. This is the usual description of all Kähler metrics in the same class as  $\omega$  (up to translations by constants). It is not an overstatement to say that the most basic problem in Kähler geometry is the following

Does there exist  $\varphi \in P(X, \omega)$  such that  $\operatorname{Scal}(\omega_{\varphi}) \equiv \mu$ ? (\*)

This is a fully nonlinear *fourth order* elliptic partial differential equation for  $\varphi$ .  $\mu$  is a constant, the average of the scalar curvature, it depends only on  $c_1(X)$  and  $[\omega]$ . When  $c_1(X) > 0$  and  $\omega$  represents the *anticanonical* class a simple application of the Hodge Theory shows that (\*) is equivalent to the *Monge-Ampere equation*.

$$\frac{\det(g_{i\overline{j}} + \varphi_{i\overline{j}})}{\det(g_{i\overline{j}})} = e^{F - \kappa \varphi} \quad (\kappa = 1) \qquad (**)$$

where F denotes the Ricci potential. When  $\kappa = 0$  this is the celebrated Calabi problem solved by S.T.Yau and when  $\kappa < 0$  this was solved by Aubin and Yau independently in the 70's. It is well known that (\*) is actually a *variational* problem. There is a natural energy on the space  $P(X, \omega)$  whose critical points are those  $\varphi$  such that  $\omega_{\varphi}$ has constant scalar curvature (csc). This energy was introduced by T. Mabuchi ([10]) in the 1980's. It is called the *K*-Energy map (denoted by  $\nu_{\omega}$ ) and is given by the following formula

$$\nu_{\omega}(\varphi) := -\frac{1}{V} \int_{0}^{1} \int_{X} \dot{\varphi_{t}}(\operatorname{Scal}(\varphi_{t}) - \mu) \omega_{t}^{n} dt$$

Above,  $\varphi_t$  is a smooth path in  $P(X, \omega)$  joining 0 with  $\varphi$ . The K-Energy does not depend on the path chosen. In fact there is the following well known formula for  $\nu_{\omega}$  where O(1) denotes a quantity which is bounded on  $P(X, \omega)$ .

$$\nu_{\omega}(\varphi) = \int_{X} \log\left(\frac{\omega_{\varphi}^{n}}{\omega^{n}}\right) \frac{\omega_{\varphi}^{n}}{V} - \mu(I_{\omega}(\varphi) - J_{\omega}(\varphi)) + O(1)$$
$$J_{\omega}(\varphi) := \frac{1}{V} \int_{X} \sum_{i=0}^{n-1} \frac{\sqrt{-1}}{2\pi} \frac{i+1}{n+1} \partial \varphi \wedge \overline{\partial} \varphi \wedge \omega^{i} \wedge \omega_{\varphi}^{n-i-1}$$
$$I_{\omega}(\varphi) := \frac{1}{V} \int_{X} \varphi(\omega^{n} - \omega_{\varphi}^{n}).$$

We have written down the K-energy in the case when  $\omega = c_1(X)$ . Observe that  $\nu_{\omega}$  is essentially the *difference* of two positive terms. What is of interest for us is that

the problem (\*) is not only a variational problem but a *minimization* problem. With this said we have the following fundamental result.

**Theorem (S. Bando and T. Mabuchi [1]).** If  $\omega = c_1(X)$  admits a Kähler Einstein metric then  $\nu_{\omega} \geq 0$ . The absolute minimum is taken on the solution to (\*\*) (which is unique up to automorphisms of X).

Therefore a *necessary* condition for the existence of a Kähler Einstein metric is a bound from below on  $\nu_{\omega}$ . In order to get a *sufficient* condition one requires that the K-energy *grow* at a certain rate. Precisely, it is required that the K-Energy be *proper*. This concept was introduced by the second author in [17].

**Definition 3.** —  $\nu_{\omega}$  is proper if there exists a strictly increasing function  $f : \mathbb{R}_+ \longrightarrow \mathbb{R}_+$  (where  $\lim_{T \longrightarrow \infty} f(T) = \infty$ ) such that  $\nu_{\omega}(\varphi) \ge f(J_{\omega}(\varphi))$  for all  $\varphi \in P(M, \omega)$ .

**Theorem ([17]).** — Assume that Aut(X) is discrete. Then  $\omega = c_1(X)$  admits a Kähler Einstein metric if and only if  $\nu_{\omega}$  is proper.

The next result was established by the second author and Xiuxiong Chen. It holds in an *arbitrary* Kähler class  $\omega$ . An alternative proof of this was given by Donaldson for polarized projective manifolds.

## **Theorem ([3]).** — If $\omega$ admits a metric of csc then $\nu_{\omega} \geq 0$ .

In this paper our interest is to test for a lower bound of  $\nu_{\omega}$  along the large but finite dimensional group G of *matrices* in the polarized case. When we restrict our attention to G we make the connection with Mumfords' Geometric Invariant Theory. The past couple of years have witnessed quite a bit of activity on this problem due to this connection.

To put things in historical perspective consider the various formulations of the Futaki invariant.

i) 1983 Futaki ([6]) introduces his invariant as a lie algebra character on a Fano manifold X

$$F_{\omega}: \eta(X) \longrightarrow \mathbb{C}.$$

ii) 1986 Mabuchi (see [10]) integrates the Futaki invariant with the introduction of the K-energy map. The linearization of the K-energy along orbits of holomorphic vector fields is the real part of the Futaki invariant.

iii) 1992 Ding and Tian ([4]) introduced the generalized Futaki invariant. Here the jumping of complex structures is introduced. The limit of the derivative of the K-Energy map is identified with the generalized Futaki invariant of  $X^{\lambda(0)}$  provided this limit has at most normal singularities.

iv) 1997 The CM polarization is defined (see [17]) for *smooth* families, as the relative canonical bundle is explicitly involved in the definition. K-Stability is defined in terms of special degenerations and the generalized Futaki invariant.