

SYMMETRIES OF THE NONLINEAR SCHRÖDINGER EQUATION

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ABSTRACT. — Symmetries of the defocusing nonlinear Schrödinger equation are expressed in action-angle coordinates and characterized in terms of the periodic and Dirichlet spectrum of the associated Zakharov-Shabat system. Application: proof of the conjecture that the periodic spectrum $\dots < \lambda_k^- \leq \lambda_k^+ < \lambda_{k+1}^- \leq \dots$ of a Zakharov-Shabat operator is symmetric, *i.e.* $\lambda_k^\pm = -\lambda_{-k}^\mp$ for all k , if and only if the sequence $(\gamma_k)_{k \in \mathbb{Z}}$ of gap lengths, $\gamma_k := \lambda_k^+ - \lambda_k^-$, is symmetric with respect to $k = 0$.

RÉSUMÉ (*Symétries de l'équation de Schrödinger non linéaire*). — Les symétries de l'équation de Schrödinger non linéaire sont exprimées dans les variables action-angles et caractérisées à l'aide du spectre périodique et du spectre de Dirichlet du système de Zakharov-Shabat associé. Comme application, nous démontrons la conjecture suivante : le spectre périodique $\dots < \lambda_k^- \leq \lambda_k^+ < \lambda_{k+1}^- \leq \dots$ de l'opérateur de Zakharov-Shabat est symétrique, *i.e.* $\lambda_k^\pm = -\lambda_{-k}^\mp$ pour tout k , si et seulement si la suite $(\gamma_k)_{k \in \mathbb{Z}}$ des longueurs des intervalles d'instabilité, $\gamma_k := \lambda_k^+ - \lambda_k^-$, est symétrique par rapport à $k = 0$.

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1. Introduction

The defocusing nonlinear Schrödinger equation NLS on the circle

$$(1.1) \quad i \partial_t \varphi = -\partial_x^2 \varphi + 2|\varphi|^2 \varphi$$

can be viewed as a completely integrable Hamiltonian system of infinite dimension. Indeed, on the phase $L^2(S^1; \mathbb{C})$, introduce the Poisson bracket

$$\{F, G\} := i \int_{S^1} \left(\frac{\partial F}{\partial \varphi(x)} \cdot \frac{\partial G}{\partial \bar{\varphi}(x)} - \frac{\partial F}{\partial \bar{\varphi}(x)} \cdot \frac{\partial G}{\partial \varphi(x)} \right) dx.$$

Equation (1.1) can then be written in Hamiltonian form as follows

$$\frac{\partial \varphi}{\partial t} = \{\mathcal{H}, \varphi\} = -i \frac{\partial \mathcal{H}}{\partial \bar{\varphi}}, \quad \frac{\partial \bar{\varphi}}{\partial t} = \{\mathcal{H}, \bar{\varphi}\} = i \frac{\partial \mathcal{H}}{\partial \varphi},$$

where the Hamiltonian $\mathcal{H}$ is given by (cf. [2])

$$\mathcal{H}(\varphi) := \int_{S^1} \left(\left| \frac{\partial \varphi}{\partial x} \right|^2 + |\varphi|^4 \right) dx.$$

Consider the following symmetry operators, acting on $L^2(S^1; \mathbb{C})$,

$$(1.2) \quad \mathcal{S}_1(\varphi) := \bar{\varphi}, \quad \mathcal{S}_2(\varphi) = \check{\varphi},$$

$$(1.3) \quad M_\alpha \varphi := e^{i\alpha} \varphi, \quad T_\tau \varphi := \varphi(\tau + \cdot),$$

where $\check{\varphi}$ is defined by $\check{\varphi}(x) = \varphi(-x)$. For convenience, we introduce $\mathcal{S}_3 := M_\pi$, i.e. $\mathcal{S}_3(\varphi) = -\varphi$. Notice that the Hamiltonian $\mathcal{H}$ is invariant under $\mathcal{S}_1, \mathcal{S}_2, M_\alpha$ and T_τ .

Denote by $U(t)$ the solution operator of (1.1) for initial data in $L^2(S^1; \mathbb{C})$ (or some Sobolev space $H^N(S^1; \mathbb{C})$) (cf [1]). It is immediate that $U(t)$ commutes with $\mathcal{S}_2, \mathcal{S}_3, M_\alpha$ and T_τ and that

$$(1.4) \quad U(t)\mathcal{S}_1 = \mathcal{S}_1 U(-t).$$

Recall that NLS admits a Lax pair representation

$$\frac{dL}{dt} = [L, A]$$

where $L = L(\varphi)$ is the Zakharov-Shabat operator

$$(1.5) \quad L(\varphi) := i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \frac{d}{dx} + \begin{pmatrix} 0 & \varphi \\ \bar{\varphi} & 0 \end{pmatrix}$$

and A is a (rather complicated) operator given in [2]. We remark that $L(\varphi)$ is unitarily equivalent to the well known AKNS-operator

$$(1.6) \quad H(\varphi) := \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \frac{d}{dx} + \begin{pmatrix} -q & p \\ p & q \end{pmatrix}$$

where $\varphi = -q + ip$, a fact which will be used throughout the paper.

Denote by $\text{spec}_{\text{per}} L(\varphi)$ the periodic spectrum of $L(\varphi)$ when considered on the interval $[0, 2]$ and by $\text{spec}_{\text{Dir}}^{\pm} L(\varphi)$ the Dirichlet spectra of $L(\varphi)$ when considered on the interval $[0, 1]$ (cf. Definitions (2.5) and (2.6) below). The operator $L(\varphi)$ is selfadjoint when considered with periodic or Dirichlet boundary conditions. Hence both $\text{spec}_{\text{per}} L(\varphi)$ and $\text{spec}_{\text{Dir}} L(\varphi)$ are real.

By elementary considerations one shows that

$$\begin{aligned}\text{spec}_{\text{per}} L(\bar{\varphi}) &= -\text{spec}_{\text{per}} L(\varphi), & \text{spec}_{\text{per}} L(\check{\varphi}) &= -\text{spec}_{\text{per}} L(\varphi), \\ \text{spec}_{\text{per}} L(M_{\alpha}\varphi) &= \text{spec}_{\text{per}} L(\varphi), & \text{spec}_{\text{per}} L(T_{\tau}\varphi) &= \text{spec}_{\text{per}} L(\varphi)\end{aligned}$$

and expresses $\text{spec}_{\text{Dir}}^{+} L(\mathcal{S}_j\varphi)$ for $j = 1, 2, 3$ in terms of $\text{spec}_{\text{Dir}}^{-} L(\varphi)$.

Recall from [7] (see also [8]) that NLS admits global Birkhoff coordinates. Denote by $\ell^2(\mathbb{Z}; \mathbb{R}^2)$ the space of ℓ^2 -sequences $(x_j, y_j)_{j \in \mathbb{Z}}$ endowed with the canonical Poisson bracket $\{x_i, x_j\} = 0$, $\{y_i, y_j\} = 0$ and $\{x_i, y_j\} = \delta_{ij}$.

THEOREM 1.1. — *There exists a canonical diffeomorphism Φ*

$$\Phi : \ell^2(\mathbb{Z}; \mathbb{R}^2) \longrightarrow L^2(S^1; \mathbb{C})$$

such that

- 1) Φ is bianalytic;
- 2) the restriction of Φ to the weighted ℓ^2 -space $\ell_N^2(\mathbb{Z}; \mathbb{R}^2)$ ($N \geq 1$) is a diffeomorphism onto the Sobolev space $H^N(S^1; \mathbb{C})$;
- 3) $(x_j, y_j)_{j \in \mathbb{Z}} = \Phi^{-1}(\phi)$ are Birkhoff coordinates for NLS and its hierarchy, i.e. any Hamiltonian in the hierarchy is a function of the actions $I_j := \frac{1}{2}(x_j^2 + y_j^2)$ only.

In this article we use the explicit formulas for action and angle variables given in [8] (see also [7]) to obtain

THEOREM 1.2. — (i) *The actions are left invariant by M_{α} and T_{τ} , i.e. for any $k \in \mathbb{Z}$*

$$I_k(M_{\alpha}\varphi) = I_k(\varphi) \quad \text{and} \quad I_k(T_{\tau}\varphi) = I_k(\varphi)$$

whereas $I_k(\bar{\varphi})$ and $I_k(\check{\varphi})$ can be computed to be ($j = 1, 2$)

$$I_k(\mathcal{S}_j\varphi) = I_{-k}(\varphi).$$

(ii) *For k with $I_k \neq 0$*

$$\begin{aligned}\theta_k(M_{\alpha}\varphi) &\equiv \theta_k + \alpha \pmod{2\pi}, \\ \theta_k(\check{\varphi}) &\equiv \theta_{-k}(\varphi) \pmod{2\pi}, \\ \theta_k(\bar{\varphi}) &\equiv -\theta_{-k}(\varphi) \pmod{2\pi}.\end{aligned}$$

As a first application of Theorem 1.2 one obtains (cf. Proposition 4.1 in Section 4)

COROLLARY 1.1. — *When evaluated at $I = (I_k)_{k \in \mathbb{Z}}$ with $I_k = I_{-k}$ for all $k \in \mathbb{Z}$, the NLS frequencies $\omega = (\omega_k)_{k \in \mathbb{Z}}$, $\omega_k = \partial \mathcal{H} / \partial I_k$, are symmetric, i.e. $\omega_k(I) = \omega_{-k}(I)$ for all $k \in \mathbb{Z}$.*

The main motivation for proving Theorem 1.2 and Corollary 1.1 comes from an application to a KAM type theorem established in [5] (see also [6]).

As a second application, Theorem 1.2 is used to prove that the periodic spectrum is symmetric if and only if the sequence of the gap lengths is symmetric, a conjecture, raised by several experts in the field. More precisely, denote by

$$\text{spec}_{\text{per}} L(\varphi) = (\lambda_k^\pm(\varphi))_{k \in \mathbb{Z}}$$

the periodic spectrum of $L(\varphi)$ when considered on the interval $[0, 2]$ where the numbers $\lambda_k^\pm(\varphi)$ are ordered so that

$$\lambda_k^-(\varphi) \leq \lambda_k^+(\varphi) < \lambda_{k+1}^-(\varphi)$$

and let $\gamma(\varphi) := (\gamma_k(\varphi))_{k \in \mathbb{Z}}$ be the sequence of gap lengths,

$$\gamma_k(\varphi) := \lambda_k^+(\varphi) - \lambda_k^-(\varphi).$$

In Section 4 we prove the following

THEOREM 1.3. — *For $\varphi \in L^2(S^1; \mathbb{C})$, the following assertions are equivalent:*

- (i) $\lambda_k^\pm(\varphi) = -\lambda_{-k}^\mp(\varphi)$ for any $k \geq 0$;
- (ii) $\gamma_k(\varphi) = \gamma_{-k}(\varphi)$ for any $k \geq 1$.

2. Symmetries and spectra

2.1. Periodic spectrum. — The periodic spectrum of the Zakharov-Shabat operator $L(\varphi)$ is given by

$$\text{spec}_{\text{per}} L(\varphi) := \left\{ \lambda \in \mathbb{C} \mid \exists F \in H_{\text{loc}}^1(\mathbb{R}; \mathbb{C}^2), F \not\equiv 0 \text{ with } L(\varphi)F = \lambda F \text{ and } F(x+2) = F(x), \forall x \in \mathbb{R} \right\}.$$

By [4], $\text{spec}_{\text{per}} L(\varphi)$ consists of a sequence of real numbers $(\lambda_k^\pm(\varphi))_{k \in \mathbb{Z}}$, which can be ordered in such a way that (for all $k \in \mathbb{Z}$)

$$(2.1) \quad \lambda_k^-(\varphi) \leq \lambda_k^+(\varphi) < \lambda_{k+1}^-(\varphi)$$

and $\lambda_k^\pm(\varphi) \sim k\pi$ for $|k|$ large. We have the following

PROPOSITION 2.1. — *Let $\varphi \in L^2(S^1; \mathbb{C})$. Then, for any $k \in \mathbb{Z}$,*

- (i) $\lambda_k^\pm(e^{i\alpha}\varphi) = \lambda_k^\pm(\varphi)$, $\lambda_k^\pm(\varphi) = \lambda_k^\pm(T_\tau\varphi)$ ($\forall \alpha \in \mathbb{R}$, $\tau \in \mathbb{R}$);
- (ii) $\lambda_k^\pm(\check{\varphi}) = -\lambda_{-k}^\mp(\varphi)$, $\lambda_k^\pm(\bar{\varphi}) = -\lambda_{-k}^\mp(\varphi)$.

Proof. — (i) For $\alpha \in \mathbb{R}$ arbitrary, define $V_\alpha = \begin{pmatrix} e^{-i\alpha/2} & 0 \\ 0 & e^{i\alpha/2} \end{pmatrix}$. One easily verifies that

$$(2.2) \quad L(e^{i\alpha}\varphi) = V_\alpha^{-1}L(\varphi)V_\alpha$$

and $L(T_\tau\varphi) = T_\tau L(\varphi)T_{-\tau}$. Thus both, $L(e^{i\alpha}\varphi)$ and $L(T_\tau\varphi)$, are unitarily equivalent to $L(\varphi)$ and the claimed statement follows. To prove (ii) notice that

$$(2.3) \quad L(-\check{\varphi}) = -W^{-1}L(\varphi)W$$

where W is the unitary operator defined by

$$W \begin{pmatrix} Y \\ Z \end{pmatrix} := \begin{pmatrix} \check{Y} \\ \check{Z} \end{pmatrix}, \quad \text{with} \quad \begin{pmatrix} Y \\ Z \end{pmatrix} \in L^2_{\text{loc}}(\mathbb{R}; \mathbb{C}^2).$$

Thus

$$(2.4) \quad \text{spec}_{\text{per}} L(-\check{\varphi}) = -\text{spec}_{\text{per}} L(\varphi).$$

Combining (2.4) and (i) we obtain $\lambda_k^\pm(\check{\varphi}) = -\lambda_{-k}^\mp(\varphi)$ for all $k \in \mathbb{Z}$. Consider $\lambda \in \text{spec}_{\text{per}} L(\varphi)$ and choose $F \in H^1_{\text{loc}}(\mathbb{R}; \mathbb{C}^2)$, satisfying $F(x+2) = F(x)$ for all $x \in \mathbb{R}$ and $L(\varphi)F = \lambda F$. As λ is real, $L(-\bar{\varphi})\bar{F} = -\lambda\bar{F}$ and thus $-\lambda \in \text{spec}_{\text{per}} L(-\bar{\varphi})$. Combined with (i), this leads to $\lambda_k^\pm(\bar{\varphi}) = -\lambda_{-k}^\mp(\varphi)$ for all $k \in \mathbb{Z}$. $\square$

2.2. Dirichlet spectra and divisors. — To study properties of the Dirichlet spectra it is convenient to consider the AKNS operator $H(\varphi)$ instead of $L(\varphi)$. Let

$$F_j(x, \lambda; \varphi) := \begin{pmatrix} Y_j(x, \lambda; \varphi) \\ Z_j(x, \lambda; \varphi) \end{pmatrix}, \quad j = 1, 2,$$

be the fundamental solutions of $H(\varphi)$, i.e. the solutions to $HF = \lambda F$ such that

$$F_1(0, \lambda; \varphi) = \begin{pmatrix} 1 \\ 0 \end{pmatrix}, \quad F_2(0, \lambda; \varphi) = \begin{pmatrix} 0 \\ 1 \end{pmatrix}.$$

For each $x \in \mathbb{R}$ and $\varphi \in L^2(S^1; \mathbb{C})$, $F_1(x, \lambda; \varphi)$ and $F_2(x, \lambda; \varphi)$ are entire functions of λ . The two Dirichlet spectra are defined as follows

$$(2.5) \quad \text{spec}_{\text{Dir}}^+ L(\varphi) = \{\lambda \in \mathbb{C} \mid Z_1(1, \lambda; \varphi) = 0\},$$

$$(2.6) \quad \text{spec}_{\text{Dir}}^- L(\varphi) = \{\lambda \in \mathbb{C} \mid Y_2(1, \lambda; \varphi) = 0\}.$$

It is proved in [4] that $\text{spec}_{\text{Dir}}^+ L(\varphi)$, resp. $\text{spec}_{\text{Dir}}^- L(\varphi)$, consists of simple, real eigenvalues $(\mu_k(\varphi))_{k \in \mathbb{Z}}$, resp. $(\nu_k(\varphi))_{k \in \mathbb{Z}}$. The numerotation is chosen in such a way that $(\mu_k(\varphi))_{k \in \mathbb{Z}}$ and $(\nu_k(\varphi))_{k \in \mathbb{Z}}$ are strictly increasing satisfying $\mu_k(\varphi) \sim k\pi$ and $\nu_k(\varphi) \sim k\pi$ for $|k|$ large. Further introduce the function $\delta(\lambda; \varphi)$, defined by

$$(2.7) \quad \delta(\lambda; \varphi) = Z_2(1, \lambda; \varphi) - Y_1(1, \lambda; \varphi).$$