

*quatrième série - tome 57    fascicule 5    septembre-octobre 2024*

*ANNALES  
SCIENTIFIQUES  
de  
L'ÉCOLE  
NORMALE  
SUPÉRIEURE*

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*Ising model and s-embeddings of planar graphs*

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SOCIÉTÉ MATHÉMATIQUE DE FRANCE

# Annales Scientifiques de l'École Normale Supérieure

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Publiées avec le concours du Centre National de la Recherche Scientifique

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YVES DE CORNULIER

Publication fondée en 1864 par Louis Pasteur

Continuée de 1872 à 1882 par H. SAINTE-CLAIRE DEVILLE  
de 1883 à 1888 par H. DEBRAY  
de 1889 à 1900 par C. HERMITE  
de 1901 à 1917 par G. DARBOUX  
de 1918 à 1941 par É. PICARD  
de 1942 à 1967 par P. MONTEL

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Édition et abonnements / *Publication and subscriptions*

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13288 Marseille Cedex 09  
Tél. : (33) 04 91 26 74 64  
Email : [abonnements@smf.emath.fr](mailto:abonnements@smf.emath.fr)

Tarifs

Abonnement électronique : 480 euros.  
Abonnement avec supplément papier :  
Europe : 675 €. Hors Europe : 759 € (\$ 985). Vente au numéro : 77 €.

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ISSN 0012-9593 (print) 1873-2151 (electronic)

Directrice de la publication : Isabelle Gallagher  
Périodicité : 6 n<sup>os</sup> / an

# ISING MODEL AND S-EMBEDDINGS OF PLANAR GRAPHS

BY DMITRY CHELKAK

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**ABSTRACT.** – We discuss the notion of s-embeddings  $\mathcal{S} = \mathcal{S}_{\mathcal{X}}$  of planar graphs carrying a nearest-neighbor Ising model. The construction of  $\mathcal{S}_{\mathcal{X}}$  is based upon a choice of a global complex-valued solution  $\mathcal{X}$  of the propagation equation for Kadanoff-Ceva fermions. Each choice of  $\mathcal{X}$  provides an interpretation of all other fermionic observables as s-holomorphic functions on  $\mathcal{S}_{\mathcal{X}}$ . We set up a general framework for the analysis of such functions on s-embeddings  $\mathcal{S}^{\delta}$  with  $\delta \rightarrow 0$ . Throughout this analysis, a key role is played by the functions  $\mathcal{Q}^{\delta}$  associated with  $\mathcal{S}^{\delta}$ , the so-called origami maps in the bipartite dimer model terminology. In particular, we give an interpretation of the mean curvature of the limit of discrete surfaces  $(\mathcal{S}^{\delta}; \mathcal{Q}^{\delta})$  viewed in the Minkowski space  $\mathbb{R}^{2,1}$  as the mass in the Dirac equation describing the continuous limit of the model.

We then focus on the simplest situation when  $\mathcal{S}^{\delta}$  have uniformly bounded lengths/angles and  $\mathcal{Q}^{\delta} = O(\delta)$ ; as a particular case this includes all critical Ising models on doubly periodic graphs via their canonical s-embeddings. In this setup we prove RSW-type crossing estimates for the random cluster representation of the model and the convergence of basic fermionic observables. The proof relies upon a new strategy and also provides a quantitative estimate on the speed of convergence.

**RÉSUMÉ.** – Nous introduisons la notion de s-plongements  $\mathcal{S} = \mathcal{S}_{\mathcal{X}}$  de graphes planaires munis d'un modèle d'Ising. La construction de  $\mathcal{S}_{\mathcal{X}}$  se base sur un choix d'une solution  $\mathcal{X}$  à valeurs complexes de l'équation de propagation pour les fermions de Kadanoff-Ceva. Chaque choix de  $\mathcal{X}$  donne une interprétation de toutes les autres observables fermioniques comme fonctions s-holomorphes sur  $\mathcal{S}_{\mathcal{X}}$ . Nous établissons un cadre général pour l'analyse de telles fonctions sur les s-plongements  $\mathcal{S}^{\delta}$  avec  $\delta \rightarrow 0$ . Tout au long de cette analyse, un rôle essentiel est joué par les fonctions  $\mathcal{Q}^{\delta}$  associées aux  $\mathcal{S}^{\delta}$ , dites applications origami dans la terminologie du modèle de dimères sur des graphes bipartis. En particulier nous donnons une interprétation de la courbure moyenne de la limite des surfaces discrètes  $(\mathcal{S}^{\delta}; \mathcal{Q}^{\delta})$  vue dans l'espace de Minkowski  $\mathbb{R}^{2,1}$  comme la masse dans l'équation de Dirac décrivant la limite continue du modèle.

Par la suite nous nous concentrons sur la situation la plus simple où les  $\mathcal{S}^{\delta}$  ont des longueurs/angles bornés uniformément et  $\mathcal{Q}^{\delta} = O(\delta)$ ; cela comprend comme cas particulier tous les modèles d'Ising critiques sur des graphes doublement périodiques via leurs s-plongements canoniques. Dans ce cadre, nous démontrons des estimations de croisement de type RSW pour la représentation de Fortuin-Kasteleyn du modèle et la convergence des observables fermioniques les plus simples. La démonstration se base sur une nouvelle stratégie et fournit une estimation quantitative de la vitesse de convergence.

## 1. Introduction, main results and perspectives

### 1.1. General context

The Ising model of a ferromagnet, introduced by Lenz in 1920, recently celebrated its centenary; probably, this is one of the most studied models in statistical mechanics. The *planar* Ising model (i.e., the 2D model with nearest-neighbor interactions) gives rise to surprisingly rich structures of correlation functions; we refer the reader to monographs [26, 44, 46, 48], lecture notes [24], as well as to the introductions of the papers [1, 18] and references therein for more background on the subject, discussed from a variety of perspectives.

In this paper we consider the planar Ising model without the magnetic field, which is known to be exactly solvable on any graph and at any temperature: the partition function can be written as the Pfaffian of a certain matrix and the entries of the inverse matrix—known under the name *fermionic observables*—satisfy a simple *propagation equation*; we refer the reader to the paper [9] for more details on various combinatorial formalisms used to study the planar Ising model during its long history.

We prefer to work with the ferromagnetic Ising model defined on *faces* of a planar graph  $G$ ; the partition function is given by

$$(1.1) \quad \mathcal{Z}(G) = \sum_{\sigma: G^\circ \rightarrow \{\pm 1\}} \exp \left[ \beta \sum_{e \in E(G)} J_e \sigma_{v_-^\circ(e)} \sigma_{v_+^\circ(e)} \right],$$

where  $G^\circ$  and  $E(G)$  denote the dual graph and the set of edges of  $G$ , respectively;  $\beta = 1/kT$  is the inverse temperature,  $J_e > 0$  are interaction constants assigned to the edges of  $G$ , and  $v_\pm^\circ(e) \in G^\circ$  denote the two faces adjacent to  $e \in E(G)$ . Passing to the *domain walls* representation (also known as the low-temperature expansion) of the model, one can rewrite  $\mathcal{Z}(G)$  as

$$\mathcal{Z}(G) = 2 \prod_{e \in E(G)} (x(e))^{-1/2} \times \sum_{C \in \mathcal{E}(G)} \prod_{e \in C} x(e), \quad x(e) := \exp[-2\beta J_e].$$

where  $\mathcal{E}(G)$  denotes the set of all even subgraphs of  $G$ . Throughout this paper we identify edges of  $G$  with faces  $z(e)$  of the graph  $\Lambda(G) := G \cup G^\circ$  and denote

$$(1.2) \quad \theta_{z(e)} := 2 \arctan x(e) \in (0, \frac{1}{2}\pi),$$

Let us emphasize that we do *not* fix an embedding of  $G$  into  $\mathbb{C}$  at this point, thus  $\theta_e$  is nothing more than another (abstract, i.e., not geometric) *parametrization* of the interaction constants  $J_e$ .

We are mostly interested in the situation when  $(G, x)$  is a big weighted graph carrying critical or near-critical Ising weights  $x(e)$  though we do *not* precise the exact sense of this condition and mention it here only to give a proper perspective. For instance, the reader can think about the setup in which  $G$  is a subgraph of an infinite doubly periodic graph, in this case the criticality condition on the collection of weights  $x(e)$  is known explicitly [41, 19]. However, this periodic setup is *not* the main motivation for our paper and should be viewed as a very particular case though our results are new even in this situation and, in particular, answer a question posed by Duminil-Copin and Smirnov in their lecture notes on the conformal invariance of 2D lattice models; see [24, Question 8.5].

Given a weighted graph  $(G, x)$  we aim to embed it into the complex plane  $\mathbb{C}$  (actually, we construct both an embedding of  $\Lambda(G) := G \cup G^\circ$  and of its dual graph  $\diamond(G)$ ) in a

way allowing to analyze (subsequential) limits of fermionic observables in the same spirit as in the seminal work of Smirnov [53, 54] on the critical Ising model on  $\mathbb{Z}^2$ . A possible analogy could be Tutte's barycentric embeddings which, among other things, provide a framework to study the convergence of discrete harmonic functions to continuous ones. Let us emphasize that *s-embeddings* introduced in our paper are *not* directly related to these barycentric embeddings; in fact both can be viewed as special cases of a more general construction called *t-embeddings* or *Coloumb gauges* appearing in the bipartite dimer model context [39, 15]. The notion of s-embeddings was first announced in [8] and motivated a closely related research project [15, 16]. We benefit from this interplay and partly rely upon results obtained in [15] in a more general context.

One of the most important (from our perspective) motivations to study special embeddings of irregular weighted planar graphs into  $\mathbb{C}$  is the conjectural convergence of critical random maps carrying a lattice model to an object called *Liouville quantum gravity*; e.g., see [25, 27] or [51] and references therein. Though this discussion goes far beyond the scope of this paper, let us briefly mention several ideas of this kind that appeared in the literature: circle packing embeddings (see [47] and references therein), embeddings as piecewise flat Riemann surfaces (e.g., see [29] or [27, Section 1.7.1]), embeddings by square tilings for UST-decorated random maps, Tutte's embeddings [32, 31] and circle packings [30] for the mating-of-trees approach to the LQG, the so-called Cardy embeddings [33] introduced recently in the pure gravity context, etc. We believe that adding s-embeddings to this list might be a proper path for the analysis of random maps carrying the Ising model.

Besides generalizing the results of [54, 18] to a large family of deterministic graphs with 'flat' functions  $\mathcal{Q}^\delta = O(\delta)$  and suggesting a general framework to attack problems coming from the random maps context, this paper also contains the following contributions to the subject:

- A new 'constructive' strategy of the proof of Theorem 1.2, which avoids compactness arguments in the Carathéodory topology on the set of planar domains and gives a quantitative estimate of the speed of convergence; see Section 4.3 for details. Due to a general framework recently developed in [52] this also allows to control the speed of convergence of FK-Ising interfaces to SLE(16/3) curves; see [52] for more details.
- In the non-flat case  $\mathcal{Q}^\delta \not\rightarrow 0$  as  $\delta \rightarrow 0$ , we reveal the importance of embeddings of planar graphs carrying a nearest-neighbor Ising model into the *Minkowski space*  $\mathbb{R}^{2,1}$  and not simply into the complex plane. As briefly discussed in Section 2.7, this leads to a rather unexpected interpretation of the *mass* appearing in the effective description of the model in the small mesh size limit as the *mean curvature* of the corresponding surface in  $\mathbb{R}^{2,1}$ .

The latter observation can be used in deterministic setups (e.g., see a recent paper [14] where such an interpretation of a near-critical model on isoradial grids is discussed) and leads to interesting questions on limits of random surfaces in  $\mathbb{R}^{2,1}$  obtained via s-embeddings of appropriate random maps carrying the Ising model.