Astérisque

B. HELFFER

J. Sjöstrand

Semiclassical expansions of the thermodynamic limit for a Schrödinger equation

Astérisque, tome 210 (1992), p. 135-181

http://www.numdam.org/item?id=AST_1992_210_135_0

© Société mathématique de France, 1992, tous droits réservés.

L'accès aux archives de la collection « Astérisque » (http://smf4.emath.fr/ Publications/Asterisque/) implique l'accord avec les conditions générales d'utilisation (http://www.numdam.org/conditions). Toute utilisation commerciale ou impression systématique est constitutive d'une infraction pénale. Toute copie ou impression de ce fichier doit contenir la présente mention de copyright.

\mathcal{N} umdam

Article numérisé dans le cadre du programme Numérisation de documents anciens mathématiques http://www.numdam.org/

Semiclassical expansions of the thermodynamic limit for a Schrödinger equation

I. The one well case

by B.Helffer and J.Sjöstrand

<u>§1 Presentation of the problem :</u>

One of the motivations of the study presented here is a statistical model introduced by M.Kac $[Ka]_2$ and called the exponential bidimensional model. This model was supposed to present phase transition. Let us just recall here (see $[Ka]_2$ or [Br-He] for details) that after some reductions M.Kac arrive to the question of studying the spectral properties of the following operator:

 $(1.1) K_{m}(h) :=$

$$= \exp\left[-V^{(m)}(x)/2\right] \cdot \exp\left[h^2 \sum_{k=1}^{m} \partial^2 / \partial x_k^2\right] \cdot \exp\left[-V^{(m)}(x)/2\right]$$

with¹:

(1.2)
$$V^{(m)}(x) = (1/4) \sum_{k=1}^{m} x_k^2 - \sum_{k=1}^{m} \log ch(\sqrt{\nu/2} (x_k + x_{k+1})).$$

^{&#}x27; In fact, the operator which appears in Kac is $\exp(-mh/2)K_m(h)$. It is easier w.l.o.g. in this article to work with this modified Kac operator.

The parameter v is here the inverse of the temperature and h is a semi-classical parameter. The two questions of interest are in this context: (1.3) If $\mu_1(m;h,v)$ is the largest eigenvalue of the Kac's operator, what is the behavior as a function of v and h of the thermodynamic quantity :

 $\operatorname{Lim}_{m\to\infty} (-\operatorname{Log} \mu_1(m;h,v)/m).$

(1.4) If $\mu_2(m;h,v)$ is the second eigenvalue (which is $< \mu_1(m;h,v)$ by standard results), <u>can we study</u> the quantity :

 $\operatorname{Lim}_{m\to\infty}(\mu_2(m;h,v)/\mu_1(m;h,v)).$

From discussions with specialists in statistical mechanics (with T.Spencer for example), we get the impression that this problem is probably well understood and that according to the value of v with respect to a critical value v_c the answer to (1.4) will be that the limit will be <1 for $v < v_c$ and will be 1 for $v > v_c$. This is a sign of a transition of phase. However, we do not have a precise reference for that and at least the problem of analyzing in detail the behavior of the different thermodynamic quantities near the critical value v_c seems to remain open.

In his interesting course in Brandeis $[Ka]_2$, M. Kac explains, at least heuristically, how to compare (in the semi-classical context) the operator K_m (h) to the exponential of (minus) a Schrödinger operator. The validity of this approximation (for m fixed) has been studied more carefully in [He-Br] and [He] using some results of [He-Sj]_{1.4}.

If we admit this approximation, we shall find the following problems for the Schrödinger equation :

(1.5)
$$P_{m}(h) = -\sum_{k=1}^{m} h^{2} \partial^{2} / \partial x_{k}^{2} + V^{(m)}(x).$$

(1.6) If $\lambda_1^{(m)}(h,v)$ is the smallest eigenvalue of the Schrödinger's operator, study as a function of v and h the thermodynamic quantity :

 $\operatorname{Lim}_{m\to\infty} (\lambda_1(m;h,v)/m).$

(1.7) If $\lambda_2(m;h,v)$ is the second eigenvalue (which is $>\lambda_1(m;h,v)$ by standard results), study the quantity :

 $\operatorname{Lim}_{m\to\infty}(\lambda_2(m;h,v) - \lambda_1(m;h,v)).$

Forgetting the initial Kac's problem, we shall start to study in this article these two questions (1.6) and (1.7). Because it is a high dimension problem, we shall use (at least in the semi-classical context) the techniques introduced by one of us (J.S). Most of the results which are given here :

(1) existence of the thermodynamic limit $\lim_{m \to \infty} (\lambda_1(m;h,v)/m)$

(2) asymptotic expansion of the limit as a formal series in h

(3) rapidity of the convergence as $m \to \infty$

are given in a relatively general framework but we shall see how it can be applied in our motivating example, in the particular case where $v < v_c$.

This is of course just the starting point (and the easiest) of a study which has to consider after the case where $v > v_c$, and then the transition around $v = v_c$. There is some hope to return later to the initial Kac's problem. This v_c can be guessed by looking carefully to the properties of $V^{(m)}$. As observed by V.Kac, for v < 1/4, the potential $V^{(m)}$ has a unique minimum at 0 and appears to be convex. For v > 1/4, we shall observe a double well problem which is certainly more difficult to analyze.

The principal result of this paper will be :

Theorem 1.1

If v < 1/4, the limit $\Lambda(h,v) = \lim_{m \to \infty} (\lambda_1(m;h,v)/m)$ exists and admit a complete asymptotic expansion : $\Lambda(h,v) \sim h\Sigma_{j \ge 0} \Lambda_j(v) h^j$ as h tends to 0. Moreover, if we denote the corresponding semiclassical expansions for $\lambda_1(m;h,v)/m$ by : $(\lambda_1(m;h,v)/m) \sim h\Sigma_{j \ge 0} \Lambda_j(m,v) h^j$, there exists k_0 s.t. for each j, there exists a constant $C_j(v)$, s.t. $|\Lambda_j(v) - \Lambda_j(m,v)| \le C_j(v)$. $exp(-k_0 m)$. $C_j(v)$ can be chosen independently of v in a compact of [0, 1/4[.

The problems, we consider here, are also connected to quantum field theory problems and a lot of results have been obtained by other techniques (see for example the new edition of [Gl-Ja] for a updated presentation).

The paper is organized in three parts.

The first part (§ 2 and §3) is essentially devoted to the proof of the existence of the thermodynamic limit. This is a non-semiclassical proof but we shall see that a control of the convergence with respect to parameters can be useful. In §3 we give additional remarks (to $[Sj]_2$) on universal estimates of the splitting of the two first eigenvalues.

The second part (§4 and §5) is the semi-classical part and the natural continuation of two papers by one of us $(J.S) [Sj]_{1,2}$.